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Particle-Wave Properties

Remarks on Neutron Interferometry

Quantum State Preparation and Measurement

Magnetic Noise Dephasing and Decoherencing

Contextuality

Topological Phases

Ultracold Neutrons and Phase Space Transformation

Particle Properties

$$m = 1.674928(1) \times 10^{-27} \text{ kg}$$

$$s = \frac{1}{2} \hbar$$

$$\mu = -9.6491783(18) \times 10^{-27} \text{ J/T}$$

$$\tau = 887(2) \text{ s}$$

$$R = 0.7 \text{ fm}$$

$$\alpha = 12.0 (2.5) \times 10^{-4} \text{ fm}^3$$

u - d - d - quark structure

m ... mass, s ... spin, μ ... magnetic moment, τ ... β -decay lifetime, R ... (magnetic) confinement radius, α ... electric polarizability; all other measured quantities like electric charge, magnetic monopole and electric dipole moment are compatible with zero

CONNECTION

de Broglie

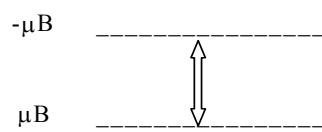
$$\lambda_B = \frac{\hbar}{m \cdot v}$$

Schrödinger

$$H\psi(\vec{r},t) = i\hbar \frac{\delta\psi(\vec{r},t)}{\delta t}$$

&

boundary conditions



Wave Properties

$$\lambda_c = \frac{\hbar}{m \cdot c} = 1.319695(20) \times 10^{-15} \text{ m}$$

For thermal neutrons
 $= 1.8 \text{ \AA}, 2200 \text{ m/s}$

$$\lambda_B = \frac{\hbar}{m \cdot v} = 1.8 \times 10^{-10} \text{ m}$$

$$\Delta_c = \frac{1}{2\delta k} \cong 10^{-8} \text{ m}$$

$$\Delta_p = v \cdot \Delta t \cong 10^{-2} \text{ m}$$

$$\Delta_d = v \cdot \tau = 1.942(5) \times 10^6 \text{ m}$$

$$0 \leq \chi \leq 2\pi (4\pi)$$

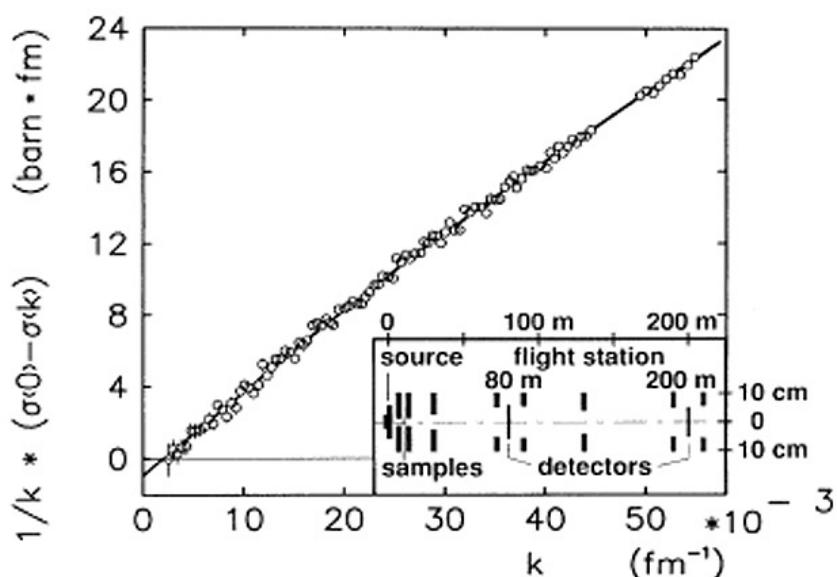
λ_c ... Compton wavelength, λ_B ... deBroglie wavelength, Δ_c ... coherence length, Δ_p ... packet length, Δ_d ... decay length, δk ... momentum width, Δt ... chopper opening time, v ... group velocity, χ ... phase.

$$\vec{d} = \alpha \vec{E}$$

$$V = -\frac{1}{2} \vec{d} \vec{E} = -\frac{1}{2} \alpha E^2$$

$$\sigma_s(k) = \sigma_s(0) + a \cdot k + b \cdot k^2 + O(k^4)$$

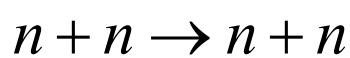
\Rightarrow Coulomb field of Pb-208



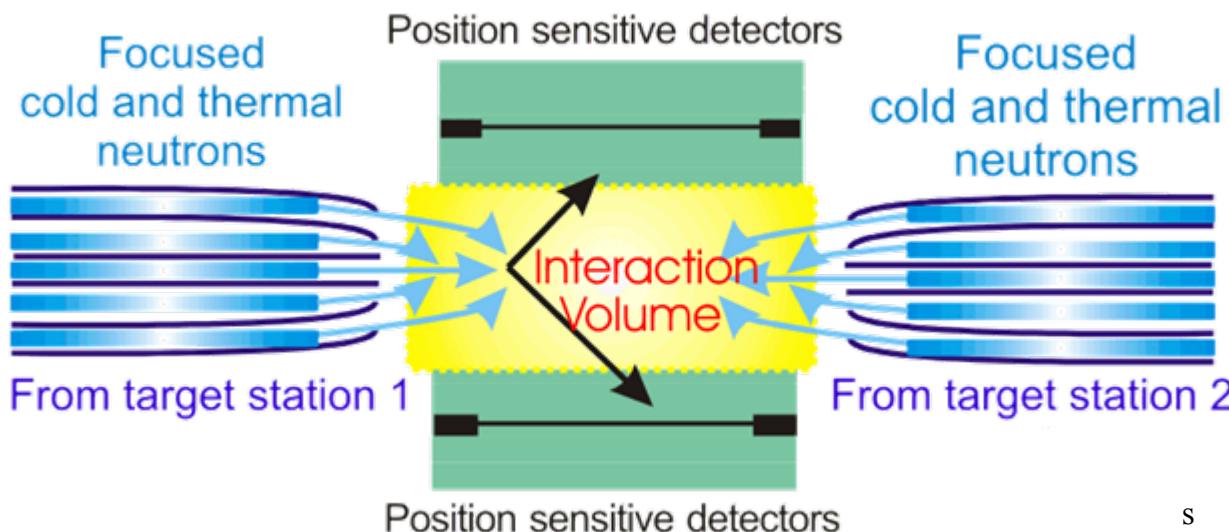
$$= (1.20 \pm 0.15 \pm 0.20) \times 10^{-3} \text{ fm}^3$$

J. Schmiedmayer, P. Rehs,
J.A. Harvey, N. W. Hill,
Phys. Rev. Lett. 66 (1991) p.1015

Charge dependence and charge symmetry



$$a_{np}^s \approx a_{pp}^s \approx a_{nn}^s \quad ?$$



$$a_{nn}^t \equiv 0 \quad ?$$

Experiment:

$$a_{np}^s = -23.678(28) \text{ fm}$$

$$a_{pp}^s = -17.3(1.0) \text{ fm}$$

$$a_{nn}^s = -16.8(1.3) \text{ fm}$$

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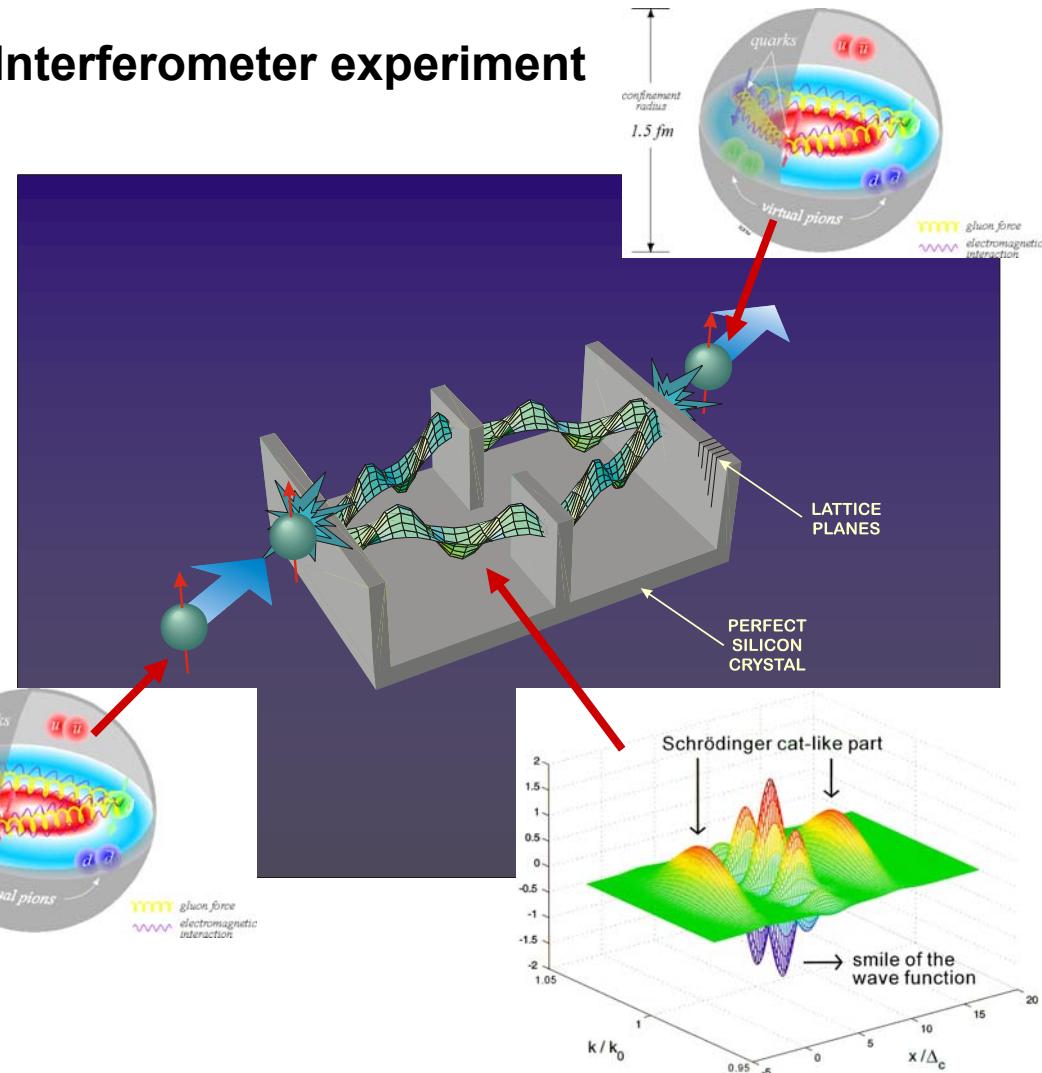
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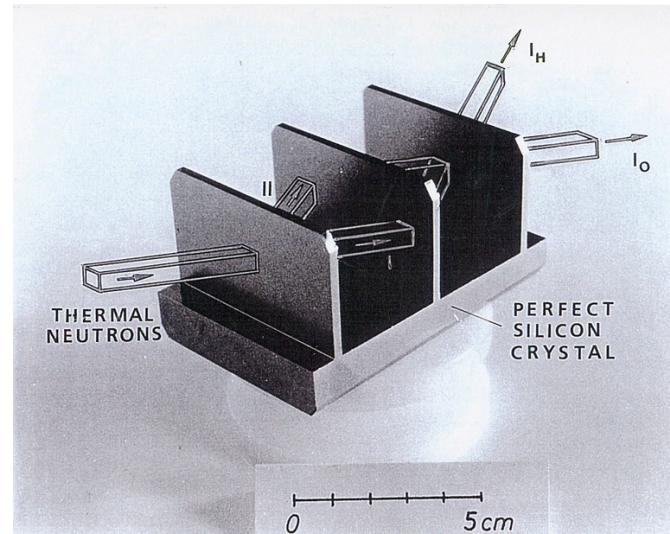


Interferometer experiment



$$I_0 \propto |\psi_0^I + \psi_0^{II}|^2 \propto A + B \cos \chi$$

$$\chi = \int \vec{k} d\vec{s} = (1-n)kD_{\text{eff}} \equiv -Nb_c\lambda D_{\text{eff}} = \Delta \cdot k = \Delta k \cdot D_{\text{eff}}$$

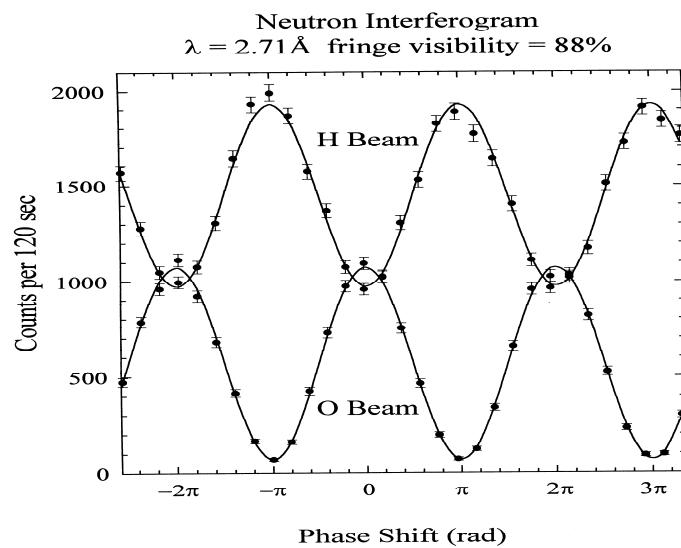


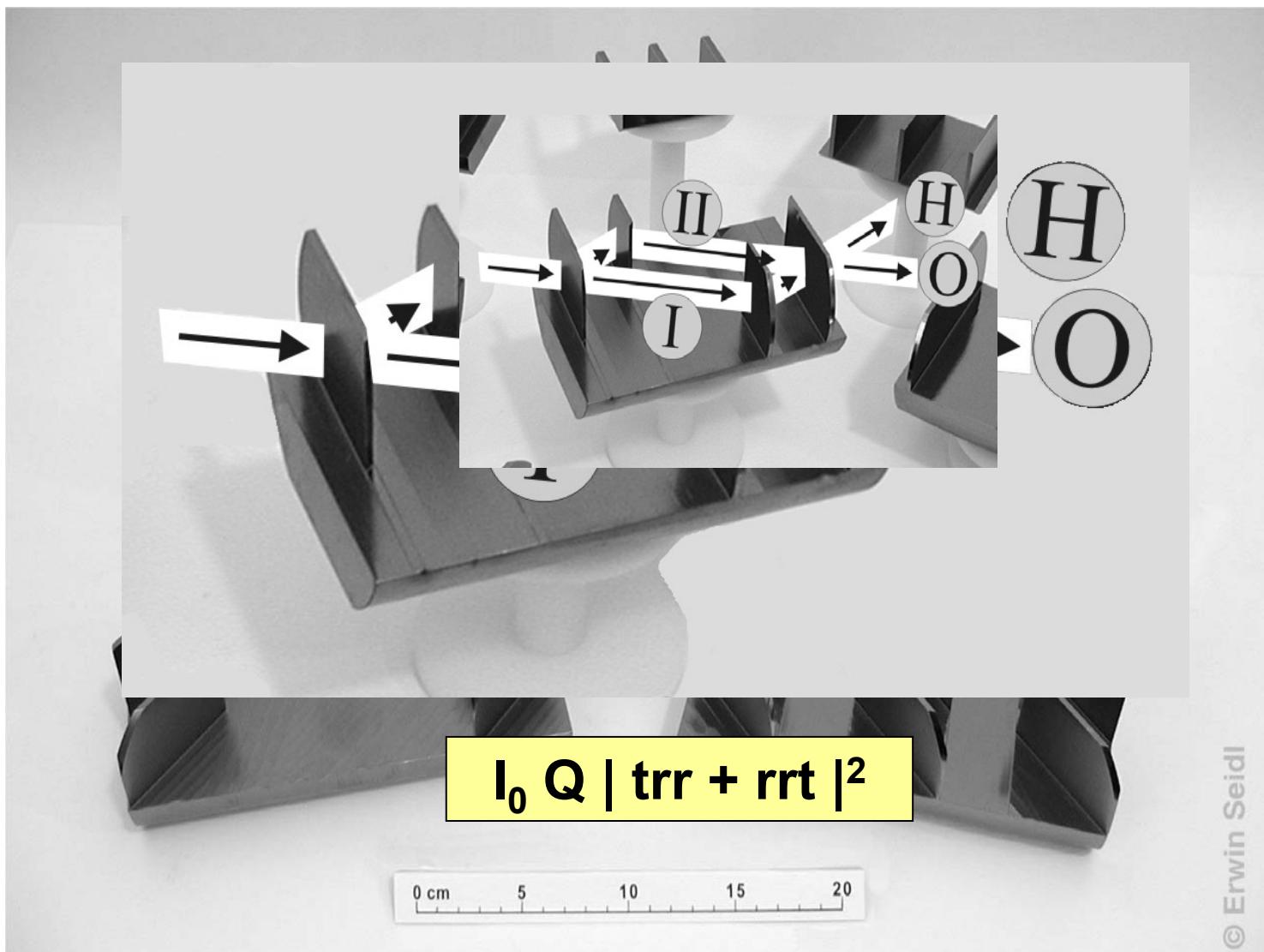
Self interference

(phase space density $\sim 10^{-14}$)

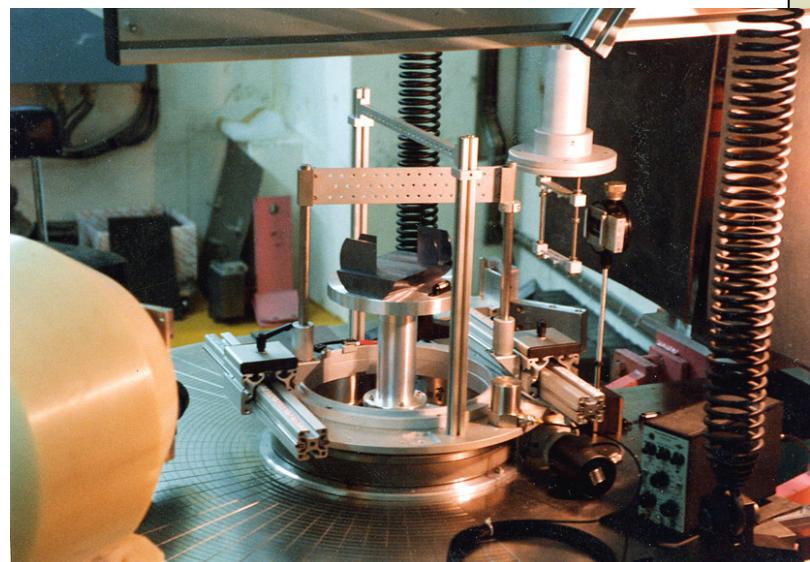
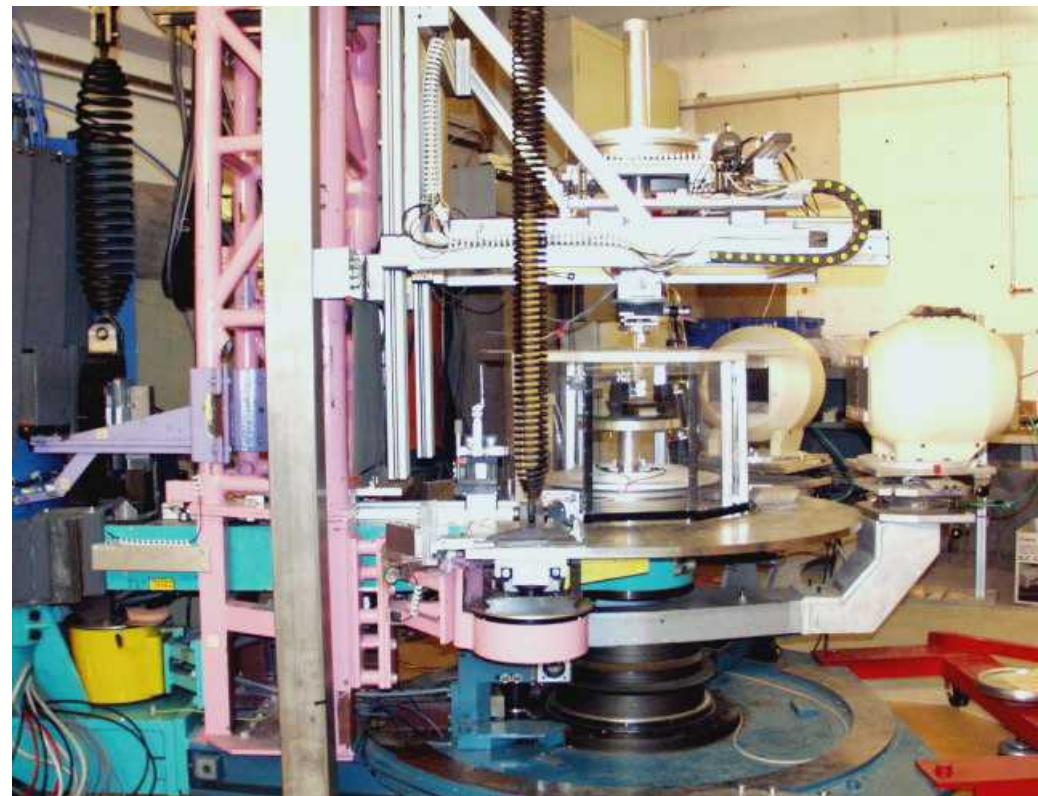
Efficiency of detectors, polarizers, flippers >99%

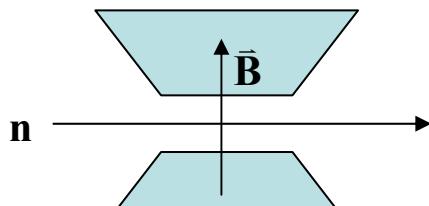
H. Rauch, W. Treimer, U. Bonse, Phys.Lett. A47 (1974) 369





© Erwin Seidl





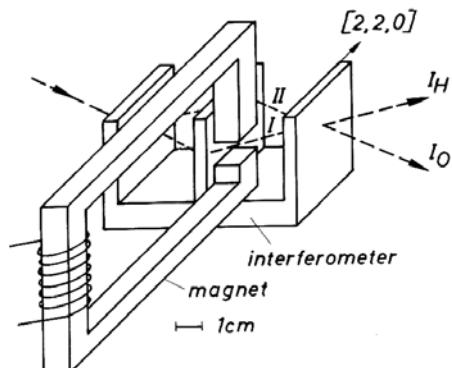
$$\begin{aligned}\psi^{\text{II}} &\rightarrow \psi^{\text{I}} e^{-iHt/\hbar} = \psi^{\text{I}} e^{-i\bar{\mu}\bar{B}t/\hbar} \\ &= \psi^{\text{I}} e^{-i\mu\bar{\sigma}\bar{B}t/\hbar} = \psi^{\text{I}} e^{-i\bar{\sigma}\alpha/2}\end{aligned}$$

$$|\alpha| = \frac{2\mu B t}{\hbar} = g B t \approx \frac{2\mu B \ell}{\hbar v} \dots \text{Larmor angle}$$

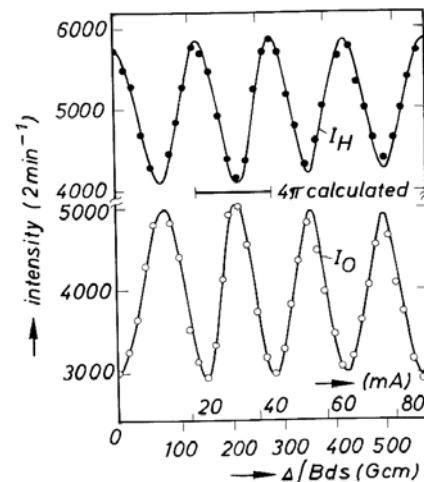
$$\psi(\alpha) = \begin{pmatrix} e^{-\alpha/2} & 0 \\ 0 & e^{\alpha/2} \end{pmatrix} \begin{pmatrix} \psi_{\uparrow}^{\text{I}}(0) \\ \psi_{\downarrow}^{\text{I}}(0) \end{pmatrix}$$

Theory: H.J.Bernstein, Phys.Rev.Lett. 18(1967)1102,
Y.Aharonov, L.Susskind, Phys.Rev. 158(1967)1237

$$\begin{aligned}\psi(2\pi) &= -\psi(0) \\ \psi(4\pi) &= \psi(0)\end{aligned}$$



$$I_0 \propto |\psi_0^{\text{I}}(0) + \psi_0^{\text{I}}(\alpha)|^2 = \frac{I_0(0)}{2} \left(1 + \cos \frac{\alpha}{2} \right)$$



