

The Karlsruhe Tritium Neutrino experiment KATRIN





Workshop on Precision Measurements at Low Energy, PSI, Switzerland, January 18&19, 2007 Christian Weinheimer

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- The neutrino mass quest
- The Karlsruhe TRItium Neutrino experiment KATRIN
- KATRIN's background suppression, statistics & systematics
- Summary



bmb+f - Förderschwerpunkt

Astroteilchenphysik

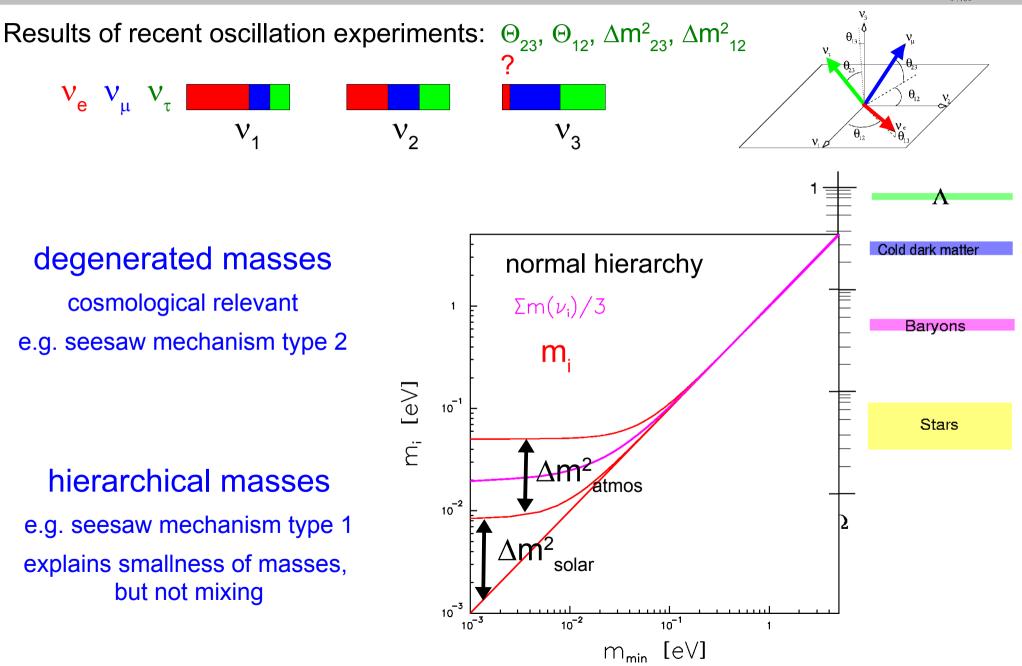
Großgeräte der physikalischen Grundlagenforschung