Experiment: H.Rauch, A.Zeilinger, G.Badurek, A.Wilfing, W.Bauspiess, U.Bonse, Phys.Lett. 54A(1975)425
S.A.Werner, R.Colella, A.W.Overhauser, C.F.Eagen, Phys.Rev.Lett. 35(1975)1053
A.G.Klein, G.I.Opat, Phys.Rev. D11(1976)523
E.Klempt, Phys.Rev. D13(1975)3125
M.E.Stoll, E.K.Wolff, M.Mehring, Phys.Rev. A17(1978)1561

$$H = \frac{p^2}{2m_i} - G \frac{Mm_g}{r} - \vec{\omega} \cdot \vec{L} \approx \frac{p^2}{2m_i} + m_g g z - \vec{\omega} \cdot \vec{L} + V_0$$

$$\vec{L} = \vec{r} \times \vec{p}$$

$$\vec{r} = \vec{r}_0 + \vec{v}_0 t - \frac{1}{2} \vec{g} t^2 + \frac{1}{3} t^3 \vec{\omega} \times \vec{g}$$

$$\beta = \oint \vec{k} d\vec{r} = \frac{m_i}{\hbar} \oint \dot{\vec{r}} d\vec{r} = -2\pi m_i m_g \frac{A g}{h^2} \lambda_0 \sin \phi + \frac{4\pi m_i}{h} \vec{\omega} \times \vec{A}$$

A ... area enclosed by the coherent beams

ϕ ... colatitude angle

$$\beta = -q_{grav} \sin \phi + q_s \sin \epsilon$$



$$2\pi m_i m_g g \lambda A / h^2$$

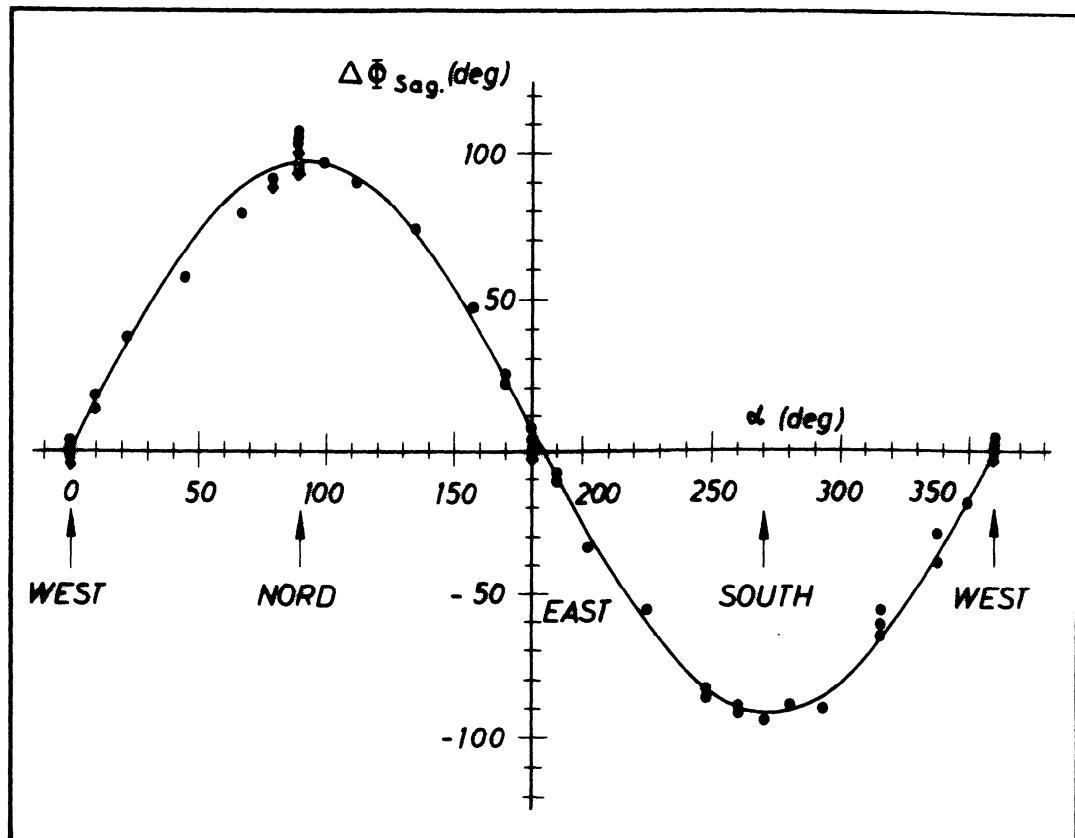
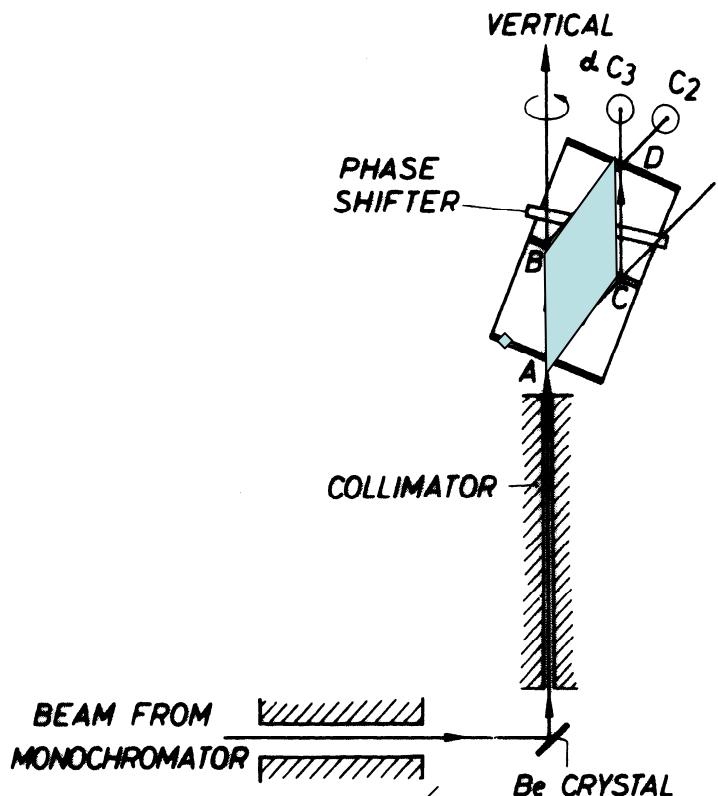


$$4\pi m_i \omega A \sin \phi / h$$

- **EXPERIMENT:**
- Colella, Overhauser, Werner, 1975
- Werner, Staudenmann, Colella, Overhauser, 1975, 1978
- Bonse & Wroblewski 1983
- Atwood, Shull, Arthur 1984

- **THEORY:**
- Page 1975;
- Anandan 1977
- Stodolsky 1979
- Audretsch & Lommerzahl 1982

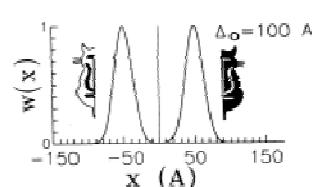
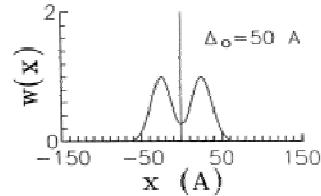
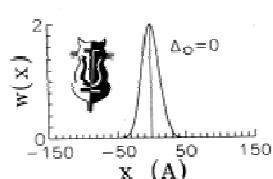
Earth Rotation Effect



J.L.Staudenmann, S.A.Werner, R.Colella, A.W.Overhauser,
 PR/A21(1980)1419

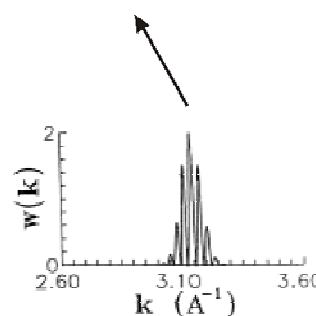
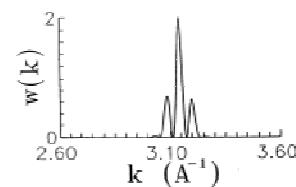
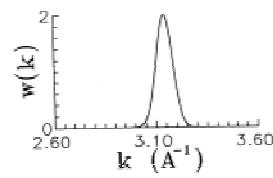
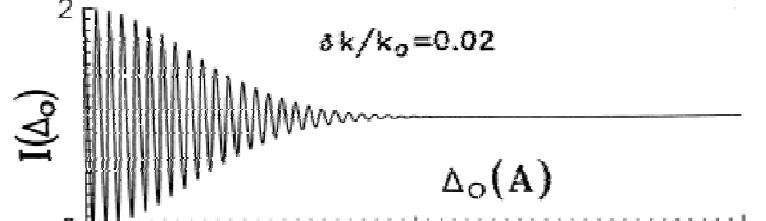
- ***Quantum State Preparation and Measurements***

Spatial distribution



Coherent state

Non-classical state



Momentum distribution



Schrödinger Equation:

$$-\frac{\hbar^2}{2m} \Delta \psi(\vec{r}, t) + V(\vec{r}, t)\psi(\vec{r}, t) = i\hbar \frac{\partial \psi(\vec{r}, t)}{\partial t}$$

Partial waves fill the whole space

Wave Function (Eigenvalue solution in free space):

$$\Psi(\vec{r}, t) = (2\pi)^{-3/2} \int \psi(\vec{k}, t) e^{i(\vec{k} \cdot \vec{r} - \omega t)} d^3 \vec{k}$$

Spatial distribution:

$$\rho(\vec{r}, t) = |\psi(\vec{r}, t)|^2$$

Coherence Function:

Stationary situation: ($\tau = 0$):

$$\Gamma(\vec{\Delta}) = \langle \psi(0) \psi(\vec{\Delta}) \rangle = (2\pi)^{3/2} \int g(\vec{k}) e^{i\vec{k}\vec{\Delta}} d^3 k$$

$$\tau = t - t'$$

$$\vec{\Delta} = \vec{r} - \vec{r}'$$

↓

and others (Wigner function etc.)

Momentum distribution:

$$g(\vec{k}, t) = |\psi(\vec{k}, t)|^2$$

Definition:

$$W(k, x) = \frac{1}{4\pi} \int e^{ik\Delta} \psi^*(x + \frac{\Delta}{2}) \psi(x - \frac{\Delta}{2}) d\Delta$$

Properties: $\int W(k, x) dk = |\psi(x)|^2$

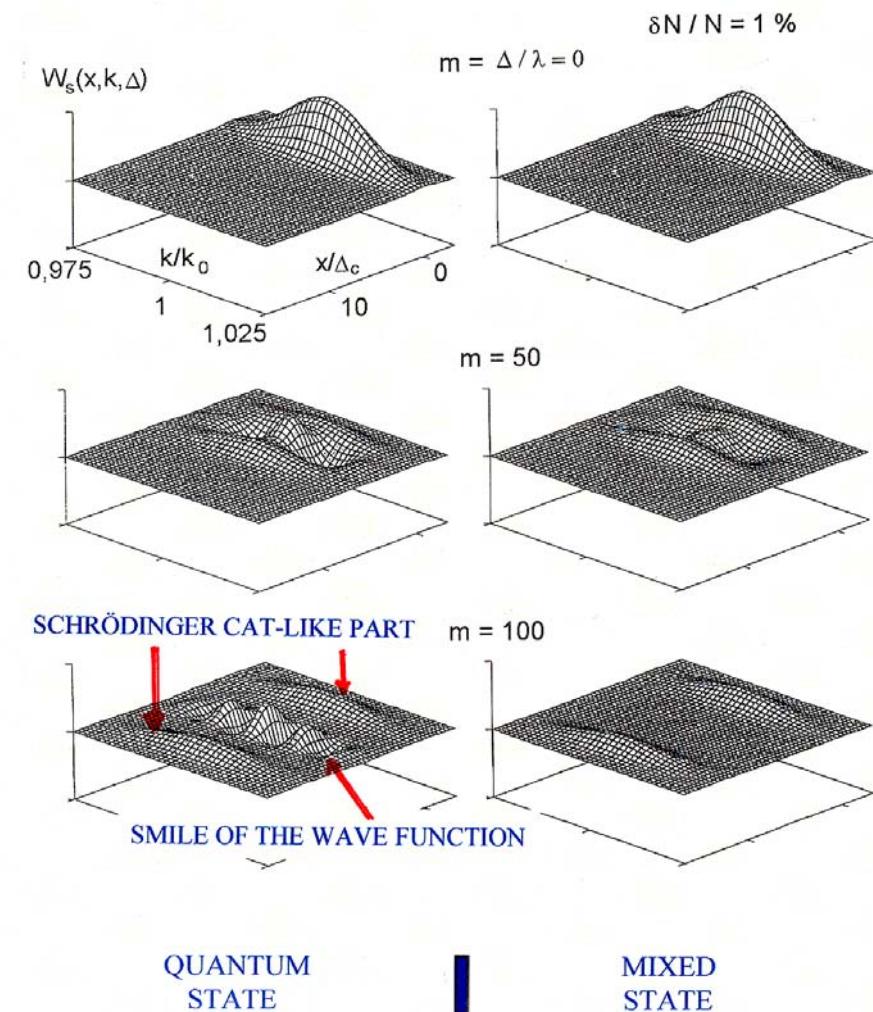
$$\int W(k, x) dx = |\psi(k)|^2$$

Interferometric Gaussian packets:

$$\psi^{I,II}(x) = (4\pi\delta x^2)^{-1/4} \exp\left[-x^2/2\delta x^2 + ixk_0\right]$$

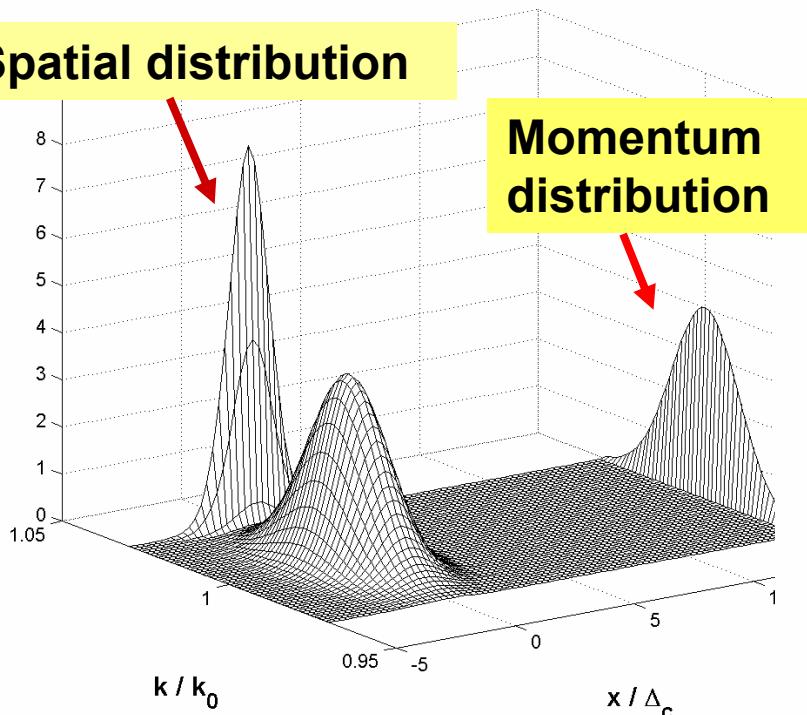
$$\psi(x) = \psi^I(x) + \psi^{II}(x + \Delta_0)$$

H. Rauch, M. Suda, Appl.Phys.B60 (1994) 181



$\Delta = 0 \text{ nm}$

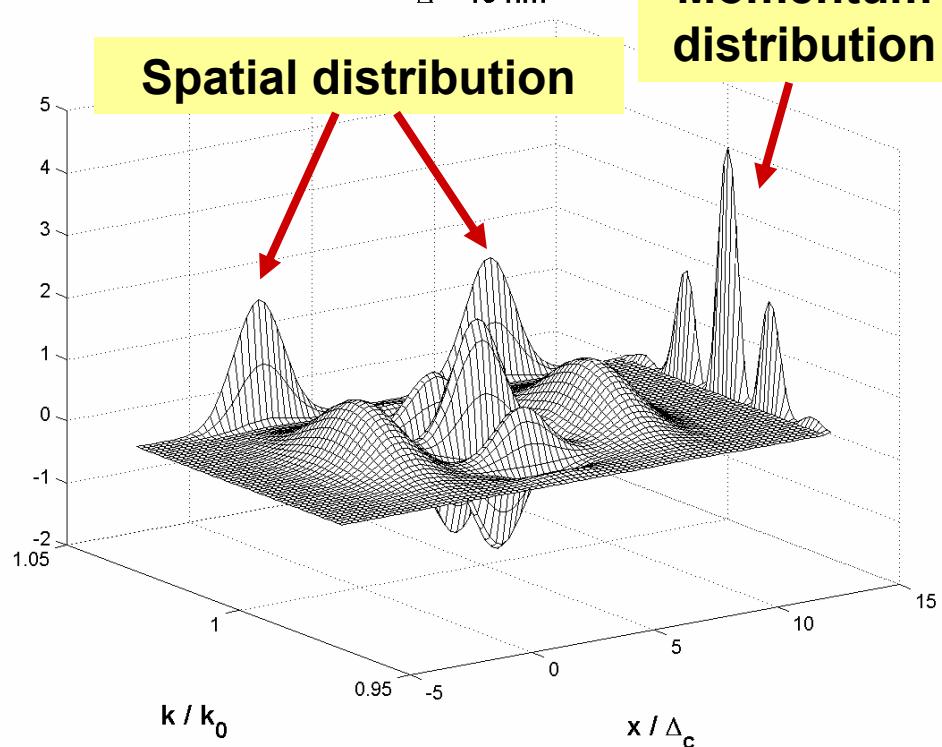
Spatial distribution



Zero order
(coherent state)

$\Delta = 15 \text{ nm}$

Momentum distribution



High order
(non-classical state)

Schrödinger Equation:

$$-\frac{\hbar^2}{2m} \Delta \psi(\vec{r}, t) + V(\vec{r}, t) \psi(\vec{r}, t) = i\hbar \frac{\partial \psi(\vec{r}, t)}{\partial t}$$

Wave Function (Eigenvalue solution in free space):

$$\Psi(\vec{r}, t) = (2\pi)^{-3/2} \int \psi(\vec{k}, t) e^{i(\vec{k} \cdot \vec{r} - \omega t)} d^3 k$$

Spatial distribution: Momentum distribution:

$$\rho(\vec{r}, t) = |\psi(\vec{r}, t)|^2 \quad g(\vec{k}, t) = |\psi(\vec{k}, t)|^2$$

Coherence Function: $\bar{\Delta} = \vec{r} - \vec{r}'; \quad \tau = t - t'$

Stationary situation: ($\tau = 0$):

$$\Gamma(\bar{\Delta}) = \langle \psi(0) \psi(\bar{\Delta}) \rangle = (2\pi)^{3/2} \int g(\vec{k}) e^{i\vec{k}\bar{\Delta}} d^3 k$$

Wigner Function:

$$W(x, k) = (2\pi)^{-1} \int e^{ikx'} \psi^*\left(x + \frac{x'}{2}\right) \psi\left(x - \frac{x'}{2}\right) dx'$$

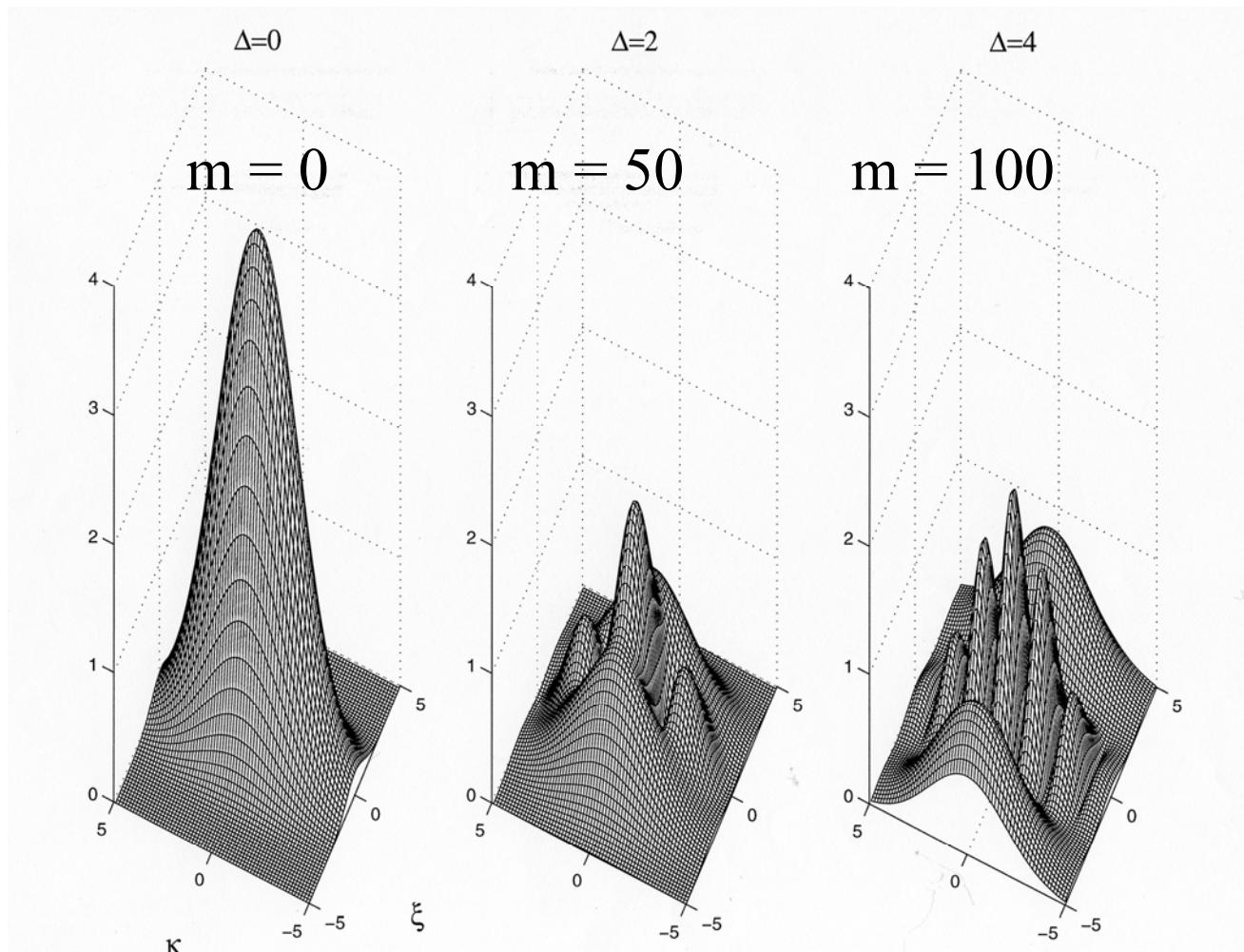
Q-Function (Husimi-Function):

$$Q(x, k) = \iint W(x', k') g(x - x', k - k', \gamma) dx' dk'$$

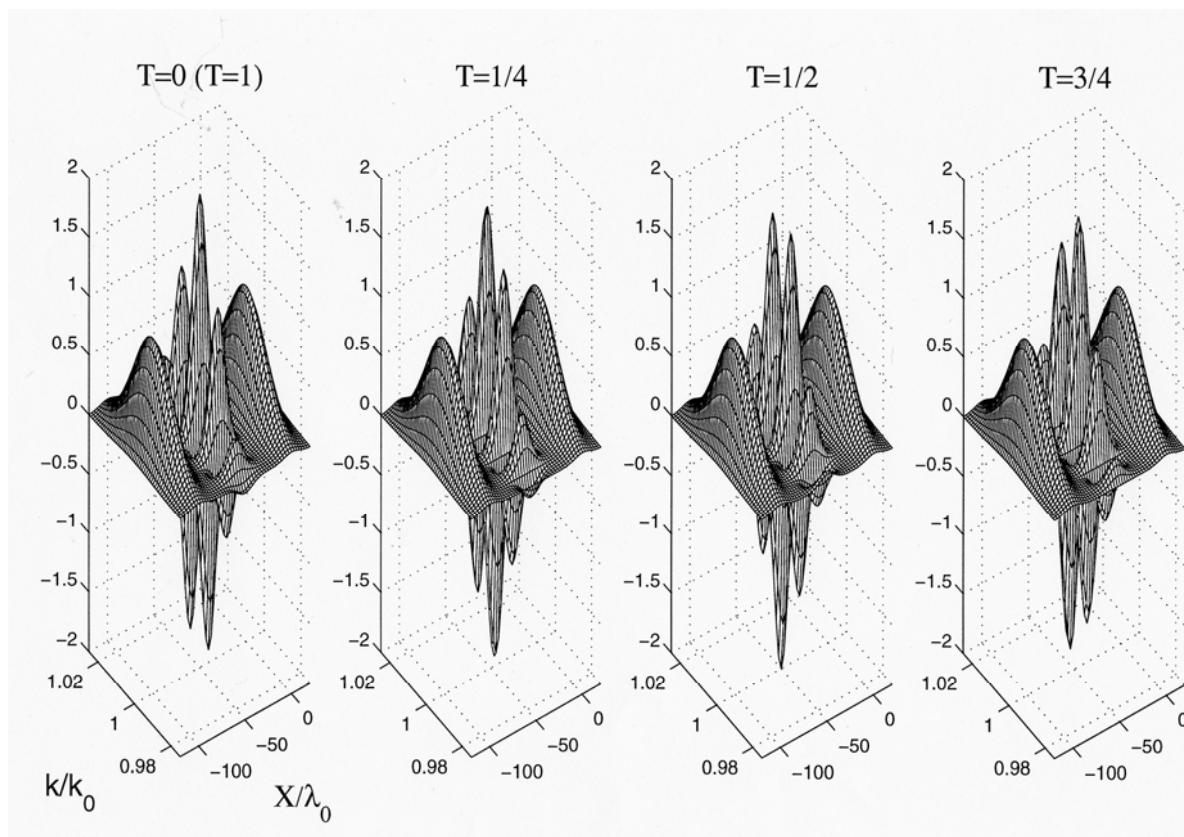
$$g(x - x', k - k') \propto \exp\left[-\frac{(x - x')^2}{\gamma} - \gamma(k - k')^2\right]$$

Weyl Function:

$$\tilde{W}(X, K) = \iint W(x, k) e^{-i(Kx - kX)} dk dx$$



$m = 100$



H.Rauch, M.Suda, Lect.Notes in Physics, Springer 2002

via Wigner functions

$$\begin{aligned} W(x, k) &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{ikx'} \psi^*(x + \frac{x'}{2}) \psi(x - \frac{x'}{2}) dx' \\ &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{ikx'} \langle x + \frac{x'}{2} | \hat{\rho} | x - \frac{x'}{2} \rangle dx' \end{aligned}$$

$$(x \doteq \Delta = -Nb_c \lambda^2 D / 2\pi)$$

Quadrature operator

$$\hat{X}_\Theta = k_0 \hat{x} \cos \Theta + \frac{\hat{k}}{k_0} \sin \Theta$$

Quadrature Wigner function

$$W(X_\Theta) = \frac{\hbar}{2\pi} \int_{-\infty}^{+\infty} dt e^{-itX_\Theta} \int_{-\infty}^{+\infty} dx \int_{-\infty}^{+\infty} dk e^{it(k_0x \cos \Theta + k \sin \Theta / k_0)} W(x, k)$$

Radon transformation

$$W(x, k) = \frac{1}{4\pi^2 \hbar} \int_{-\infty}^{+\infty} dt |t| \int_0^\pi d\Theta \int_{-\infty}^{+\infty} dX_\Theta e^{it(X_\Theta - k_0 x \cos \Theta - k \sin \Theta / k_0)} W(X_\Theta)$$

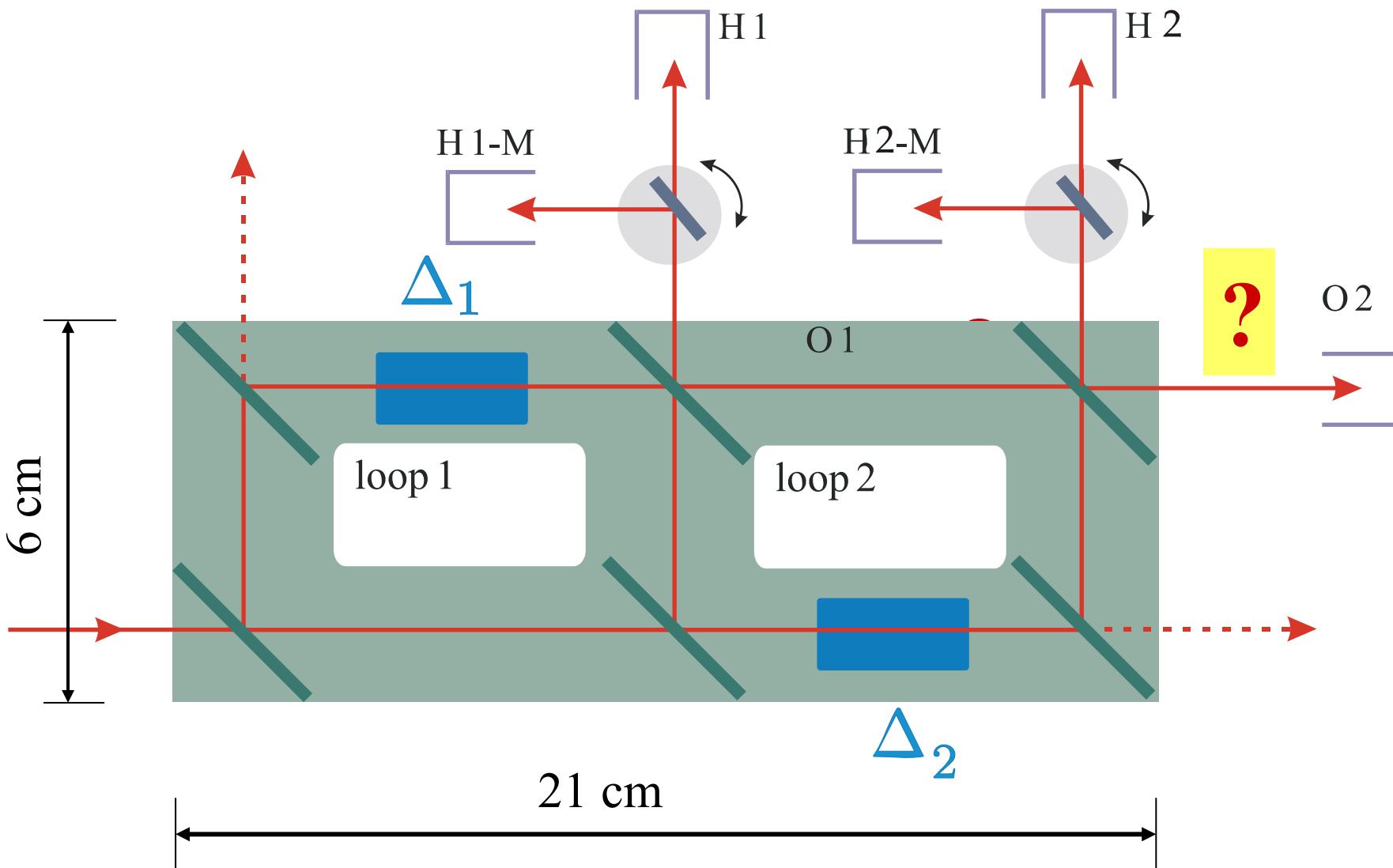
Neutron interferometry case (Gaussian packets)

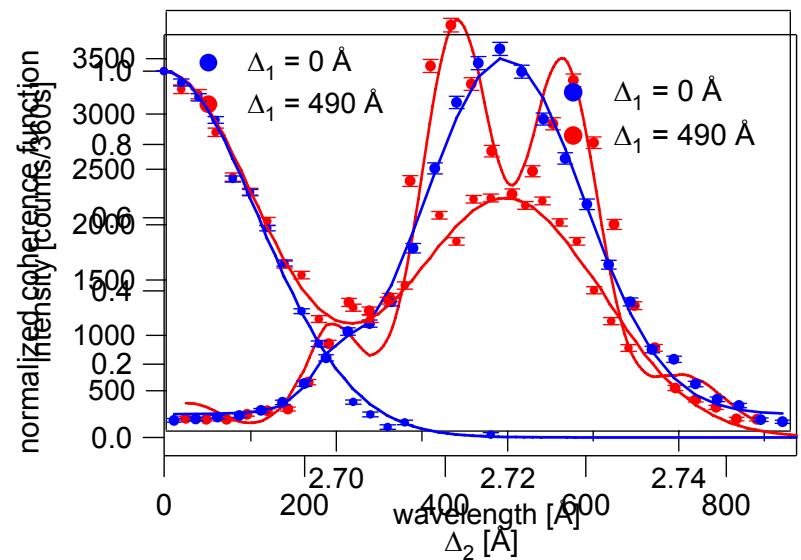
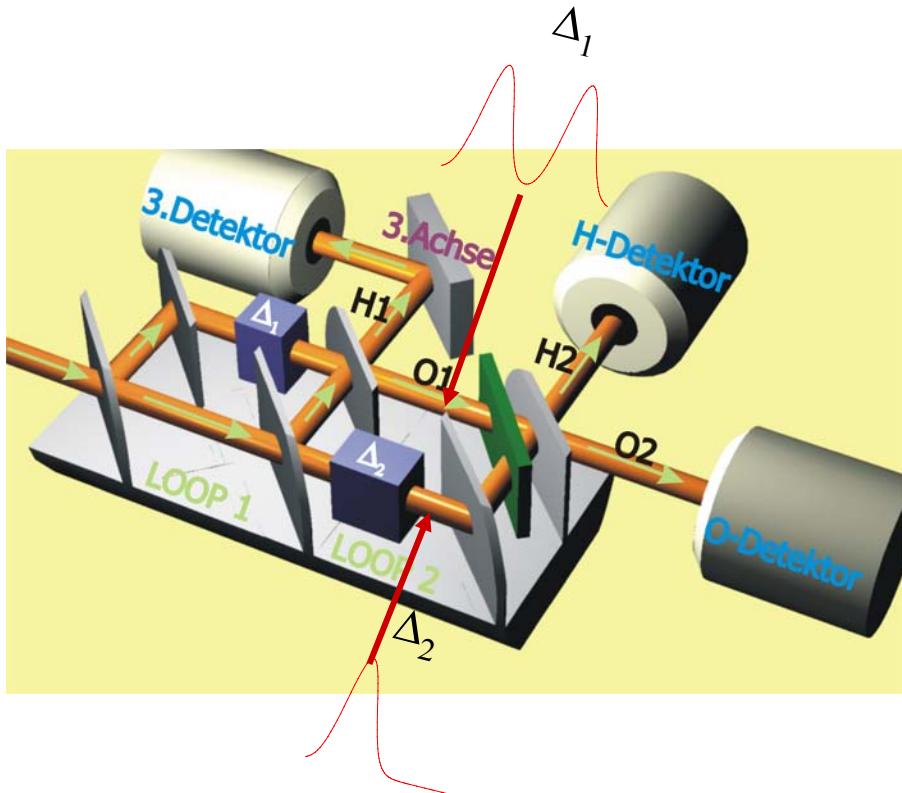
$$W(x, k, \Delta) = W(x, k) + W(x + \Delta, k) + 2 \cos k\Delta W(x + \frac{\Delta}{2}, k)$$

$$\begin{aligned} W(X_\Theta) &= \frac{\hbar}{2\pi} \sqrt{\frac{\pi}{b}} \left\{ e^{-(X_\Theta - \sin \Theta)^2 / 4b} \right. \\ &\quad + e^{-(X_\Theta - \sin \Theta + k_0 \Delta \cos \Theta)^2 / 4b} \\ &\quad + 2e^{-(X_\Theta - \sin \Theta + k_0 \Delta \cos \Theta)^2 / 4b} \\ &\quad \cdot e^{-\sigma^2 (k_0 \Delta)^2 / 2 + \sigma^4 (k_0 \Delta)^2 \sin^2 \Theta / 4b} \\ &\quad \cdot \left. \cos[(k_0 \Delta) + (X_\Theta - \sin \Theta + k_0 \Delta \cos \Theta / 2) \sigma^2 (k_0 \Delta) \sin \Theta / 2b] \right\} \end{aligned}$$

with: $b = \cos^2 \Theta / 8\sigma^2 + \sigma^2 \sin^2 \Theta / 2$

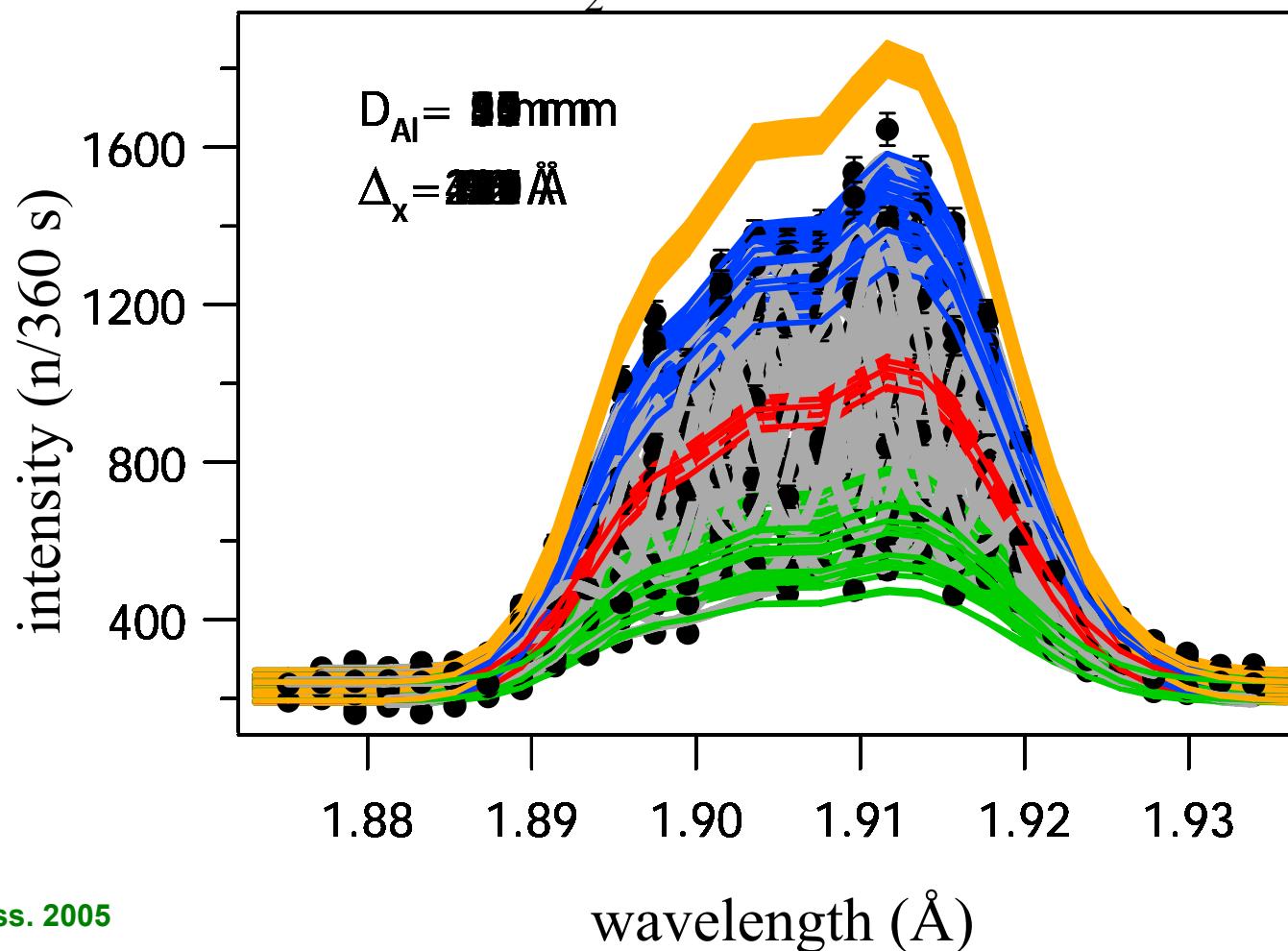
$\sigma = \delta k / k_0$





M.Baron, H.Rauch, M.Suda, J.Opt.B5 (2003) S341

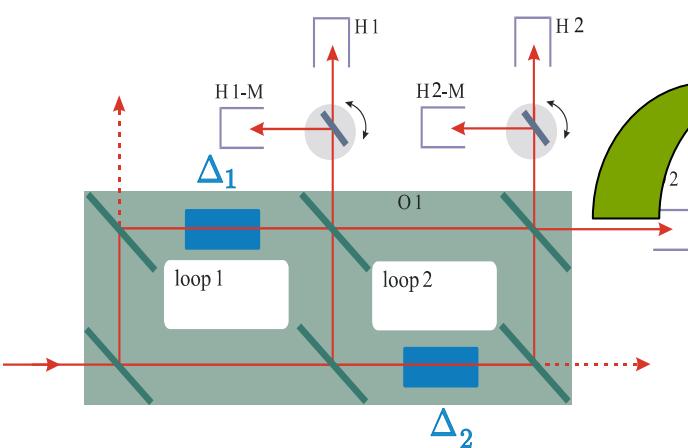
$$I_0(\vec{k}) = \frac{1}{2} g(\vec{k}) \cdot (1 + C \cos \Delta \vec{k})$$



M.Baron, Diss. 2005

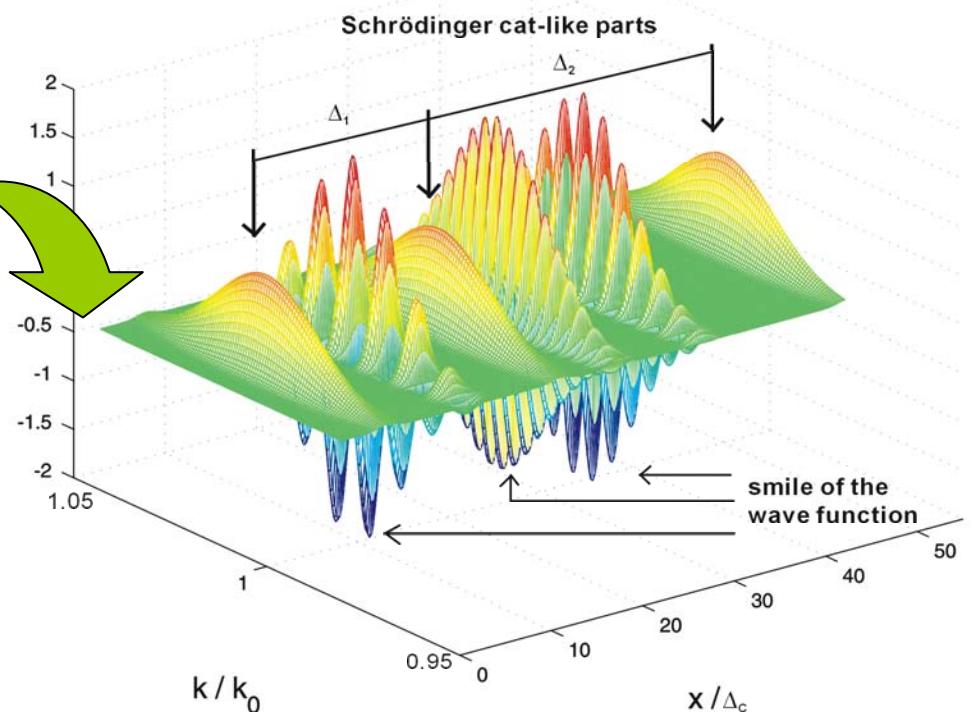
wavelength (\AA)

TWO LOOP INTERFEROMETER

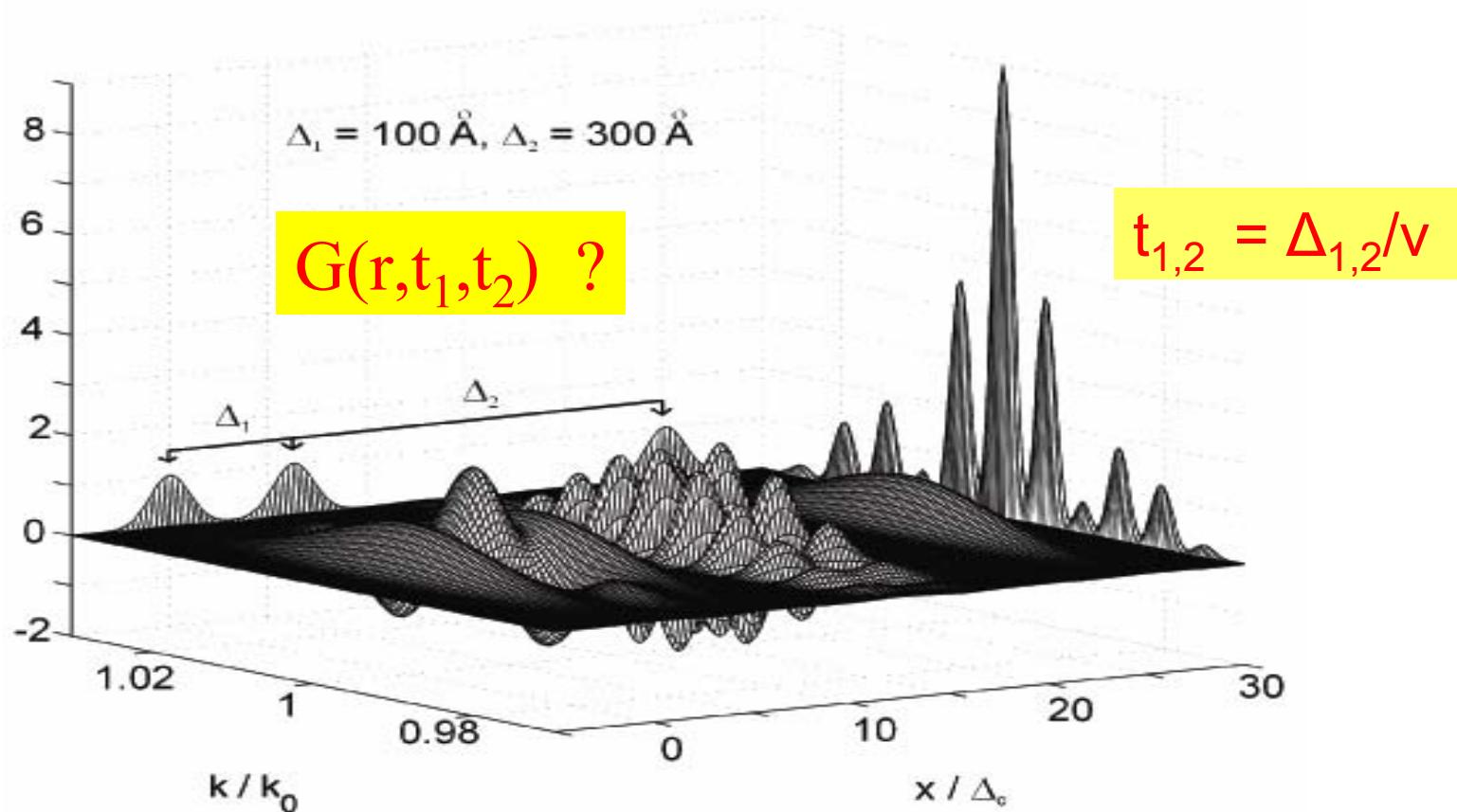


Wigner Function of a Four-Plate Interferometer

$W_s(x, k, \Delta_1=300 \text{ Å}, \Delta_2=500 \text{ Å}), \Delta_c = 15.9 \text{ Å}, \delta k / k_0 = 1 \%$



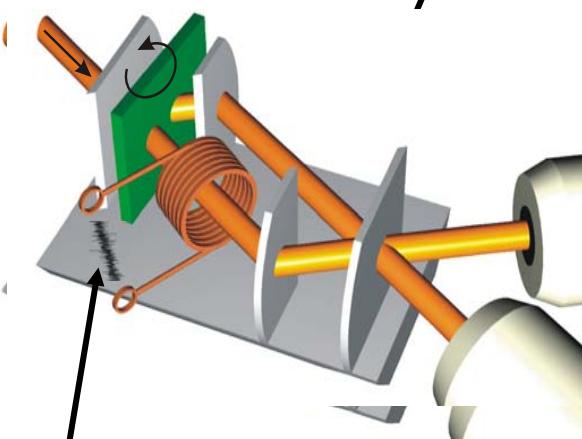
Quantum state engineering !!!