Need for the absolute ν mass determination

Westfälische Wilhelms-Universität

Münster



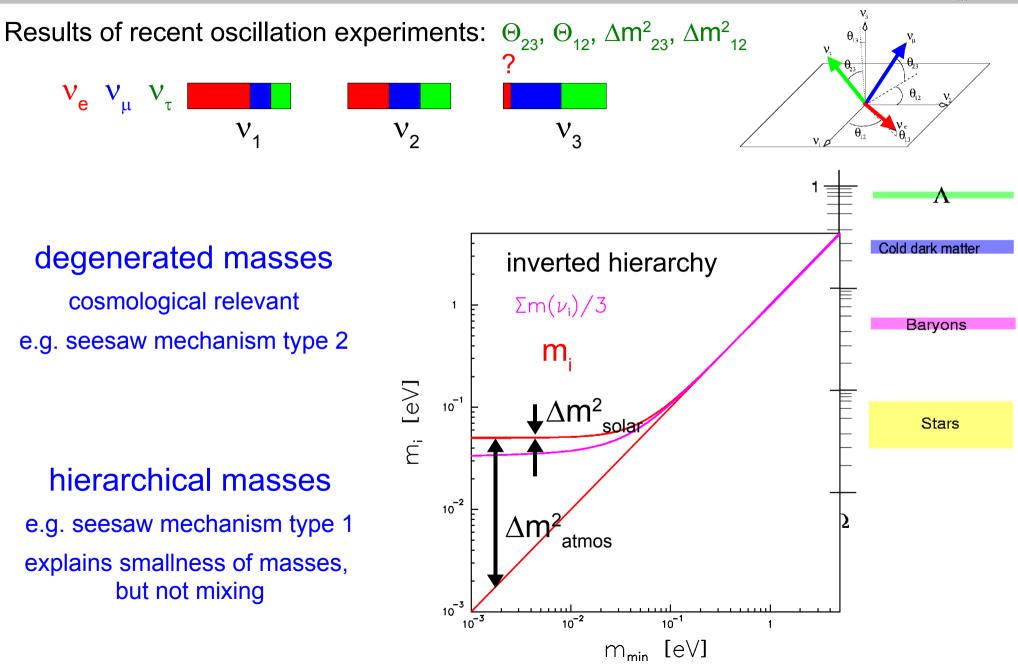


Need for the absolute ν mass determination

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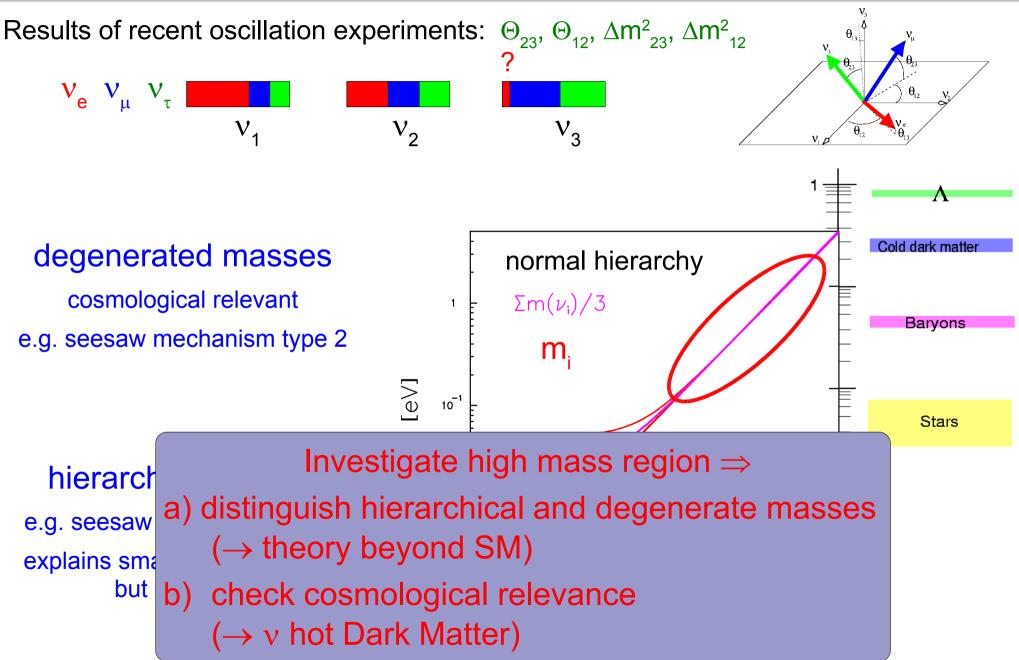




Need for the absolute ν mass determination

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Search for the absolute neutrino mass scale



1) Cosmology

very sensitive, but model dependent current sensitivity: $\Sigma m(v_i) \approx 1 eV$

2) Search for $0\nu\beta\beta$

very sensitive, but needs v to be of Majorana-type sensitive to coherent sum: $m_{ee}(v) = |\Sigma| |U_{ei}^{2}| e^{i\alpha(i)} m(v_{i})|$

 \Rightarrow partial cancelation possible

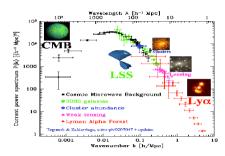
Evidence for $m_{ee}(v) \approx 0.4 \text{ eV}$ (Klapdor-Kleingrothaus et al.)?