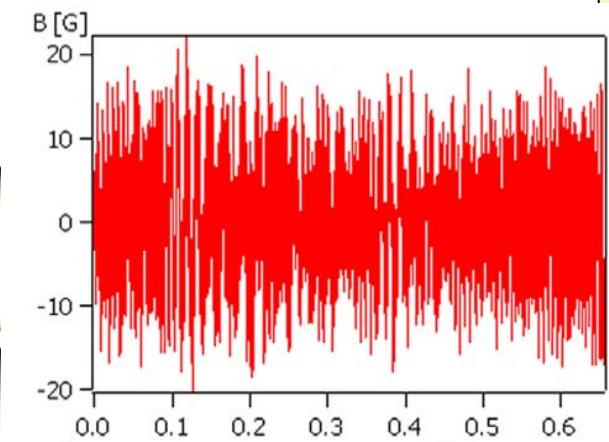
H.Rauch, M.Suda 2001

- ***Magnetic Noise Dephasing or Decoherencing***

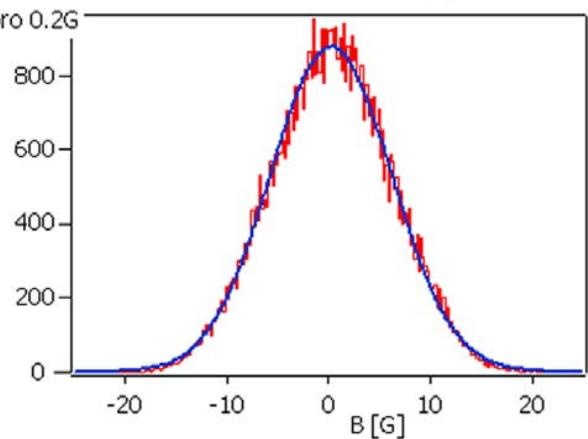
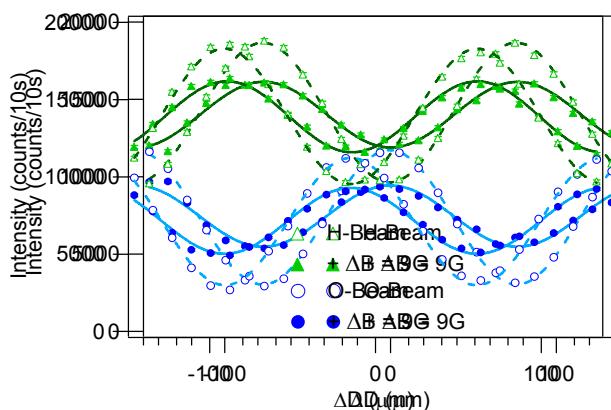
Magnetic noise field



Magnetic noise field

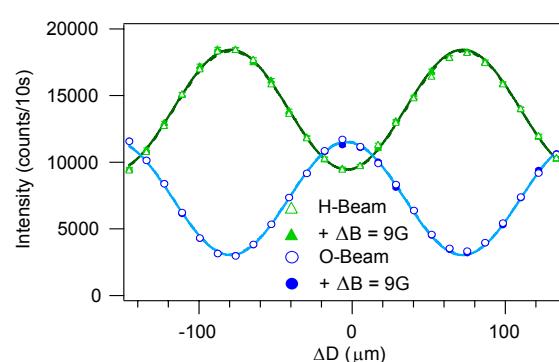
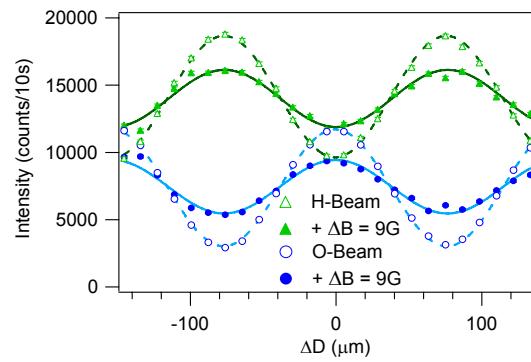
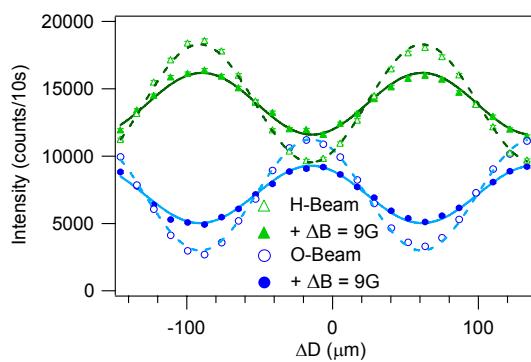
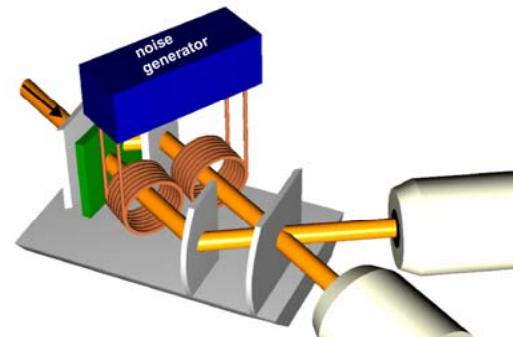
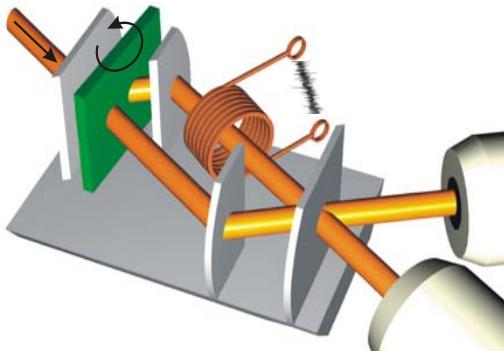
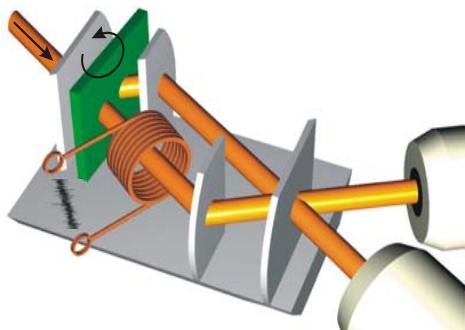


Anzahl
pro 0.2G

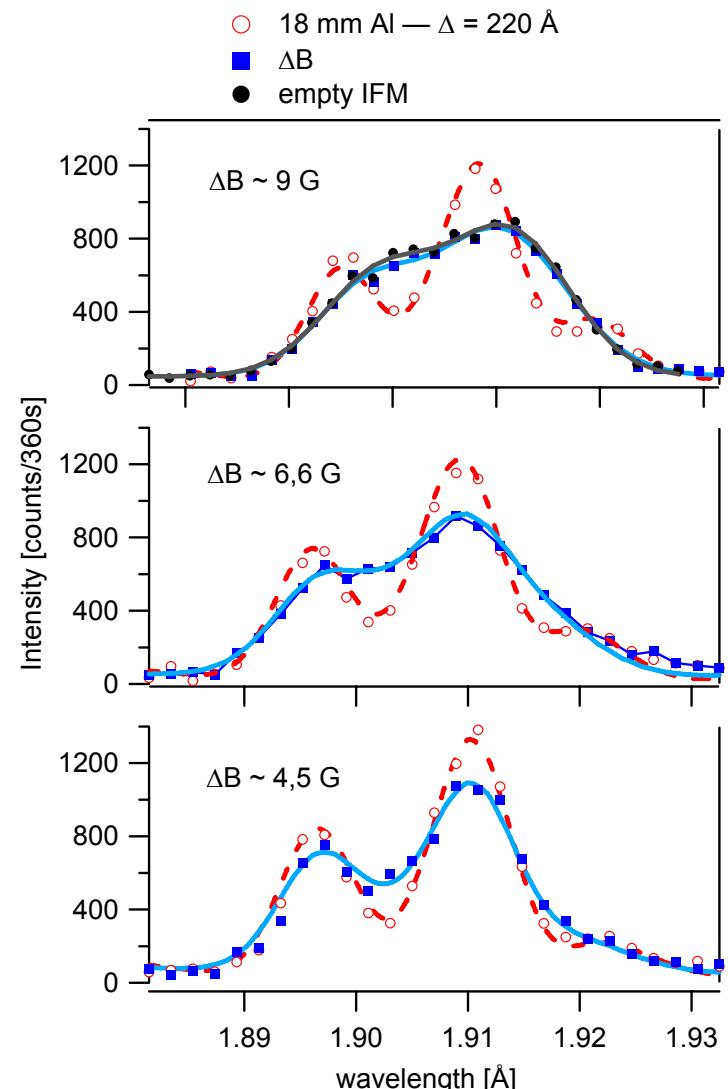
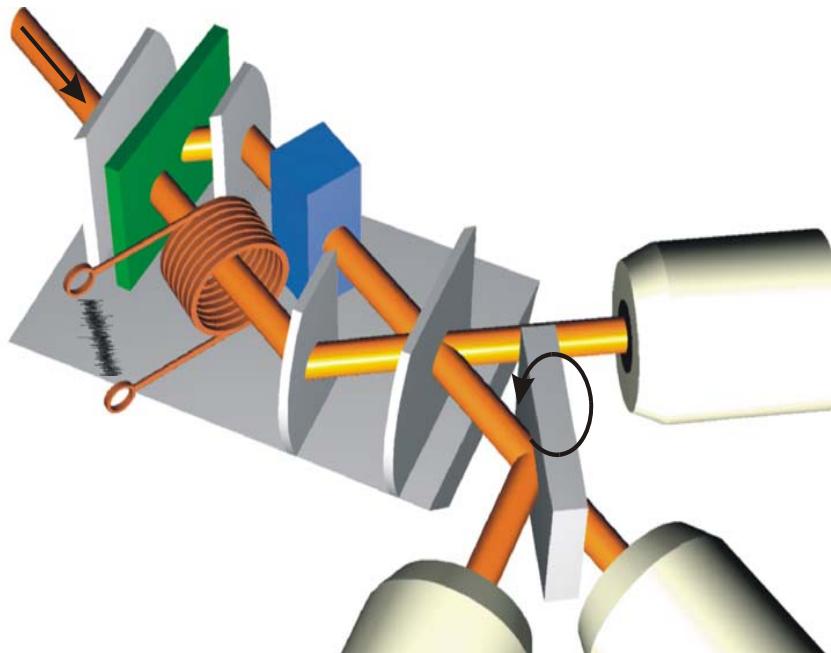


M.Baron, H.Rauch, M.Suda (in progress)

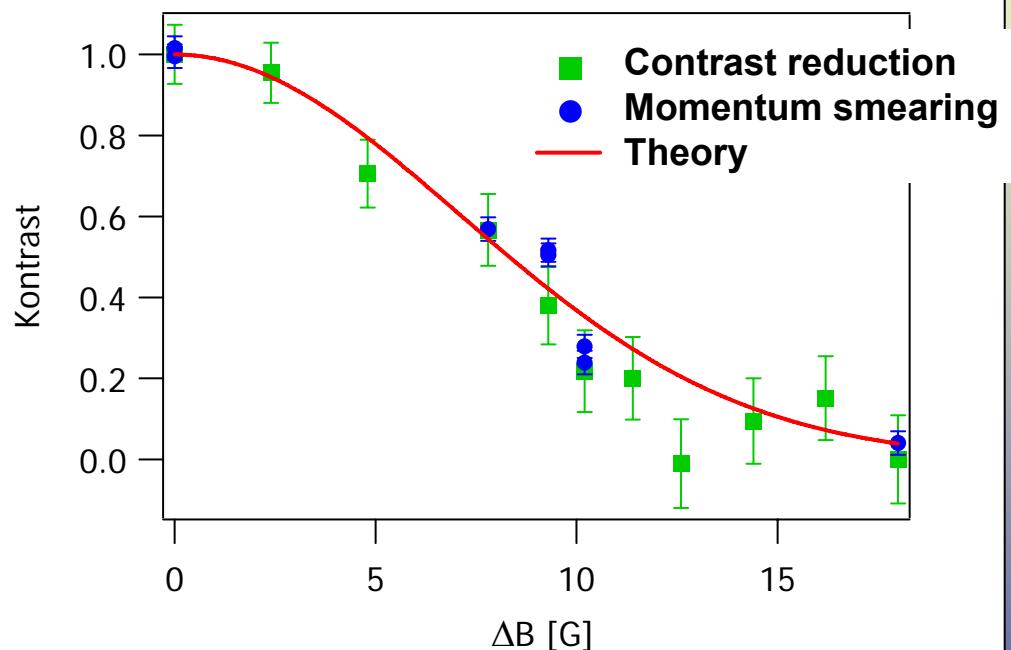
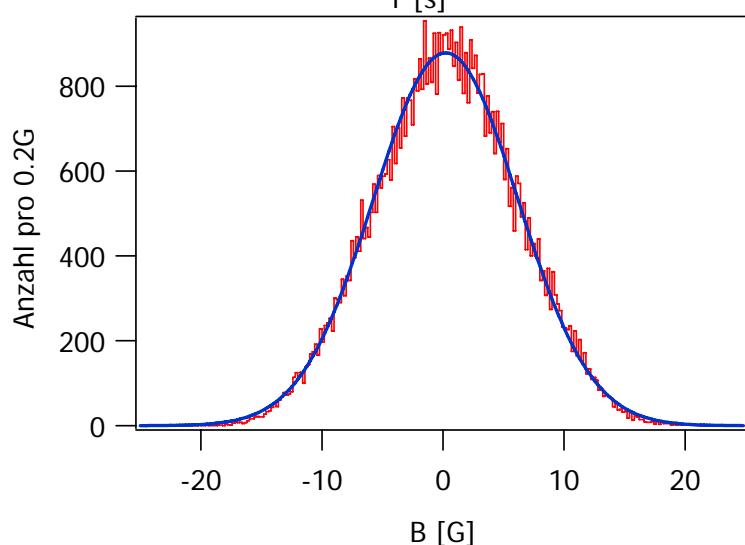
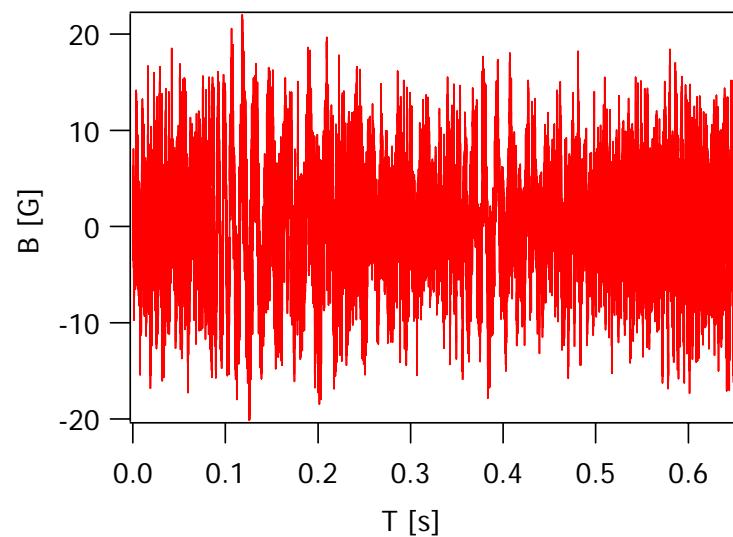
Magnetic noise fields



M.Baron, H.Rauch, M.Suda (in progress)

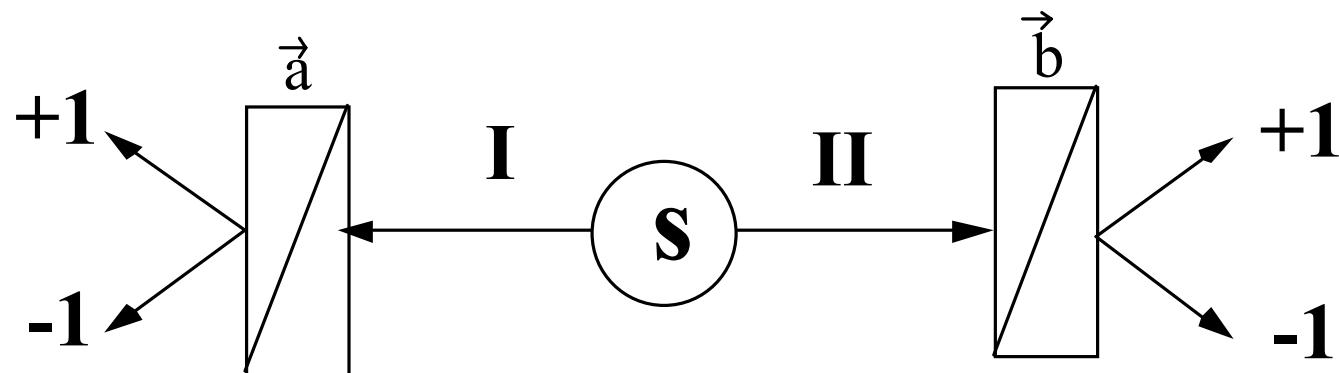


M.Baron, H.Rauch, M.Suda, J.Opt.B5 (2003) S244



$$C = C_0 \exp[-(\mu \Delta B D_{\text{eff}} / \hbar v)^2 / 2]$$

- *Quantum Contextuality*



Entanglement of two photon polarizations

$$|\Psi\rangle = \frac{1}{\sqrt{2}} \{ | \uparrow \rangle_I \otimes | \downarrow \rangle_{II} + | \downarrow \rangle_I \otimes | \uparrow \rangle_{II} \}$$

⇒⇒⇒ Entanglement between *Two-Particles*

A. Einstein, B. Podolsky and N. Rosen, Phys. Rev. 47 (1935) 777.

2-Particle Bell-State

$$|\Psi\rangle = \frac{1}{\sqrt{2}} \{ | \uparrow \rangle_I \otimes | \downarrow \rangle_{II} + | \downarrow \rangle_I \otimes | \uparrow \rangle_{II} \}$$

I, II represent **2-Particles**

Measurement on each particle

$$\begin{cases} \hat{A}^I(\xi) \\ \hat{B}^{II}(\xi) \end{cases} \quad -2 < S < 2$$

$$S = E(\alpha_1, x_1) - E(\alpha_1, x_2) + E(\alpha_2, x_1) + E(\alpha_2, x_2)$$

$$\text{where } \hat{P}_{(\xi; \pm)} = \frac{1}{2} (| \uparrow \rangle \pm e^{i\xi} | \downarrow \rangle) (\langle \uparrow | \pm e^{-i\xi} \langle \downarrow |)$$

$$\text{Then, } [\hat{A}^I, \hat{B}^{II}] = 0$$

\implies **(Non-)Contextuality**

(In)Dependent Results for commuting Observables

2-Space Bell-State

$$|\Psi\rangle = \frac{1}{\sqrt{2}} \{ | \uparrow \rangle_s \otimes | I \rangle_p + | \downarrow \rangle_s \otimes | II \rangle_p \}$$

s, p represent **2-Spaces**, e.g., spin

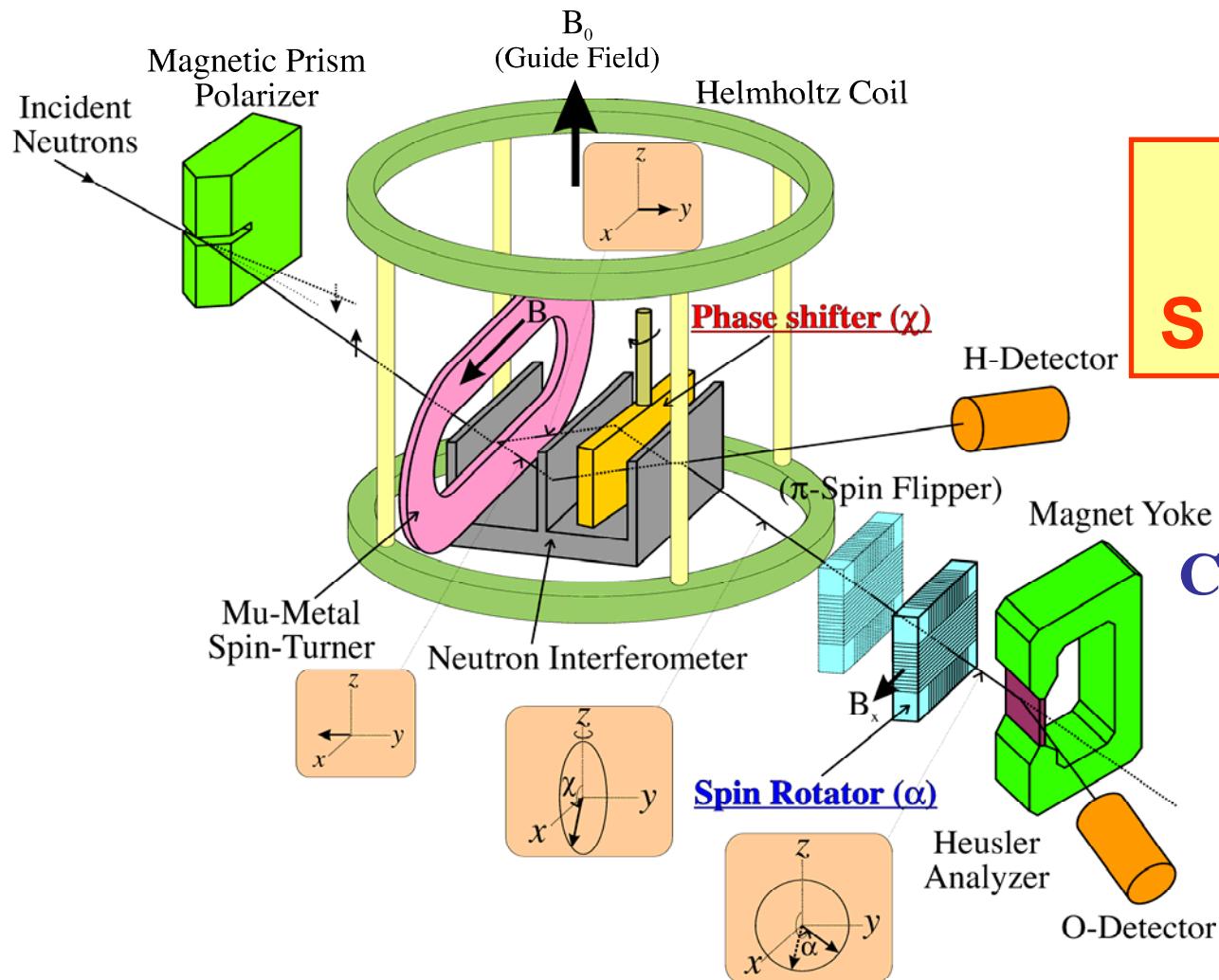
Measurement on each property

$$\begin{cases} \hat{A}^s(\phi) \\ \hat{B}^p(\phi) \end{cases} \quad -\pi < S < \pi$$

$$\text{where } \hat{P}_{(\phi)} = \frac{1}{2} (| \phi \rangle + e^{i\phi} | \bar{\phi} \rangle) (\langle \phi | + e^{-i\phi} \langle \bar{\phi} |)$$

$$\text{Then, } [\hat{A}^s, \hat{B}^p] = 0$$

Contextuality Experiment



Result:

$$S = 2.051 \pm 0.019$$

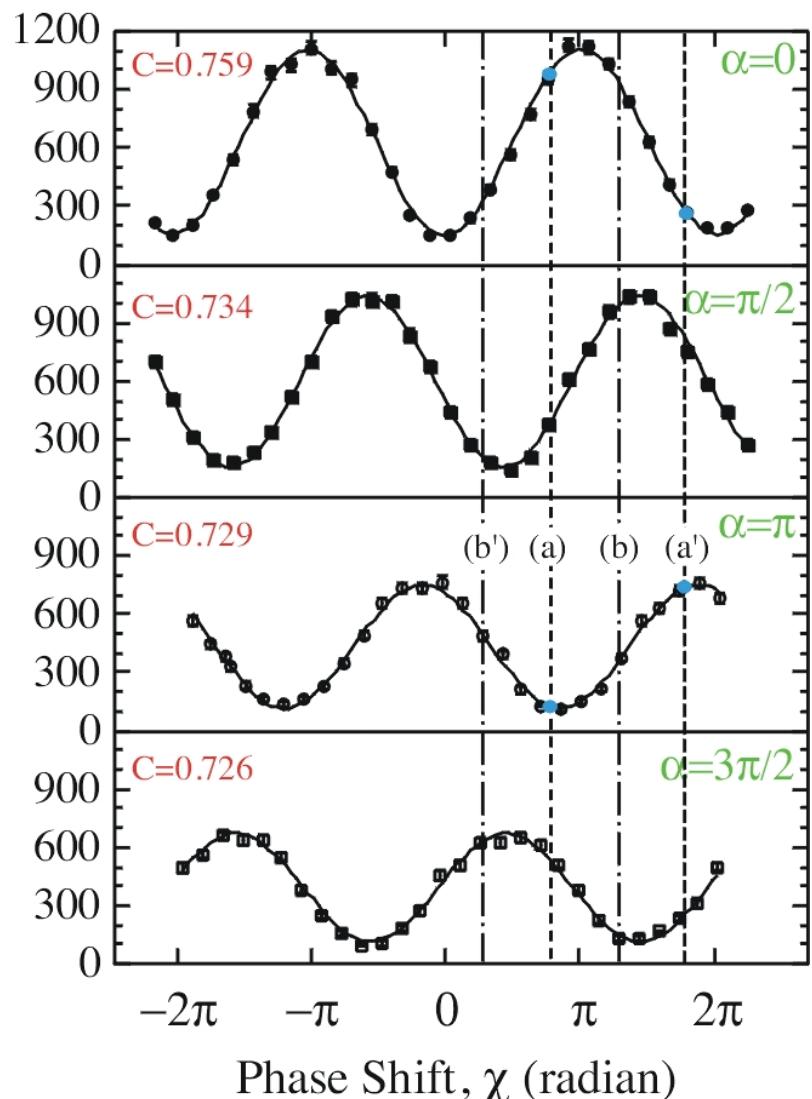
Theory: $S_{\text{Max}} = 2.82$

Classical correlation:
(hidden variables)

$$-2 < S' < 2$$

.Y.Hasegawa, R.Loidl, G.Badurek, M.Baron, H.Rauch, Nature 425 (2003) 45 and
Phys.Rev.Lett.97 (2006) 230401

Intensity (neutrons/120s)



$$\begin{aligned}
 & E'(\alpha=0, \chi = 0.79\pi) \\
 &= [N'(0, 0.79\pi) + N'(\pi, 1.79\pi) \\
 &\quad - N'(0, 1.79\pi) - N'(\pi, 0.79\pi)] \\
 &\div [N'(0, 0.79\pi) + N'(\pi, 1.79\pi) \\
 &\quad + N'(0, 1.79\pi) + N'(\pi, 0.79\pi)] \\
 &= 0.542
 \end{aligned}$$

In the same manner,

$$\left\{
 \begin{array}{l}
 E'(\alpha=0, \chi = 1.29\pi) \\
 E'(\alpha=0.5\pi, \chi = 0.79\pi) \\
 E'(\alpha=0.5\pi, \chi = 1.29\pi)
 \end{array}
 \right.$$

were determined.

Kochen-Specker phenomenon

$$C_{\text{non-contextual}} = C_{\text{classic}} = 2$$

$$C_{\text{contextual}} = C_{\text{quantum}} = 4$$

$$C_{\text{experimental}} = 3.138(15)$$

A cartoon-like

The colour of a skier is undetermined, represented by a question mark. After a ‘measurement’ on the path in our experiment, it takes its own colours (the direction), depending on what was measured. Basically no correlation is expected!

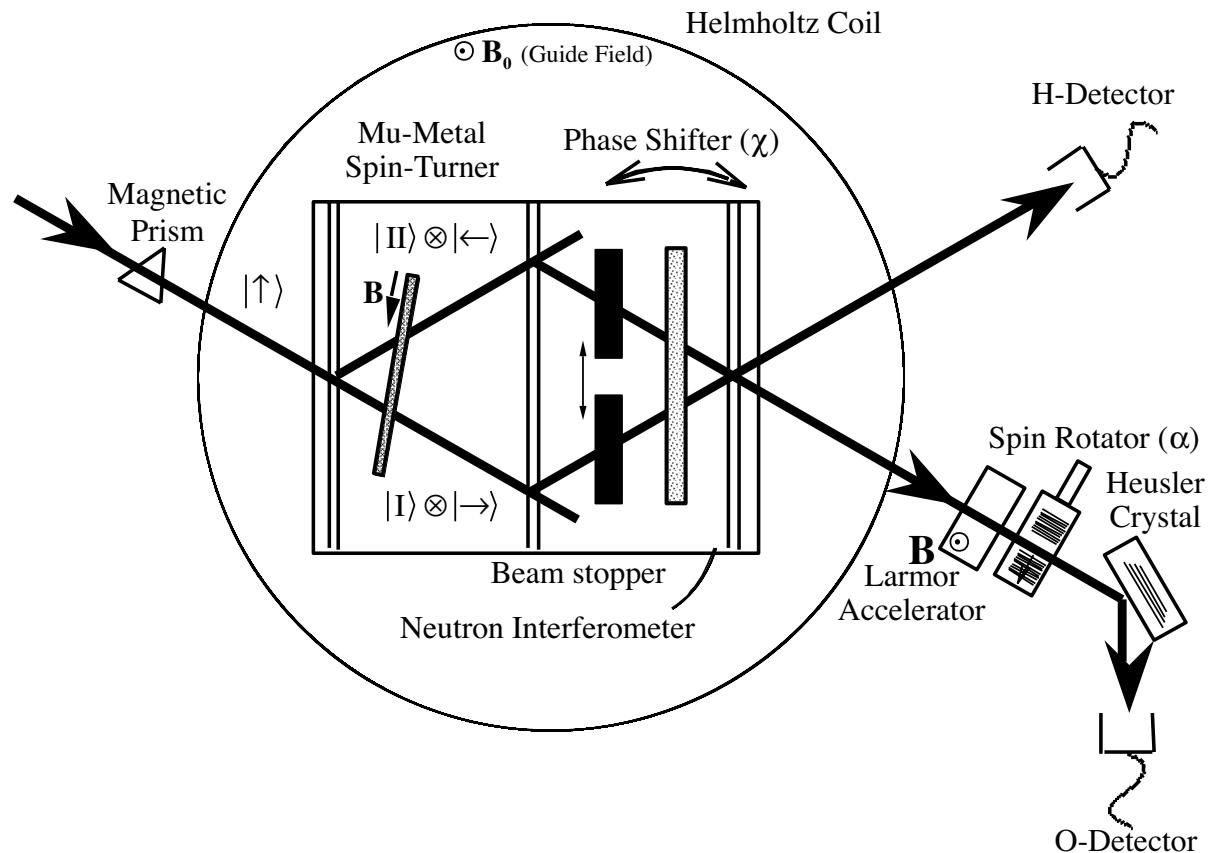
Y.Hasegawa, R.Loidl, G. Badurek, M.Baron, H.Rauch,
PRL 97 (2006) 23040

the colours of jacket and trousers

- *Spin State Reconstruction*

Experimental setup

Measurement principle: D.F.V.James, P.G.Kwiat, W.J.Munro, A.G.White, PR/A,64 (2001) o52312



Beam I
Beam II
Beam I+II, $\chi = 0, \pi/2$
 $|z\rangle, |-z\rangle, |x\rangle, |y\rangle$
 → **16 positions**

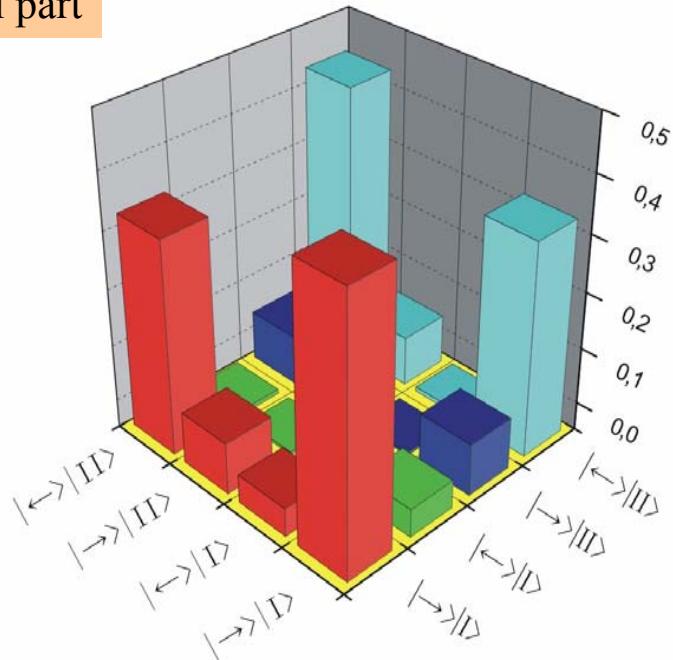
Y.Hasegawa, J.Klepp, S.Filipp, R.Loidl (in progress)

Quantum state tomography of neutron's Bell-state

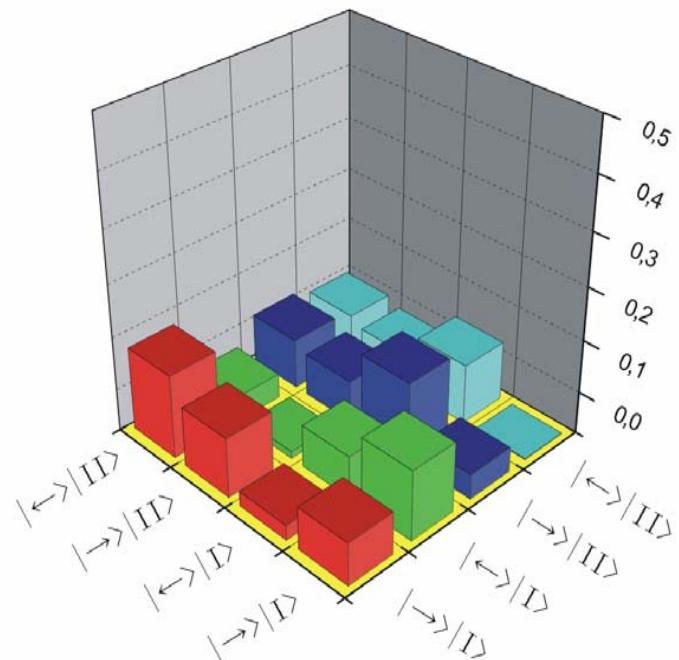
$$\rho = |\Psi\rangle\langle\Psi|$$

where $|\Psi\rangle = \{|\uparrow\rangle, |\downarrow\rangle\} \otimes \{|\text{I}\rangle, |\text{II}\rangle\}$

real part



imaginary part



Y.Hasegawa, J.Klepp, S.Filipp, R.Loidl (in progress)

- ***Geometrical Phases***

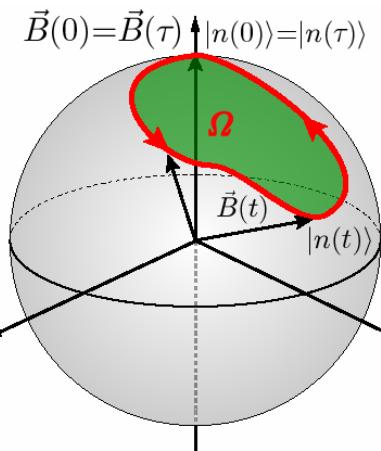
Berry Phase (adiabatic & cyclic evolution)

[Berry; Proc.R.S.Lond. A 392, 45 (1984)]

$$|\Psi(t)\rangle = e^{-i\phi_d} e^{i\phi_g} |n(R(t))\rangle$$

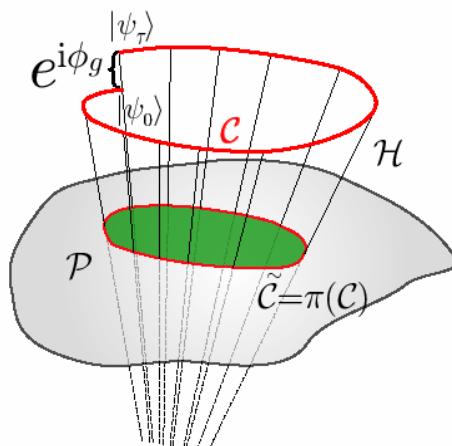
$$\phi_d(t) = \frac{1}{\hbar} \int_0^t dt' E_n(t')$$

$$\phi_g = -\frac{\Omega}{2} \quad (\text{for 2-level systems})$$



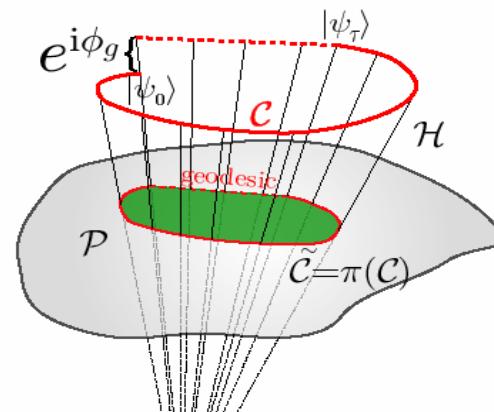
Non-adiabatic evolution

[Aharonov & Anandan, PRL 58, 1593 (1987)]



Non-adiabatic & non-cyclic evolution

[Samuel & Bhandari, PRL 60, 2339 (1988)]



Theory:

S.Pancharatnam, Proc.Ind.Acad.Sci. A44(1956)247

M.V.Berry, Proc.Roy.Soc.London, A392(1984)415

J.Anandan,Nature360(1992)307; R.Bhandari,Phys.Rep.281(1977)1

$$|\psi(T)\rangle = e^{i\phi} |\psi(0)\rangle$$



$$\phi = -\frac{1}{\hbar} \int_0^T \langle \psi(t) | H | \psi(t) \rangle dt + i \int_0^T \langle \psi(t) | \frac{d}{dt} | \psi(t) \rangle dt = \delta + \gamma$$

$$\delta = \frac{\alpha}{2} \cos \Theta \text{ dynamical phase}$$

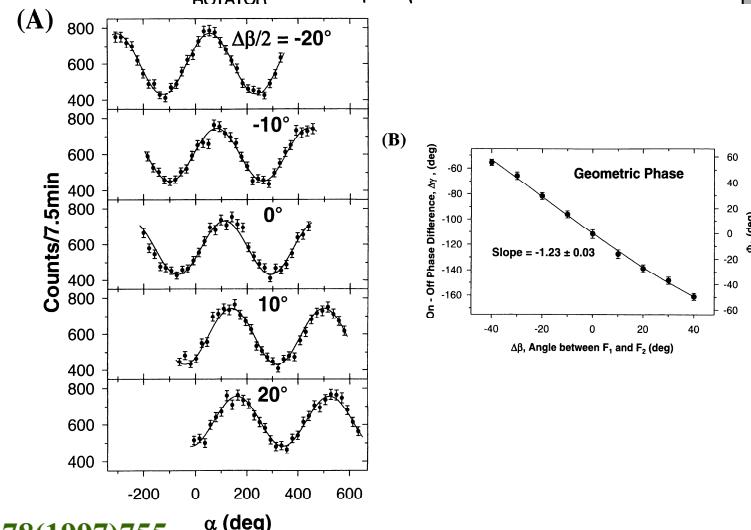
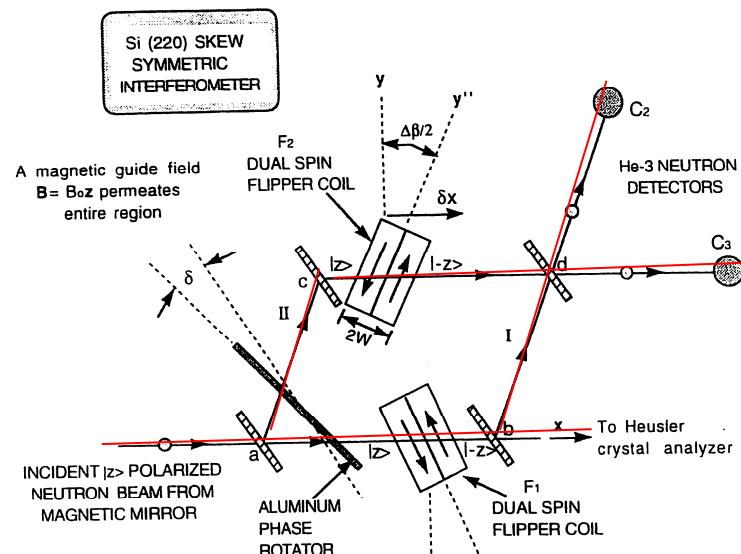
$$\gamma = \frac{\alpha}{2} (1 - \cos \Theta) \text{ ... geometric phase}$$

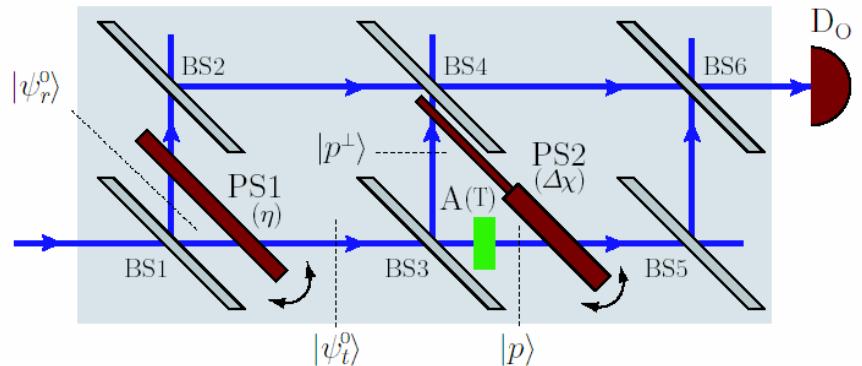
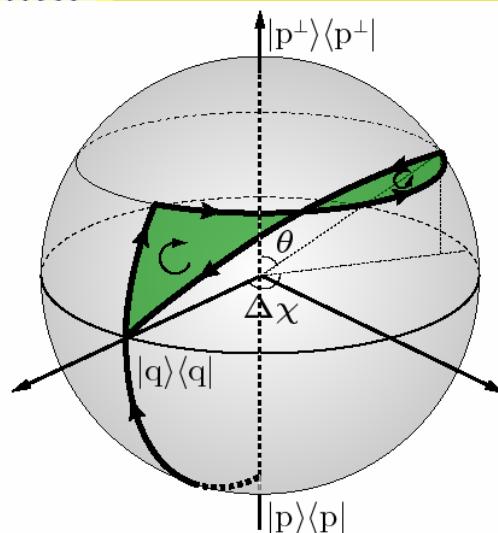


$$\begin{aligned} I(\chi, \alpha) &= |\psi_0(0,0) + \psi_0(\chi, \alpha)|^2 \propto D + \cos \chi \cos \frac{\alpha}{2} + \sin \chi \sin \frac{\alpha}{2} \cos \Theta \\ &= D + A \cos(\chi + \phi) \end{aligned}$$

$$\cos \phi = \frac{\cos \frac{\alpha}{2}}{\sqrt{\cos^2 \frac{\alpha}{2} + \cos^2 \frac{\alpha}{2} \cos^2 \Theta}}$$

Exp.: A.G.Wagh, V.C.Rakhecha, J.Summhammer, G.Badurek, H.Weinfurter,
B.E.Allman, H.Kaiser, K.Hamacher, D.L.Jacobson, S.A.Werner, Phys.Rev.Lett. 78(1997)755





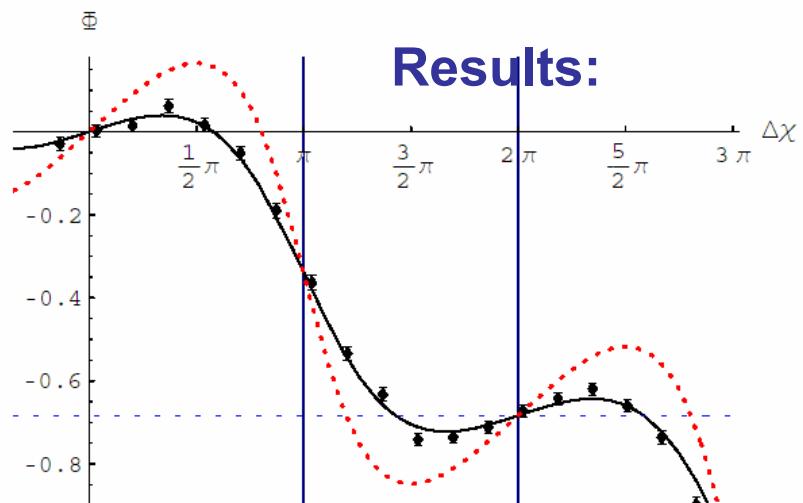
$$\Phi \equiv \arg\langle\psi'_r|\psi'_t\rangle = \frac{\chi_1 + \chi_2}{2} - \arctan \left[\tan \frac{\Delta\chi}{2} \left(\frac{1 - \sqrt{T}}{1 + \sqrt{T}} \right) \right]$$

$$\Phi_g \equiv \Phi - \Phi_d = \Phi$$

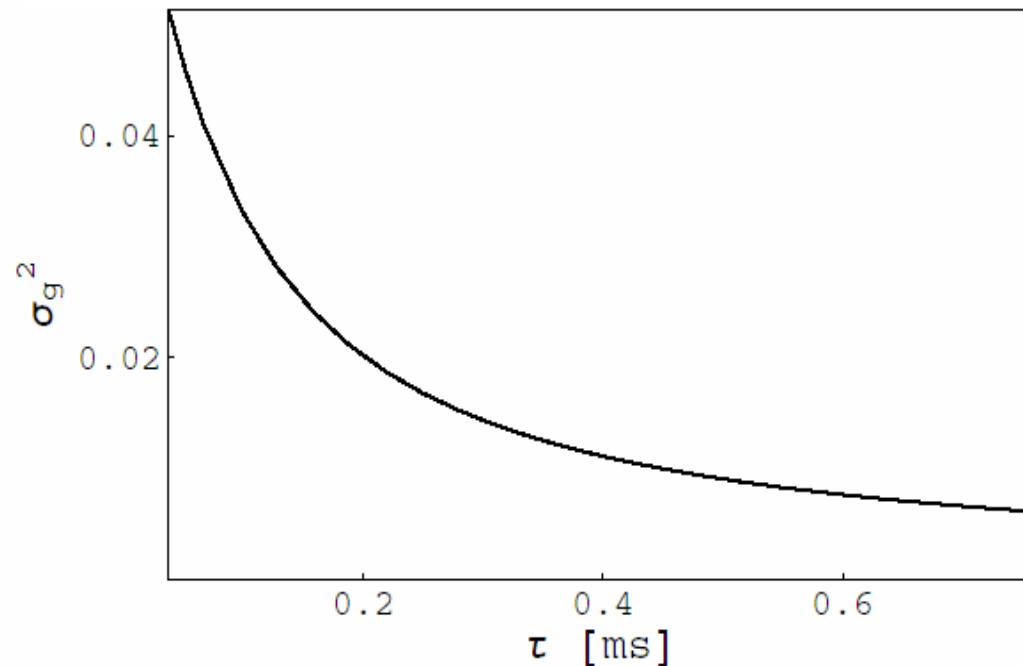
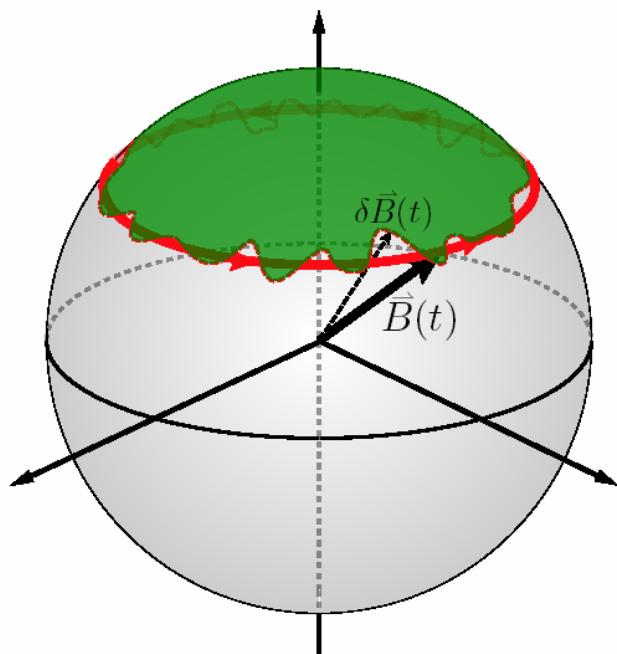
Cancelling dynamical phase, if

$$\Phi_d = \frac{\chi_1 + T\chi_2}{1 + T} = 0$$

S. Filipp, Y. Hasegawa, R. Loidl and H. Rauch, Phys.Rev. A72 (2005) 021602



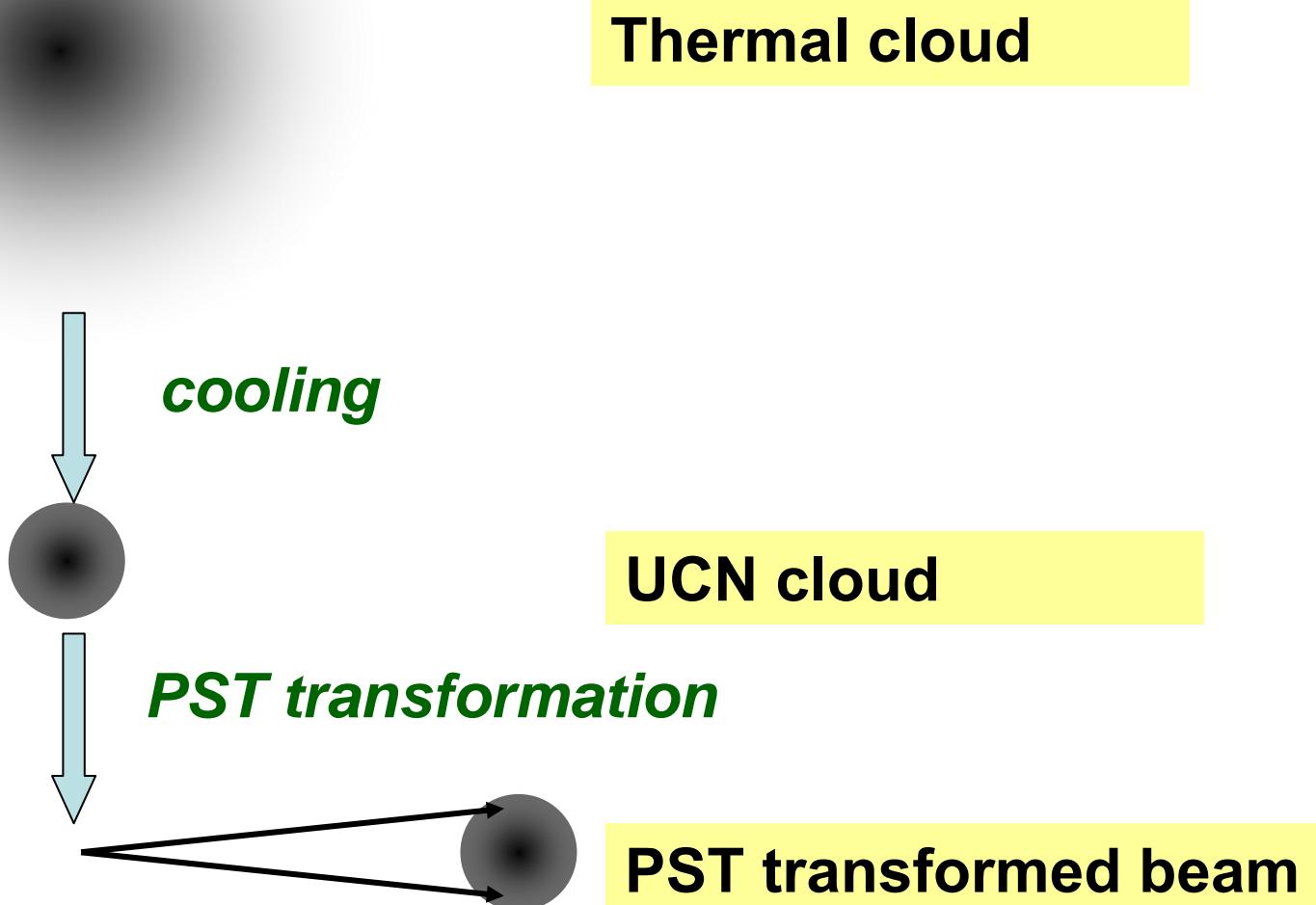
[De Chiara and Palma, PRL 91, 090404 (2003)]



Variance of geometric phase (σ_g^2) tends to 0 for increasing time of evolution in magnetic field.