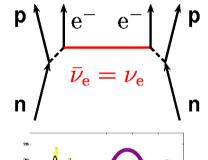
3) Direct neutrino mass determination:

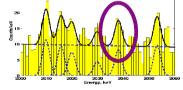
No further assumptions needed, use $E^2 = p^2c^2 + m^2c^4 \Rightarrow m^2(v)$

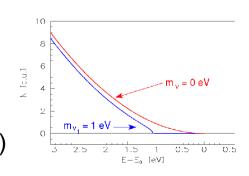
- Time-of-flight measurements (v from supernova) SN1987a (large Magellan cloud) $\Rightarrow m(v_e) < 5.7 \text{ eV}$ (PDG 2003)
- Kinematics of weak decays measure charged decay prod., energy/momentum conserv. $\Rightarrow m^2(v)$ β -decay searchs for m(v_e) - tritium β decay spectrometers

¹⁸⁷Re bolometers











3 complementary ways to the absolute v mass scale



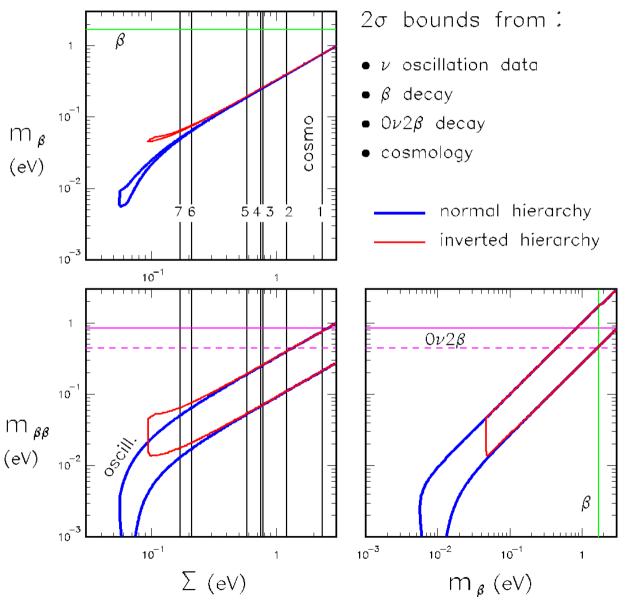
1) cosmology: very sensitive $\Sigma m(v_i)$ but model dependent

2) 0νββ: evidence? Majorana neutrinos

$$m_{ee}(v) = |\Sigma| |U_{ei}^2| e^{i\alpha(i)} m(v_i)|$$

 3) direct neutrino mass determination: no further assumptions no cancellations:

$$m^{2}(v_{e}) = \Sigma |U_{ei}^{2}| m^{2}(v_{i})$$



Fogli et al., hep-ph/0608060



Direct determination of $m(v_e)$ from β decay

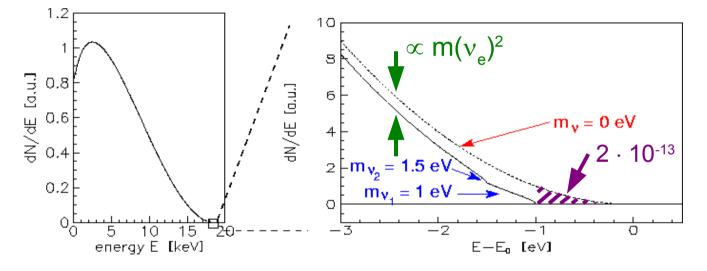


 β decay: (A,Z) \rightarrow (A,Z+1)⁺ + (e⁻)

 β electron energy spectrum:

dN/dE = K F(E,Z) p E_{tot} (E₀-E_e) $\Sigma |U_{ei}|^2 [(E_0-E_e)^2 - m(v_i)^2]^{1/2}$