- *Ultracold Neutrons and*
- *Phase Space*
- Transformation*

UCN – PST beam tailoring



$$dN = w(\vec{x}, \vec{k}) N d\vec{x} d\vec{k}$$



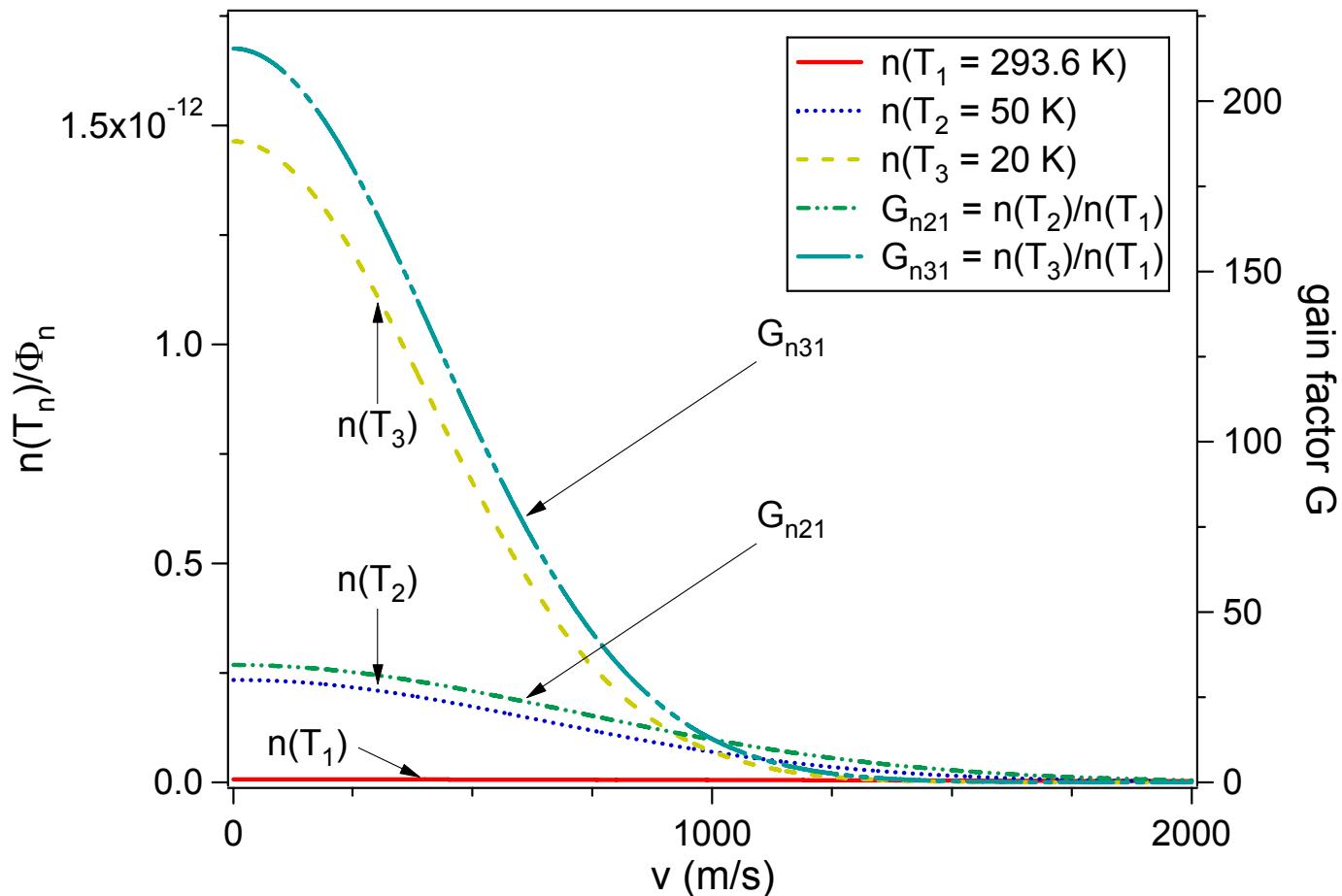
$$dN = \frac{1}{\pi^{3/2} k_T^3} N e^{-\frac{\hbar^2 k^2}{2m k_B T}} d\vec{x} d\vec{k}$$

$$\Phi = \frac{2N}{\sqrt{\pi}} v_T$$

Maier-Leibnitz Formel

$$dN = \frac{\Phi}{2\pi v_T k_T^3} e^{-\frac{\hbar^2 k^2}{2m k_B T}} d\vec{x} d\vec{k}$$

PHASE SPASE DENSITY AND GAIN FACTORS



$$n(T) = \frac{\rho(v, T)}{4\pi} = \frac{\Phi_n}{2\pi v_T^4} \exp\left(-\frac{v^2}{v_T^2}\right)$$

$$G(T_1, T_2) = \frac{n(T_2)}{n(T_1)} = \frac{v_{T_1}^4}{v_{T_2}^4} \exp\left[-v^2 \left(\frac{1}{v_{T_2}^4} - \frac{1}{v_{T_1}^4}\right)\right]$$

LUMINOSITY

$$L_z(T) = n(T) \frac{dV_P}{dt dA d\Omega} = \frac{\Phi_n}{2\pi} \frac{v_z v^2}{v_T^4} \exp\left(-\frac{v^2}{v_T^2}\right) dv_z$$

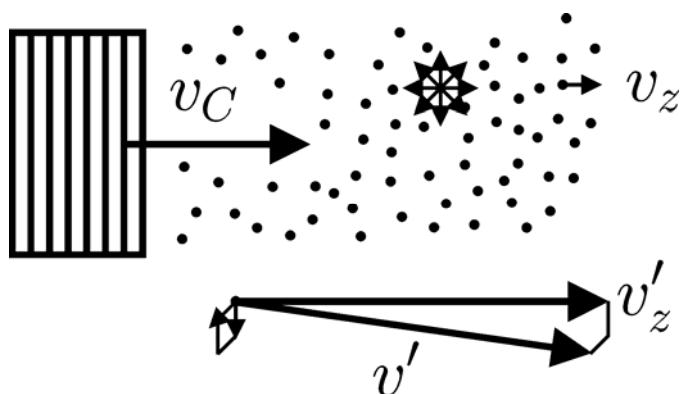
$$dV_P = dx dy dz dv_x dv_y dv_z$$

$$dt = \frac{dz}{v_z}$$

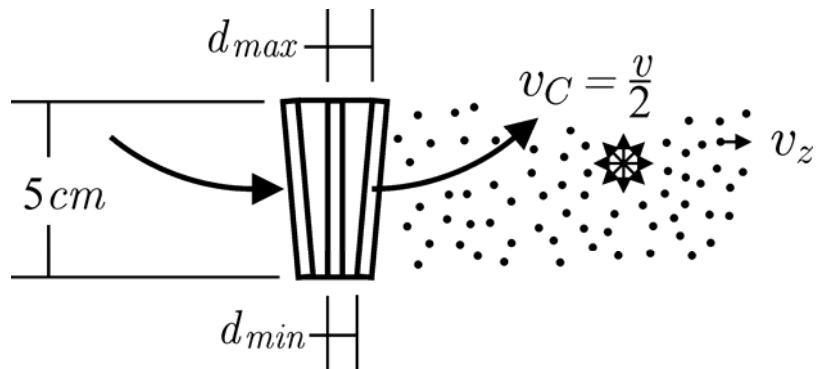
$$dA = dx dy$$

$$d\Omega = dv_x dv_y / v^2$$

linear



rotating

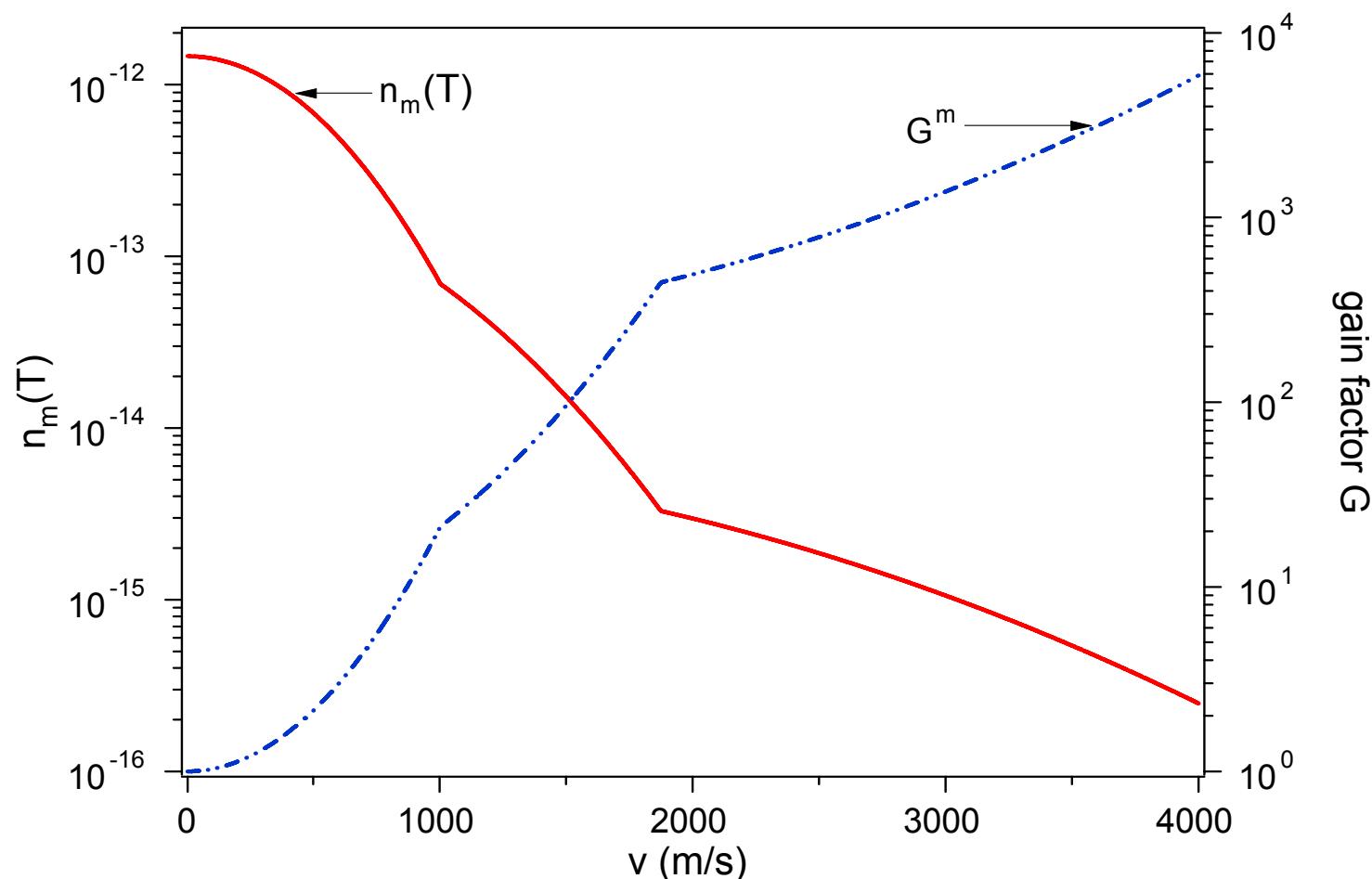


$$v_c = 2\pi r v = h/m d_{hkl}$$

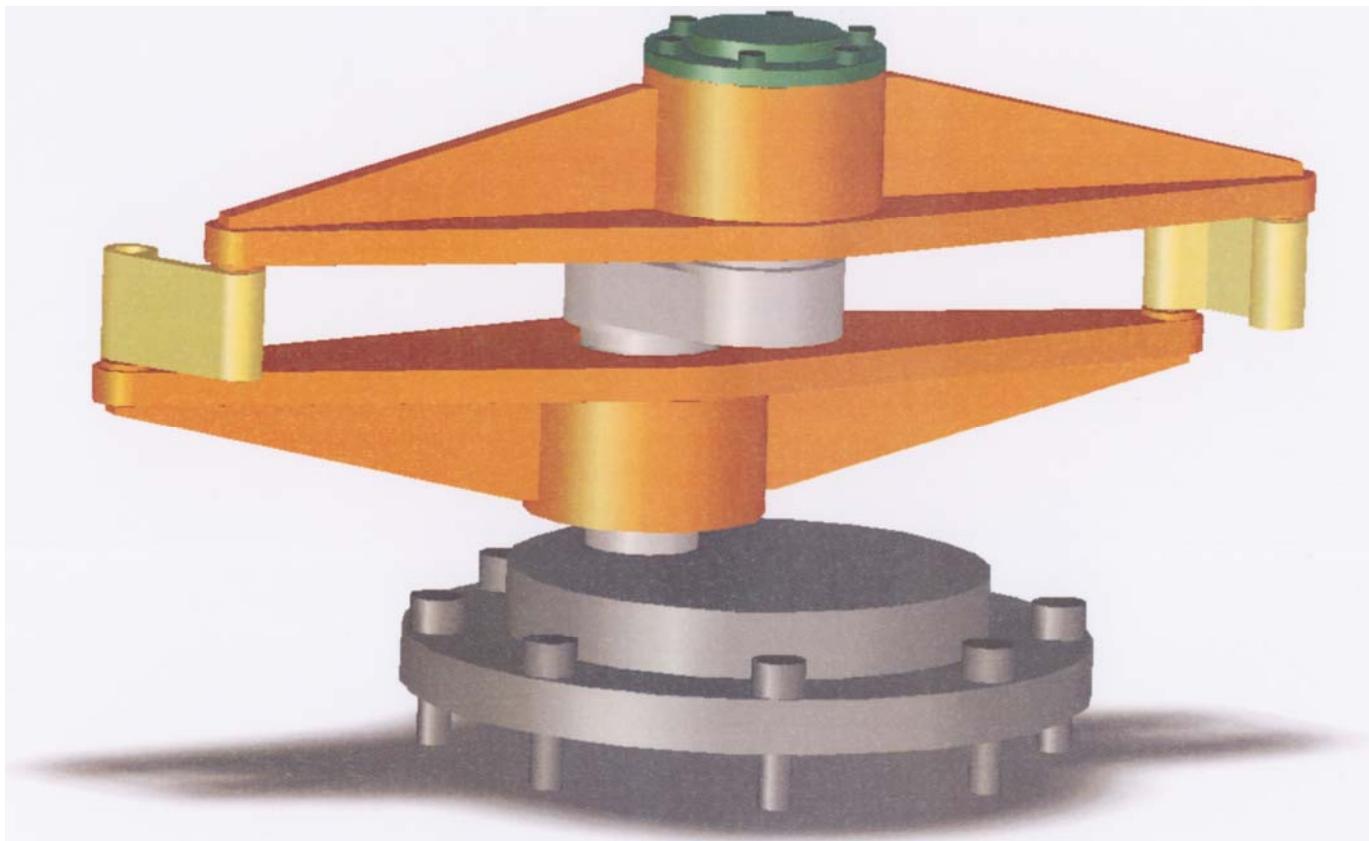
Example: $r = 35 \text{ cm}$, $v = 13,648 \text{ rpm}$

$$\rightarrow v_c = 500 \text{ m/s} \quad (v = 1000 \text{ m/s})$$

$$d_0 = 4.0 \text{ \AA}, \quad \Delta d/d_0 = 6\%$$



$$G_S^m = \frac{n_S(T_3)}{\max\{n(T_1), n(T_2), n(T_3)\}}$$



© Fa.Merk, Darmstadt



Thank You

Remarks on Neutron Interferometry

Quantum State Preparation and Measurements

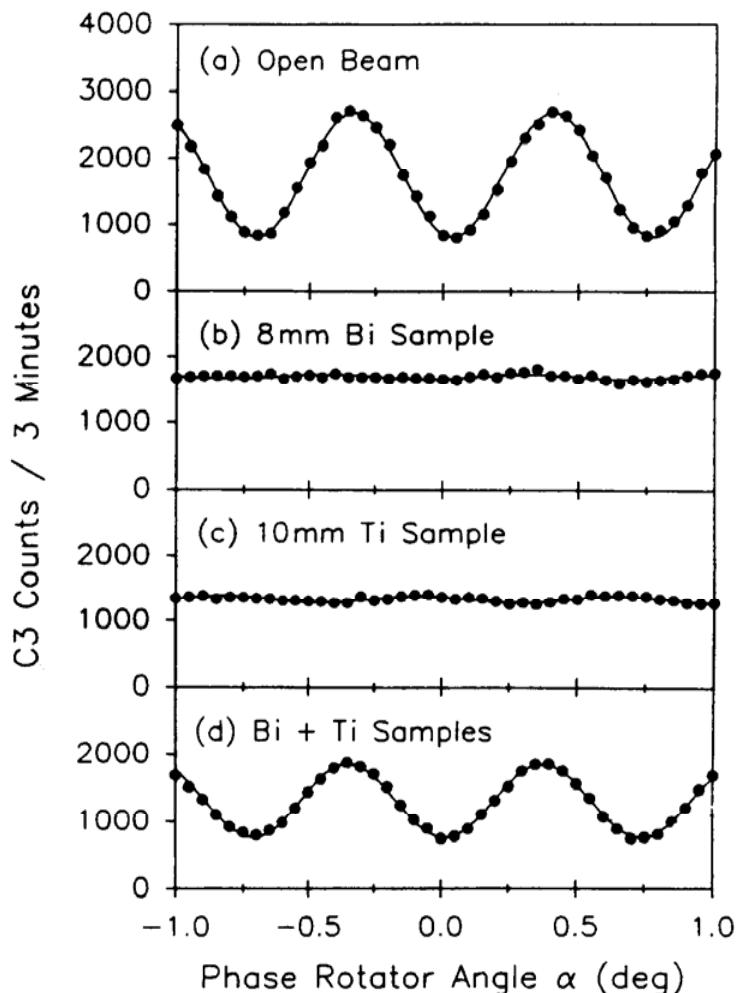
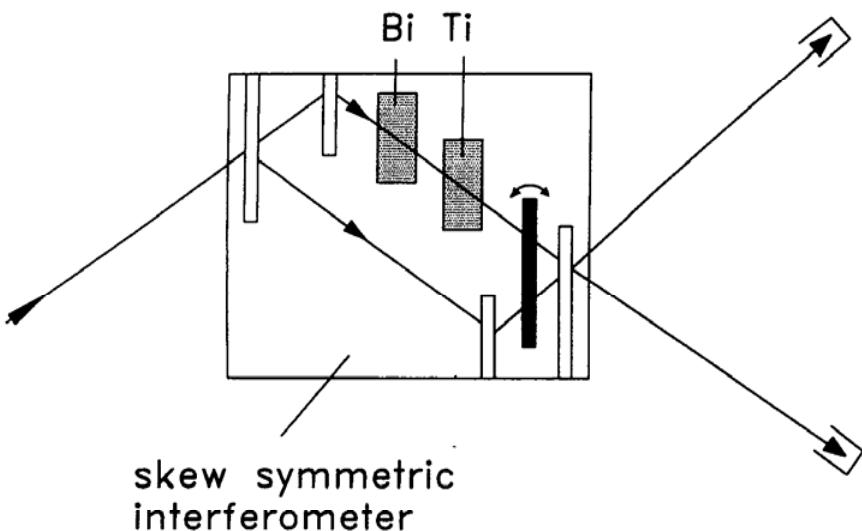
Magnetic Noise Dephasing

Confinement induced phase

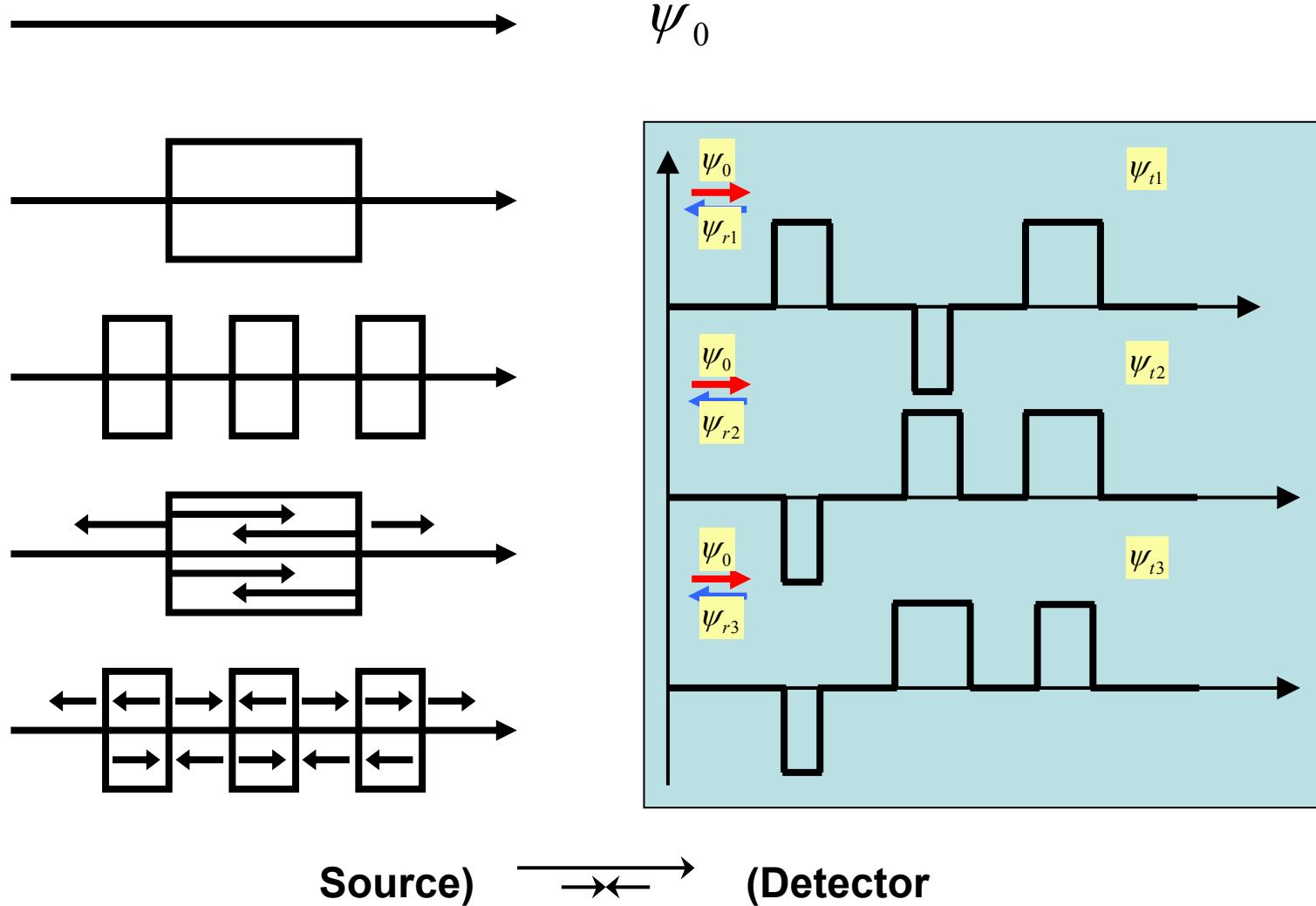
Contextuality

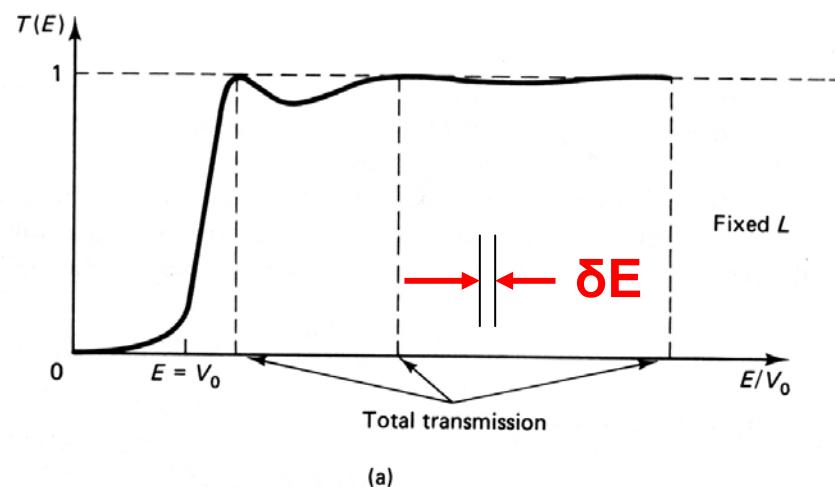
Quantum State Tomography

Unavoidable Losses



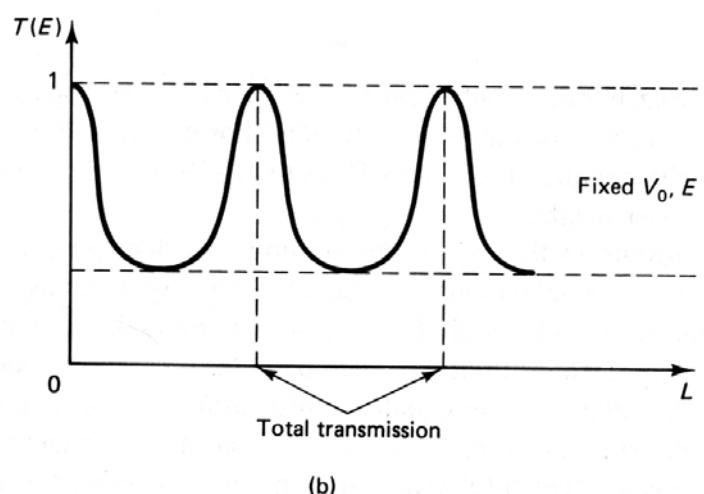
Clothier R., Kaiser H., Werner S.A., Rauch H., Wölwitsch H.,
Phys.Rev.A44 (1991)5357





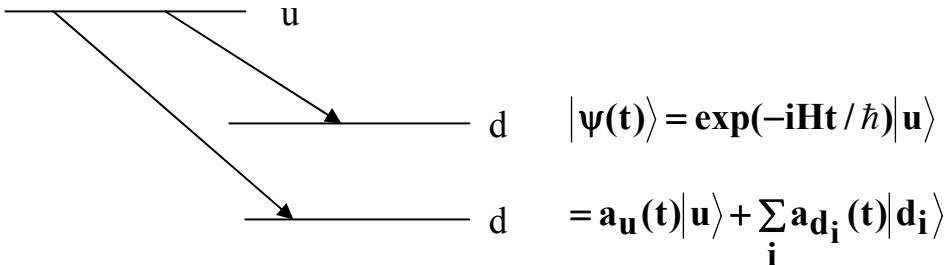
$$T + R = 1$$

→ $T < 1$



$$R_{\text{Min}} = (V/2E)^2 \delta k^2 L^2 > 0$$

B. Misra and E.C.G. Sudarshan 1977



Probability of finding the system undecayed:

$$P(t) = |a_u(t)|^2 = \langle u | \exp(-Ht/\hbar) | u \rangle^2$$

$$\approx 1 - (\Delta H / \hbar)^2 t^2 + O(t^4) \dots$$

$$(\Delta H)^2 = \langle u | H^2 | u \rangle - \langle u | H | u \rangle^2$$

Measurements: N-times in [0,t]

$$P_N(t) = \left[1 - (\Delta H / \hbar)^2 \left(\frac{t}{N} \right)^2 \right]^N$$

$$N \rightarrow \infty$$

$$= 1 - (\Delta H / \hbar)^2 \left(\frac{t^2}{N} \right) + \dots \rightarrow 1$$

a) Spin rotation

$$P_+ = \cos^2 \left(\frac{\omega_L \ell_0}{2v} \right) \rightarrow 0$$

$$\ell_0 = (2m+1)\pi v / \omega_L$$

$$(\omega_L = 2|\mu|B / \hbar)$$

b) Zeno situation (ideal)

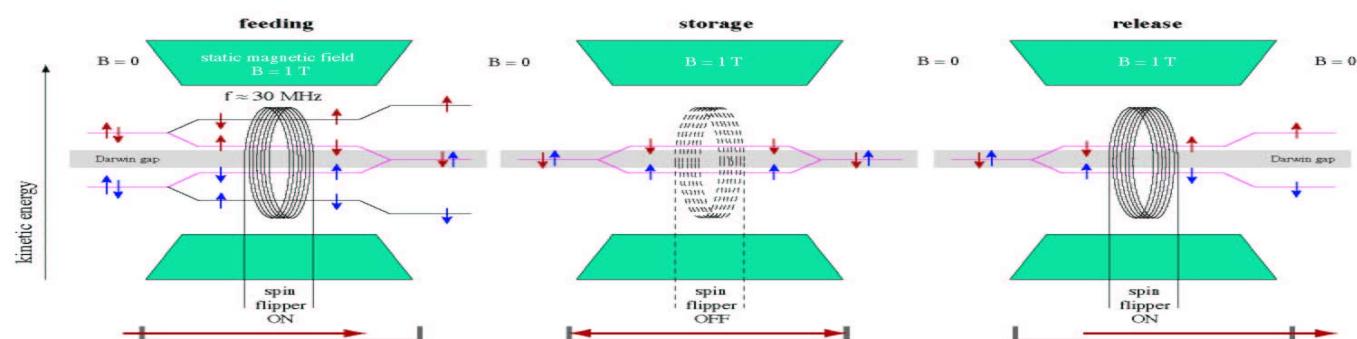
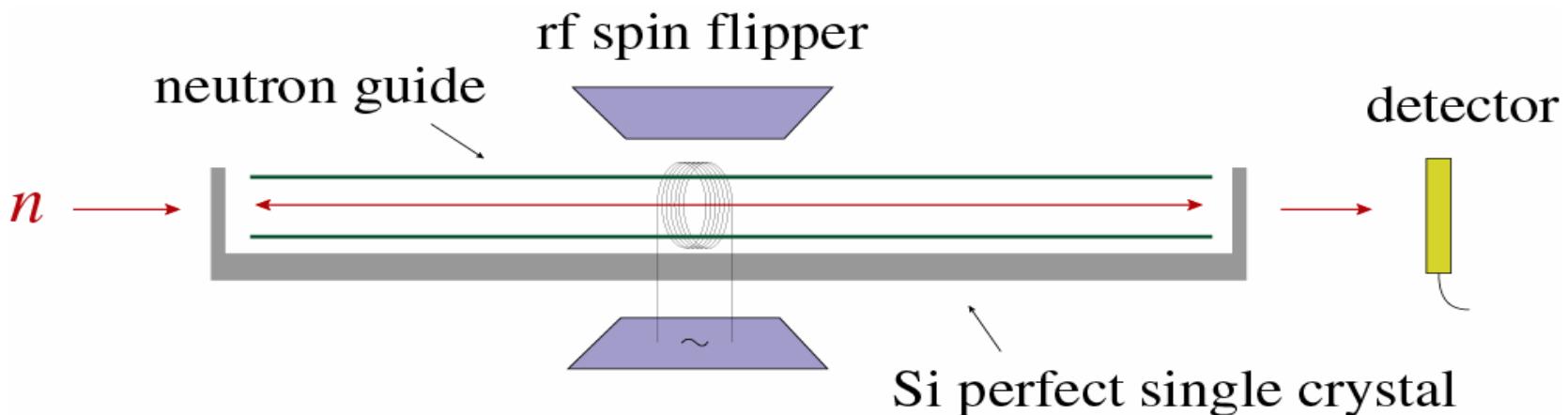
$$P_+ = \left[\cos^2 \left(\frac{\omega_L \ell}{2v} \right) \right]^n = \left[\cos^2 \frac{\pi}{2n} \right]^n \xrightarrow{n \rightarrow \infty} 1$$

$$\ell = \ell_0 / n$$

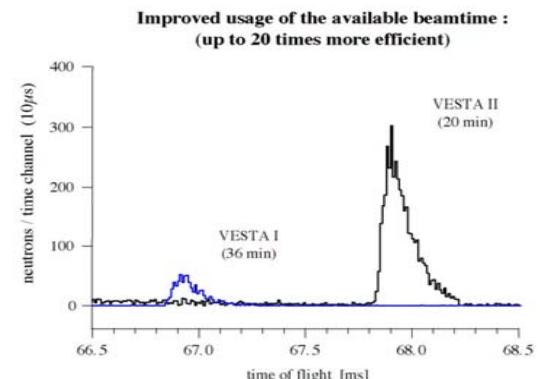
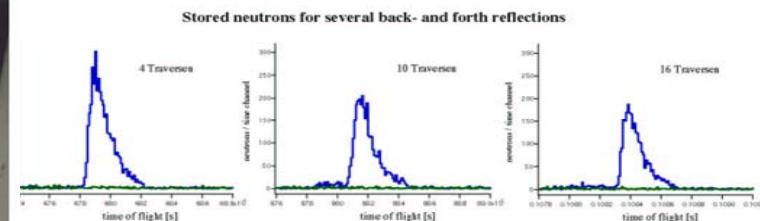
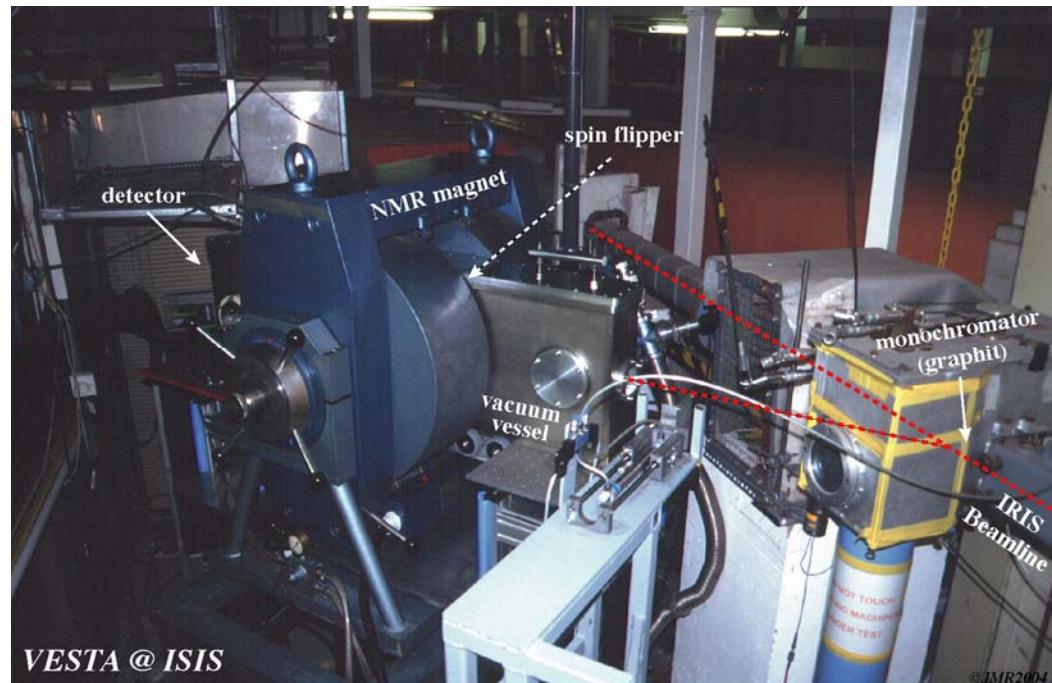
c) Zeno situation (real)

$$I_+ = P_+ \bar{T}^n = P_+ (1 - \bar{R})^n = P_+ \left[1 - \frac{1}{2} \left(\frac{V_0}{2E} \right) \right]^n$$

$$\cong 1 - \frac{n}{2} \left(\frac{V_0}{2E} \right)^2 + \dots \xrightarrow{n \rightarrow \infty} 0$$



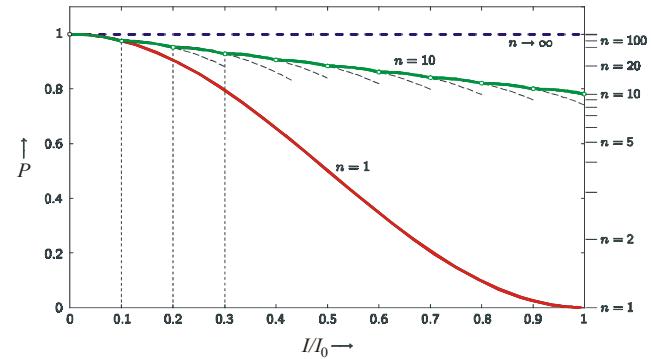
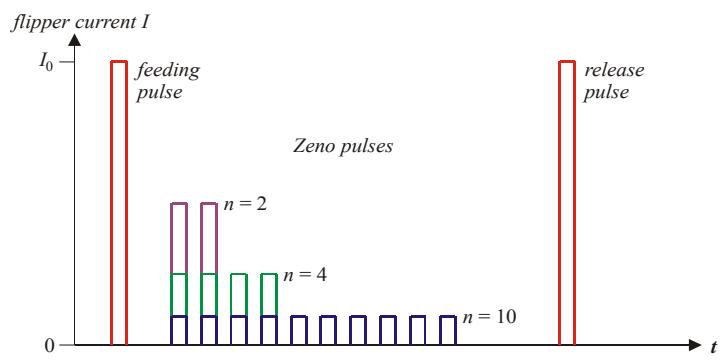
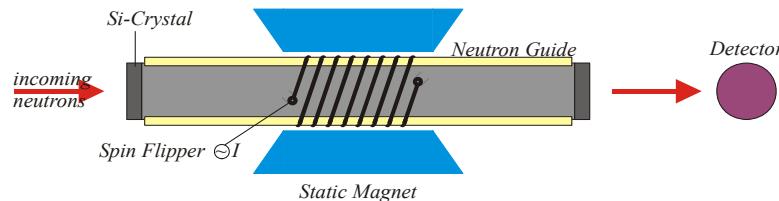
Set-up at ISIS Spallation Source



M.R.Jaekel, E.Jericha, H.Rauch, Nucl.Instr.Meth. A539 (2005) 335

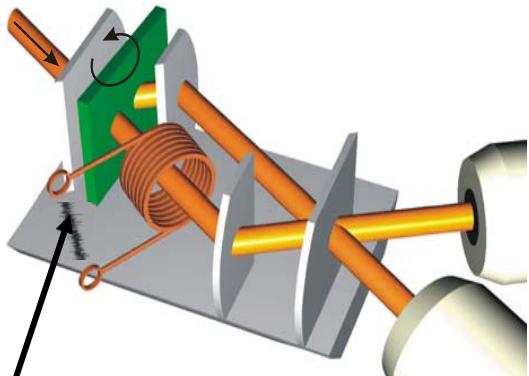
Quantum Zeno Effect

$$P_{Survival} = \cos^2\left(\frac{\mu Bl}{\hbar v}\right) = \cos^2\left(\frac{\pi I}{2I_0}\right) \xrightarrow{I \rightarrow I_0} 0$$

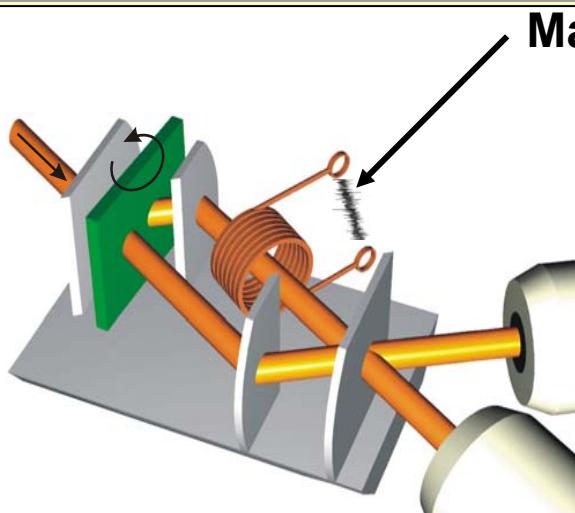


E.Jericha, M.Jaekel, S.Pascazio, H.Rauch (in progress)

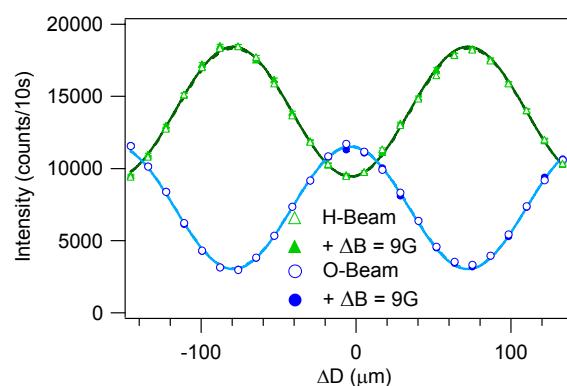
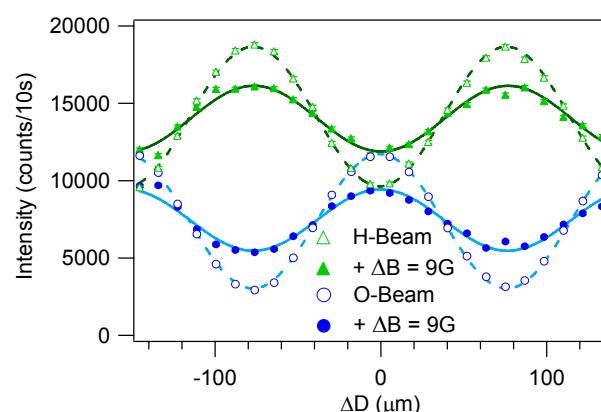
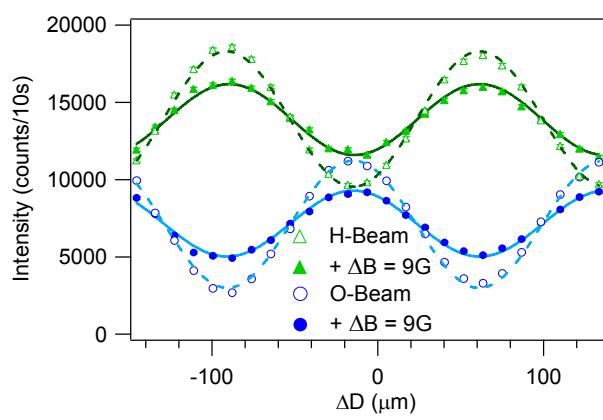
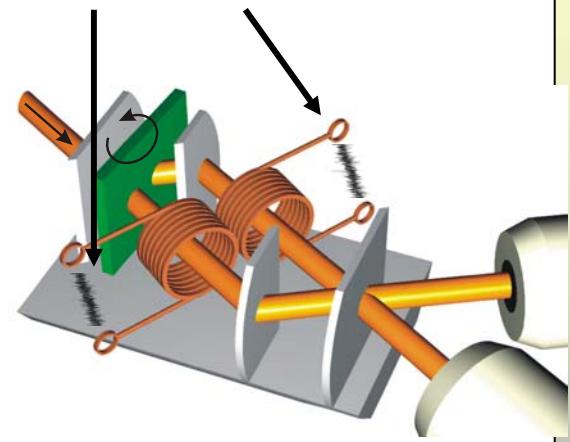
- *The neutron is an ideal tool for quantum experiments*
- *Quantum mechanics has been verified illustrating some of its strange features*
- *There is no natural limit between quantum and classical world*
- *Non-locality and contextuality are fundamental laws*
- *There is much more information in a quantum system than usually extracted*
- *Unavoidable quantum losses may play an important role in the understanding of decoherence phenomena*



Magnetic noise field



Magnetic noise fields



M.Baron, H.Rauch, M.Suda (in progress)

Helmut Rauch

Atominstut der Österreichischen Universitäten, Wien

Remarks on Neutron Interferometry
Quantum State Preparation and Measurement
Magnetic Noise Dephasing
Confinement induced phase
Contextuality
Quantum State Tomography and Topological Phases
Unavoidable Losses

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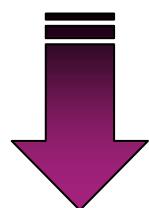
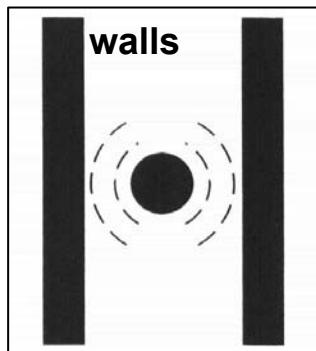
Confinement induced phase

Contextuality

Quantum State Tomography

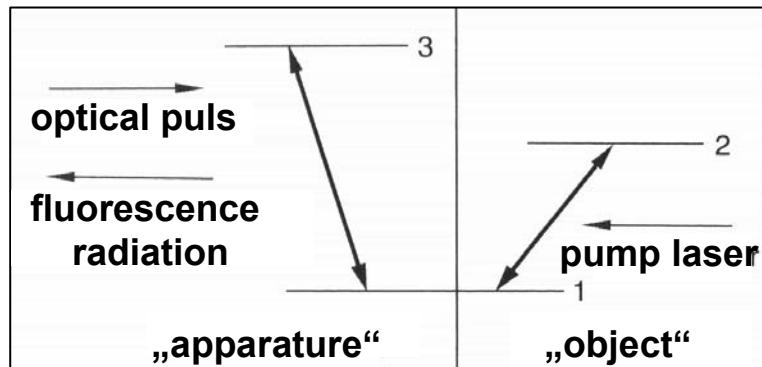
Unavoidable Losses

Spatial Confinement (Casimir-Effect)



Decay depends on
Distance

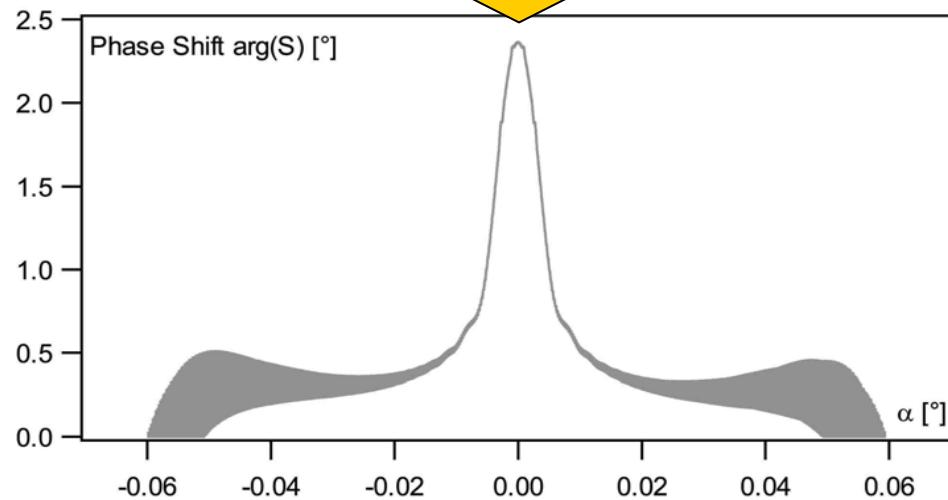
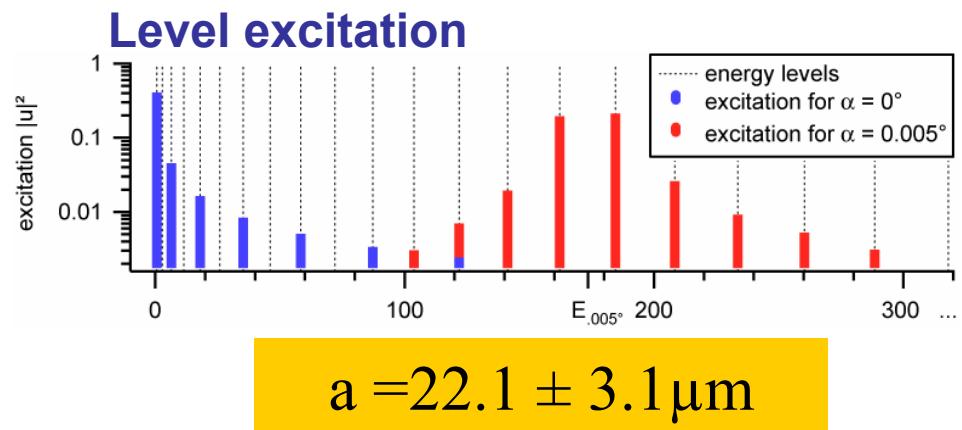
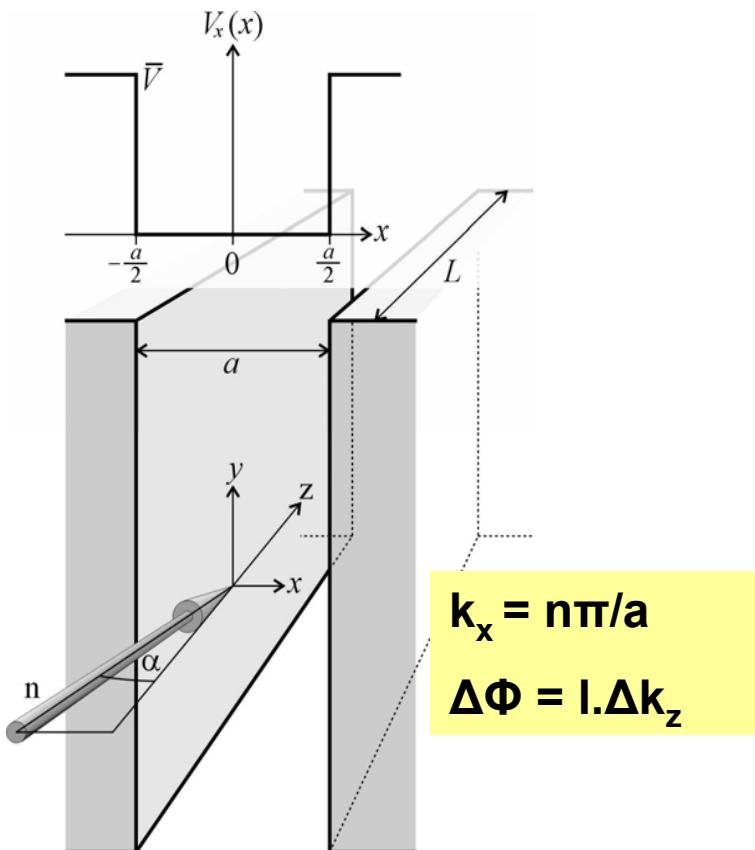
Temporal Confinement (Zeno-Effect)