(modified by electronic final states, recoil corrections, radiative corrections)



oscillation exp: small Δm_{ij}^2 \Rightarrow see only average neutrino mass squared: $m(v_e)^2 := \Sigma |U_{ei}|^2 m(v_i)^2$

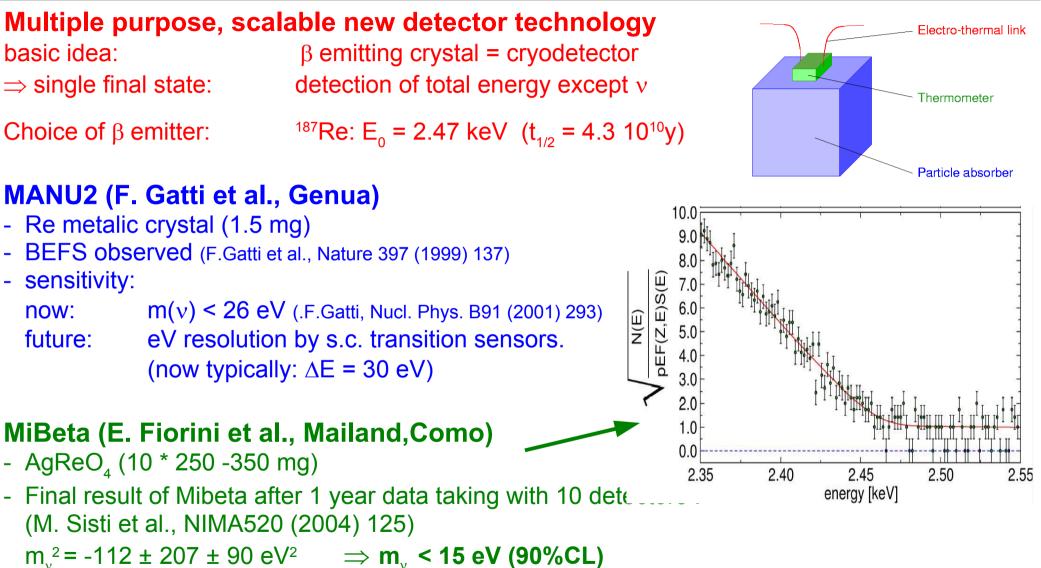
 Need:
 low endpoint energy very high energy resolution & very high luminosity & avery high luminosity & avery low background
 ⇒ Tritium ³H, (¹⁸⁷Re)

 ⇒ MAC-E-Filter (or bolometer for ¹⁸⁷Re)

Cryo bolometer experiments with ¹⁸⁷Re

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"Common future: MARE I: sensitivity 2-3 eV expected by 300 detectors MARE II: better DE, Dt, 50000 detectors: sub-eV sensitivity

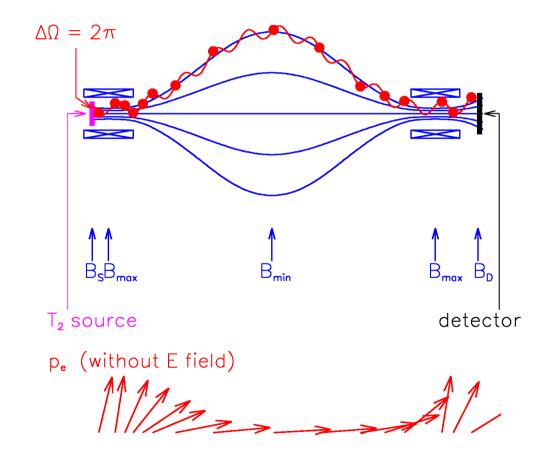


Principle of the MAC-E-Filter



<u>Magnetic Adiabatic Collimation + Electrostatic Filter</u> (A. Picard et al., Nucl. Instr. Meth. 63 (1992) 345)