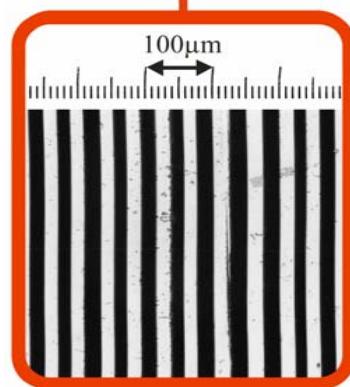
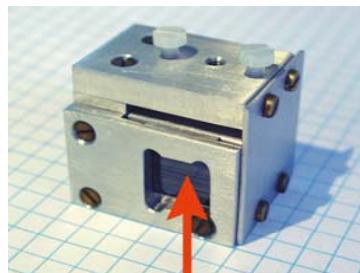


Decay depends on
Observation Frequency

Proposal by: J.M.Levy-Leblond 1987; D.M.Greenberger 1988

- 187 levels
- $E_0 = 0.4872 \text{ peV}$

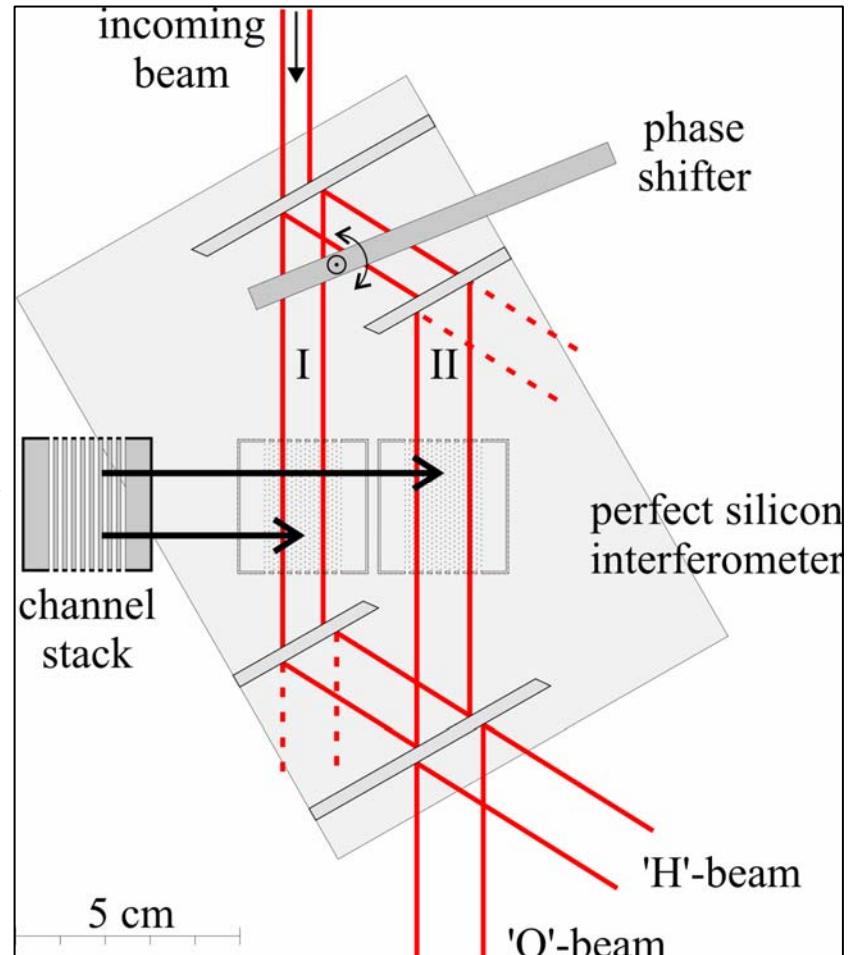


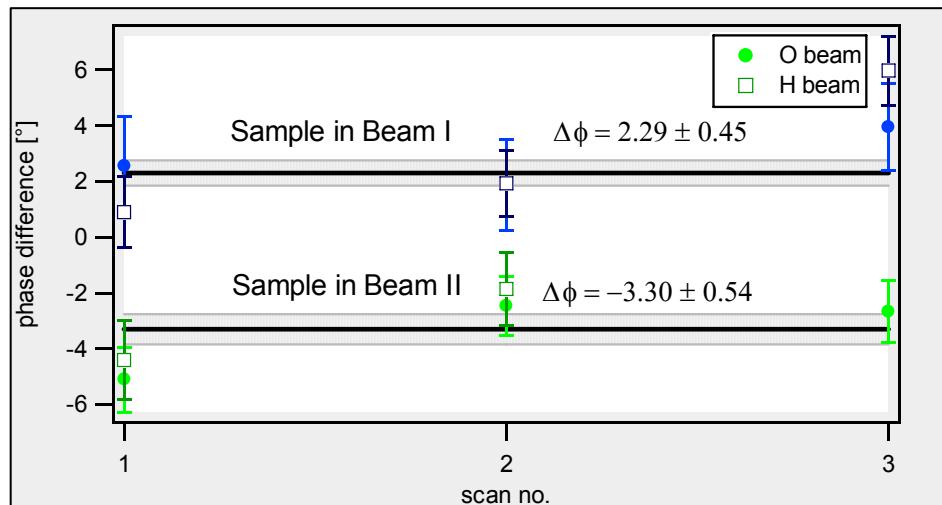
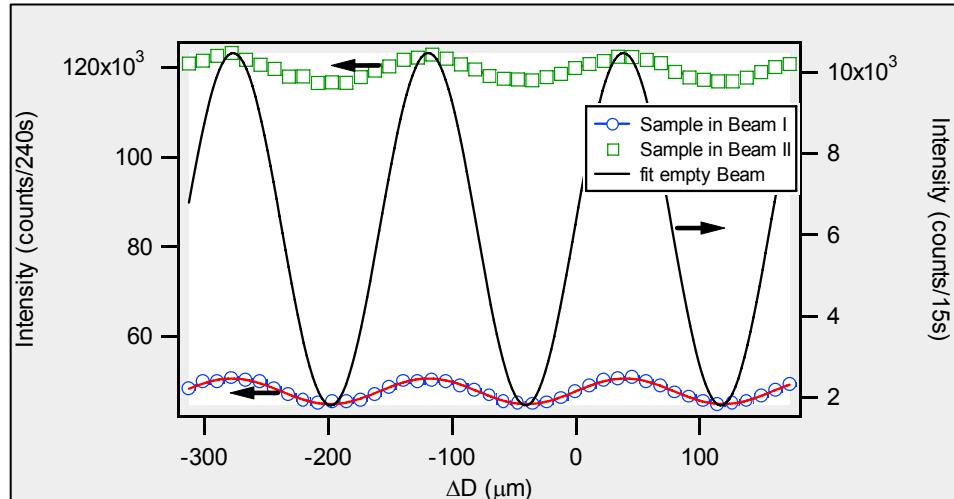


Calc. $\Delta\Phi = 2.50^\circ$
Exp. $\Delta\Phi = 2.8(4)^\circ$

H.Rauch, H.Lemmel, M.Baron, R.Loidl,
 Nature 417 (2002) 630

Experimental Setup





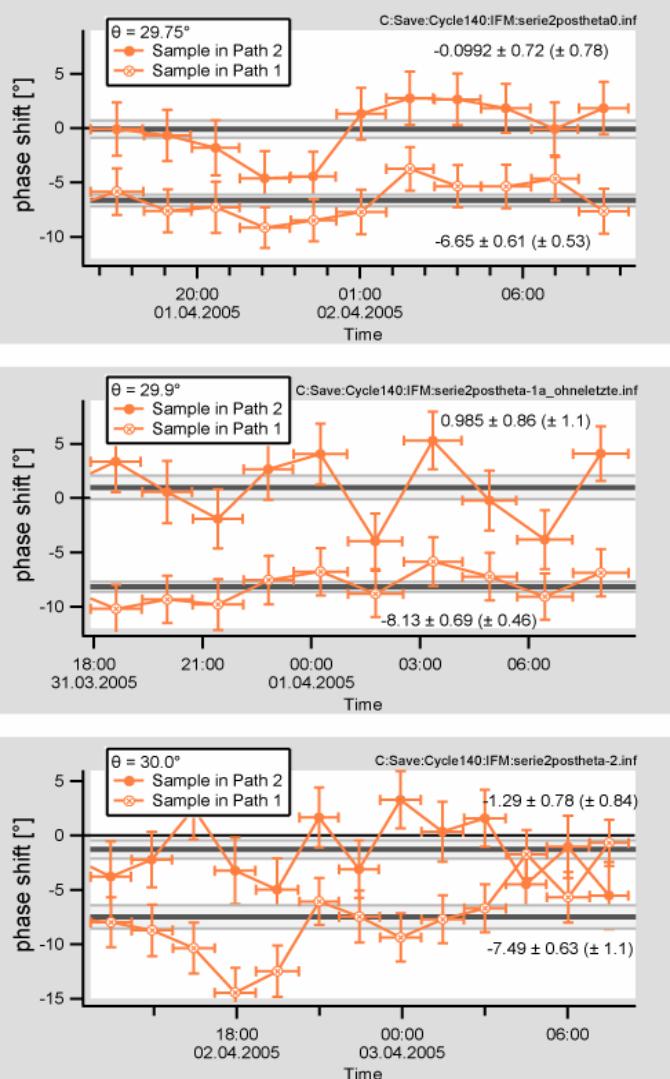
Measured Interference Pattern

Summary of
Measured Phase
Shifts

Calc. $\Delta\Phi = 2.5^\circ$

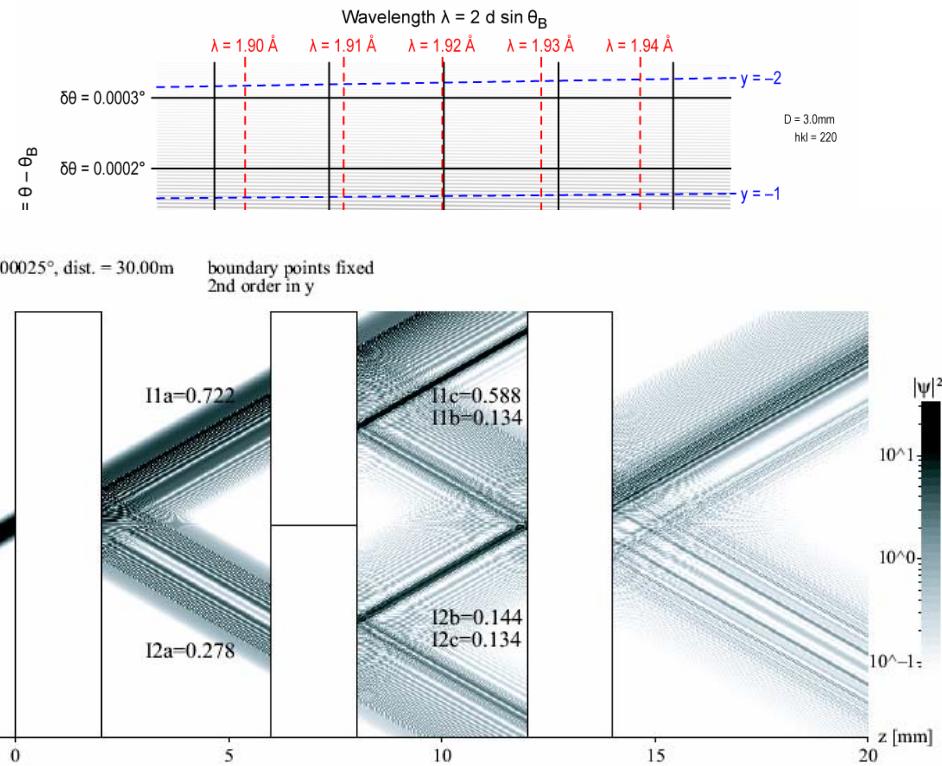
Exp. $\Delta\Phi = 2.8(4)^\circ$

H.Rauch, H.Lemmel, B.Baron, R.Lemmel; Nature 417 (2002) 630



Wave field inside slits

Beam Intensity after Two Laue Reflections



H.Lemmel, R.Loidl, H.Rauch (in progress)

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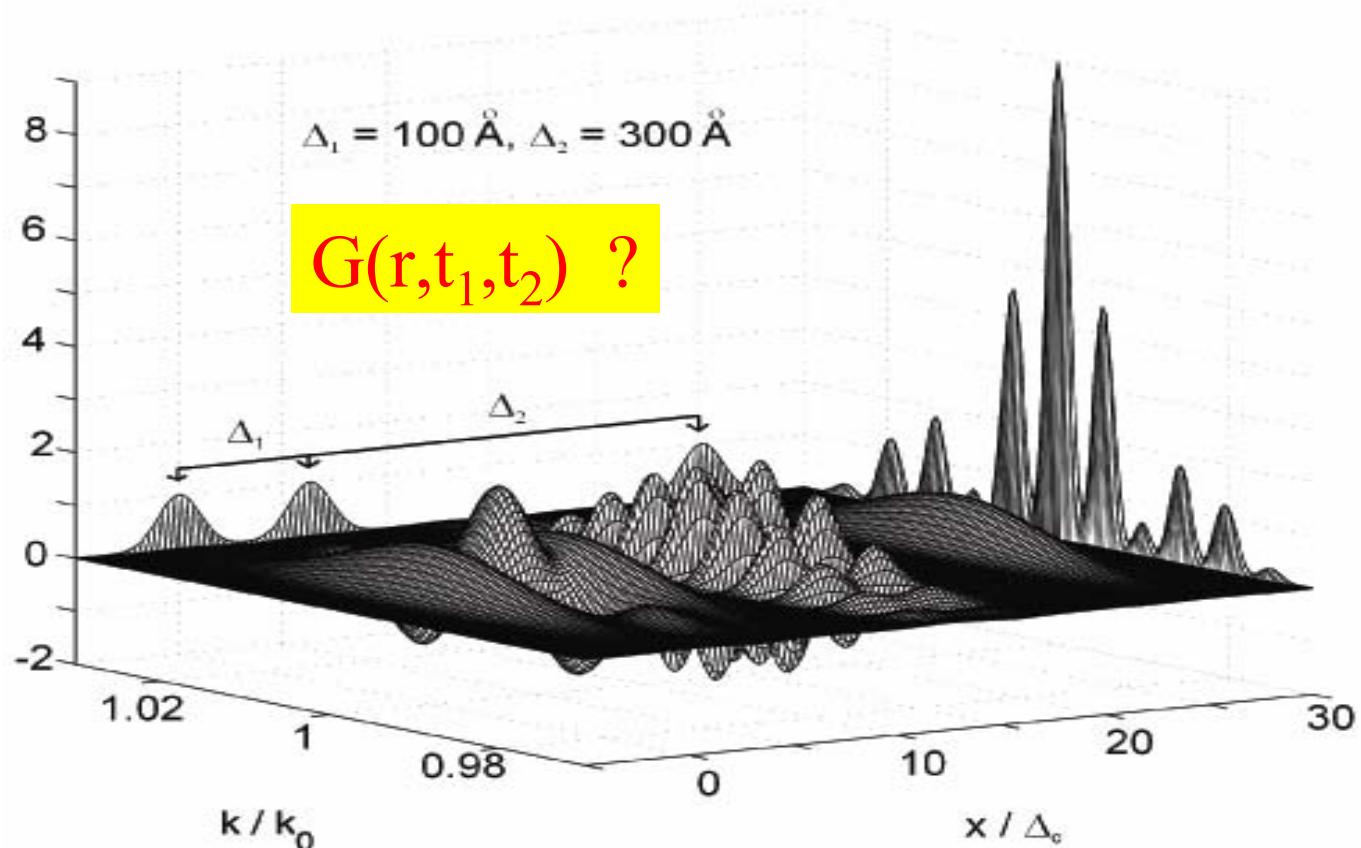
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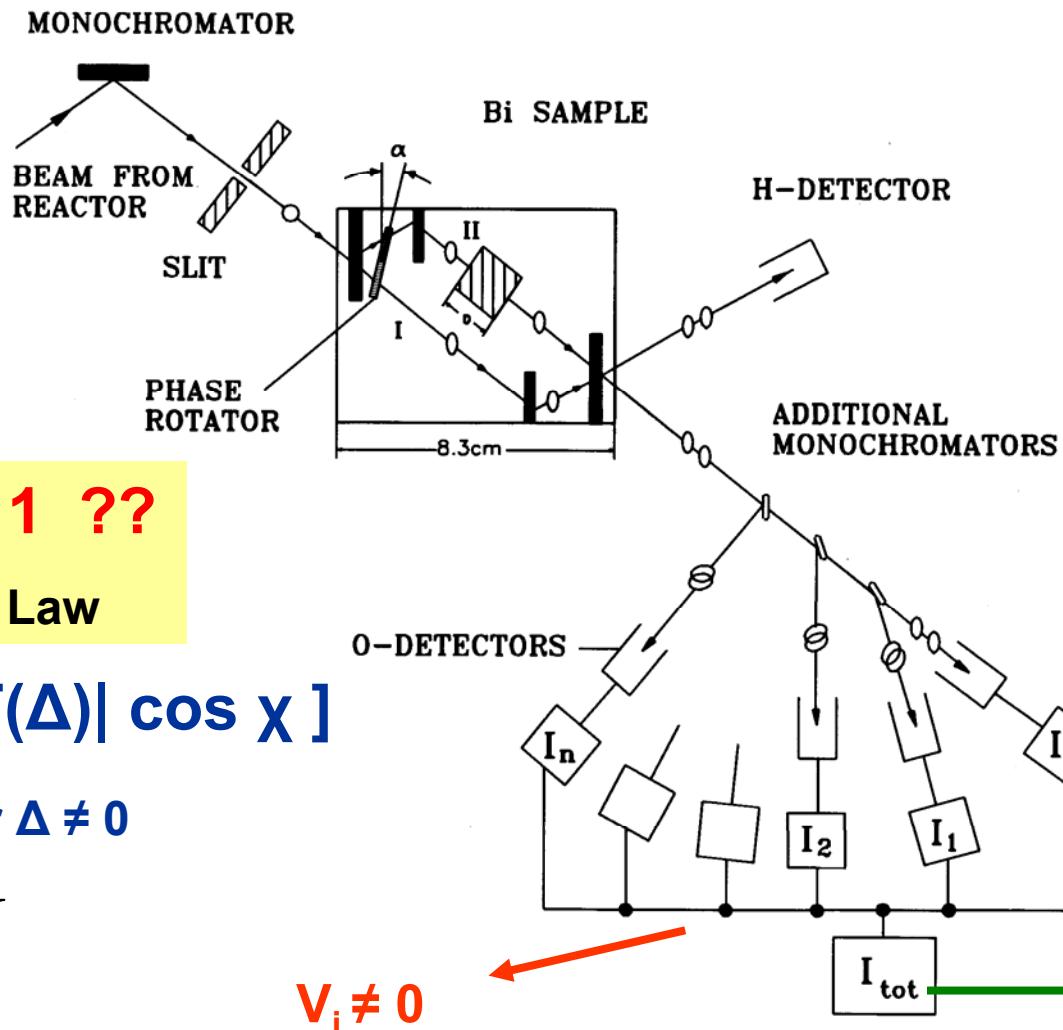
Quantum State Tomography and Topological Phases

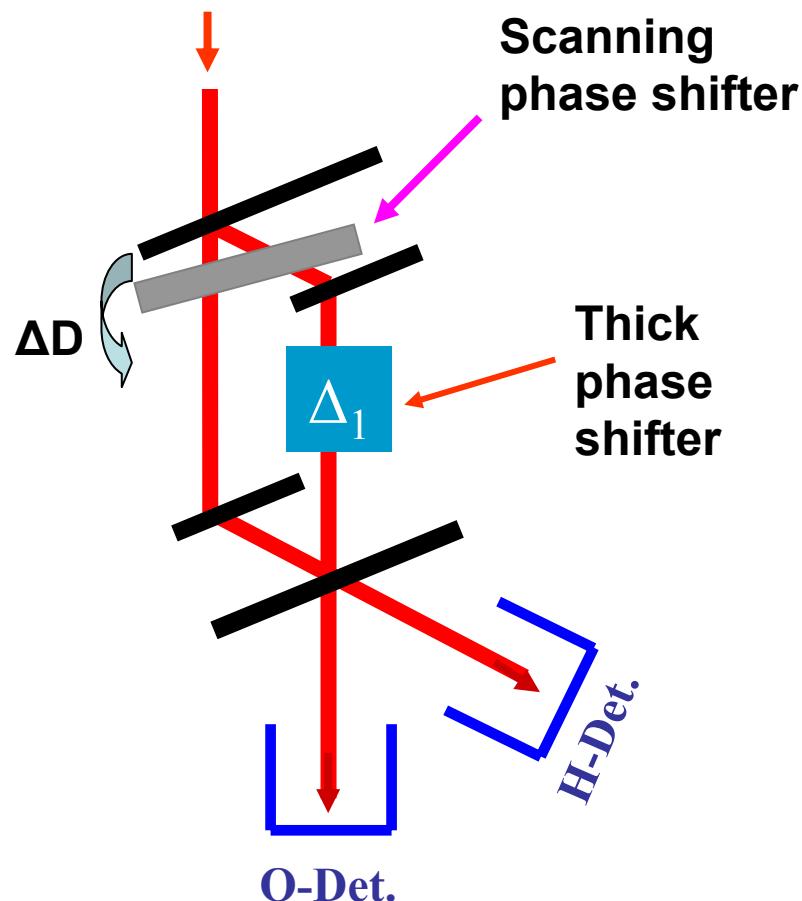
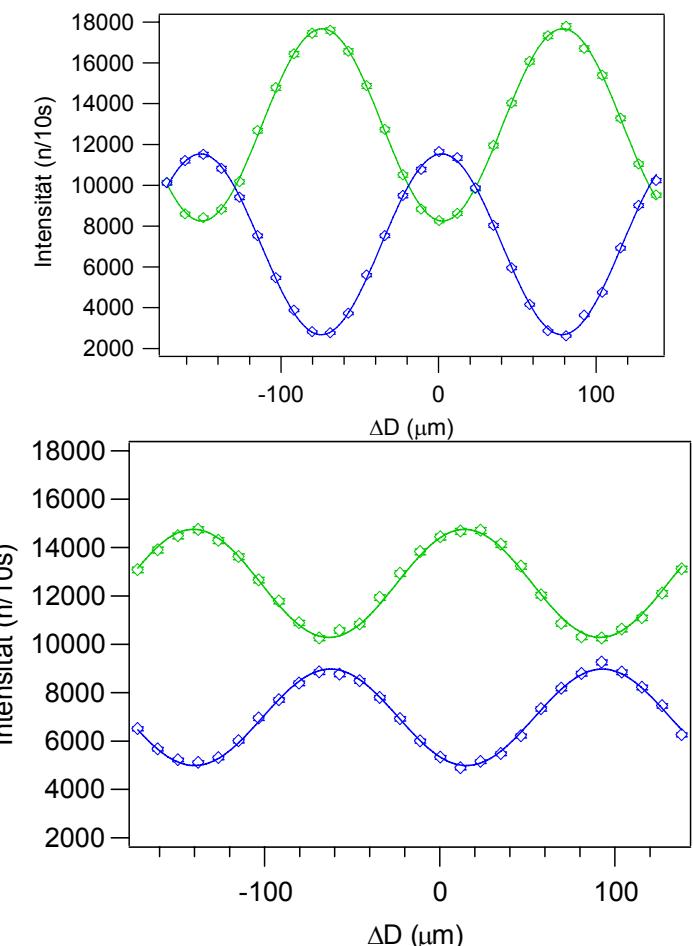
Unavoidable Losses

TRIPLE PEAK WIGNER FUNCTION



H.Rauch, M.Suda 2001





$$I_O \propto |\psi_O^I + \psi_O^{II}|^2 \propto A + B \cos \chi$$

$$\chi = \oint \vec{k} d\vec{s} = (1-n)k D_{eff} \equiv -Nb_c \lambda D_{eff} = \Delta_1 \cdot k$$