- Two supercond. solenoids compose magnetic guiding field
- Electron source (T₂) in left solenoid
- e⁻ in forward direction: magnetically guided
- adiabatic transformation: µ = E_⊥/B = const.
 ⇒ parallel e⁻ beam



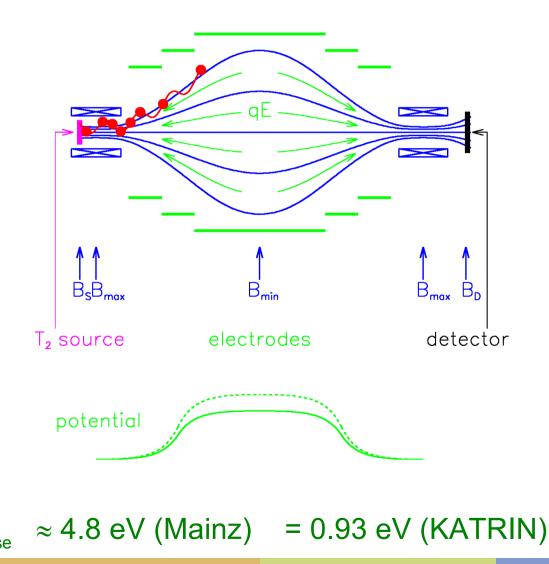


Principle of the MAC-E-Filter



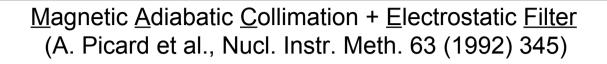
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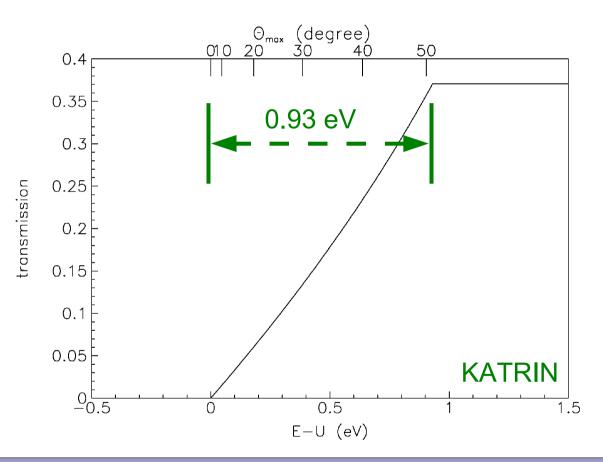
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- adiabatic transformation: μ = E_⊥/B = const.
 ⇒ parallel e⁻ beam
- Energy analysis by electrostat. retarding field $\Delta E = E \cdot B_{min}/B_{max} = E \cdot A_{s,eff}/A_{analyse}$



Principle of the MAC-E-Filter







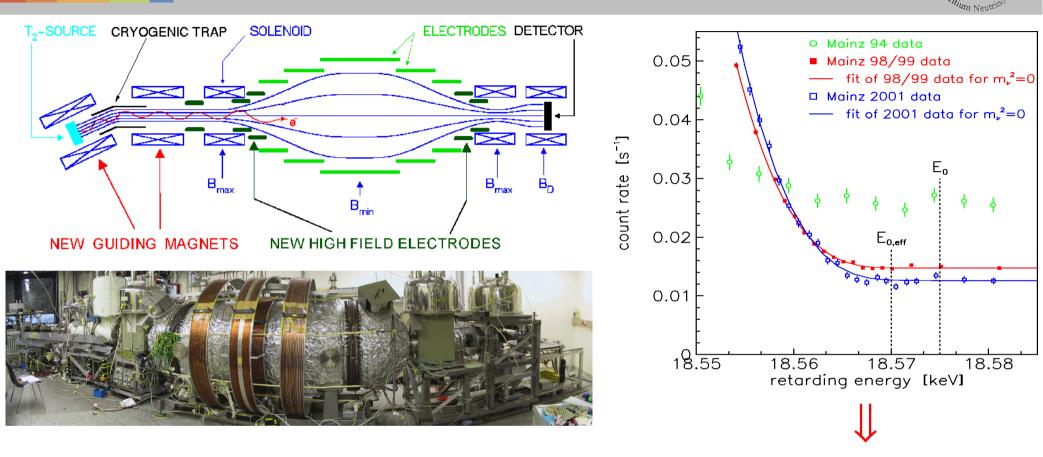
 \Rightarrow sharp integrating transmission function without tails:

 $\Delta E = E \cdot B_{min} / B_{max} = E \cdot A_{s,eff} / A_{analyse} = 0.93 \text{ eV}, \text{ KATRIN}$ (4.8 eV, Mainz)

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The Mainz Neutrino Mass Experiment Phase 2: 1997-2001





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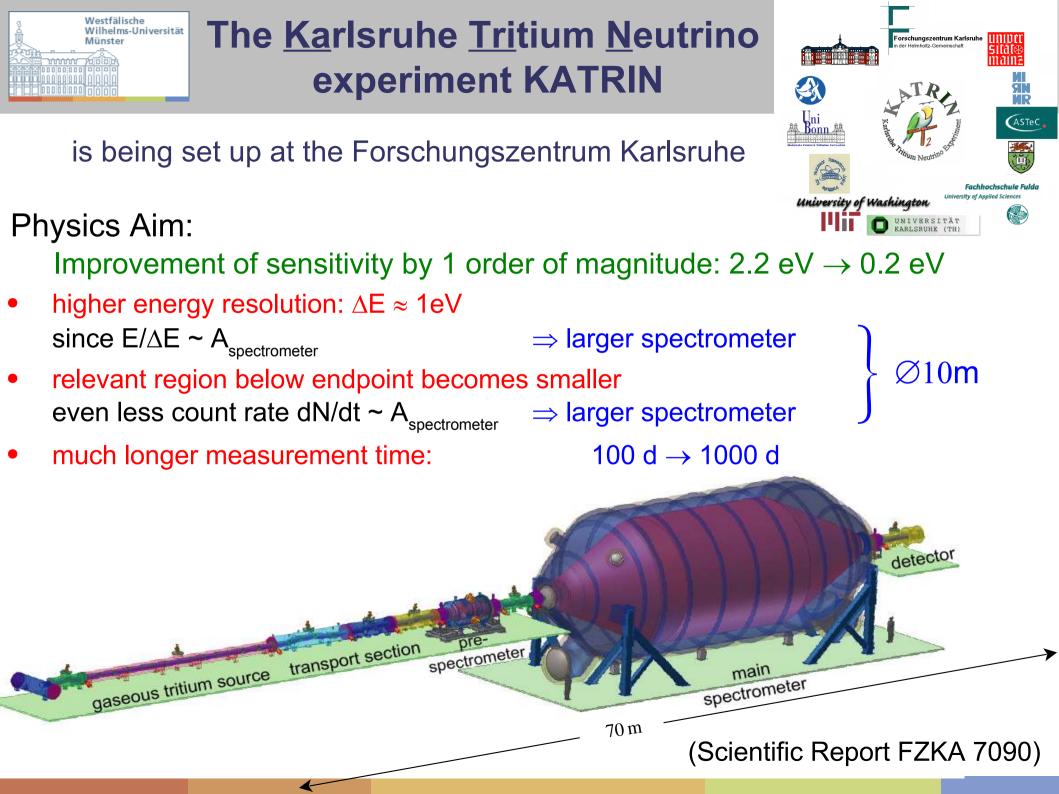
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After all critical systematics measured by own experiment (inelastic scattering, self-charging, neighbor excitation):

 $m^{2}(v) = -0.6 \pm 2.2 \pm 2.1 \text{ eV}^{2} \implies m(v) < 2.3 \text{ eV} (95\% \text{ C.L.})$

C. Kraus et al., Eur. Phys. J. C 40 (2005) 447

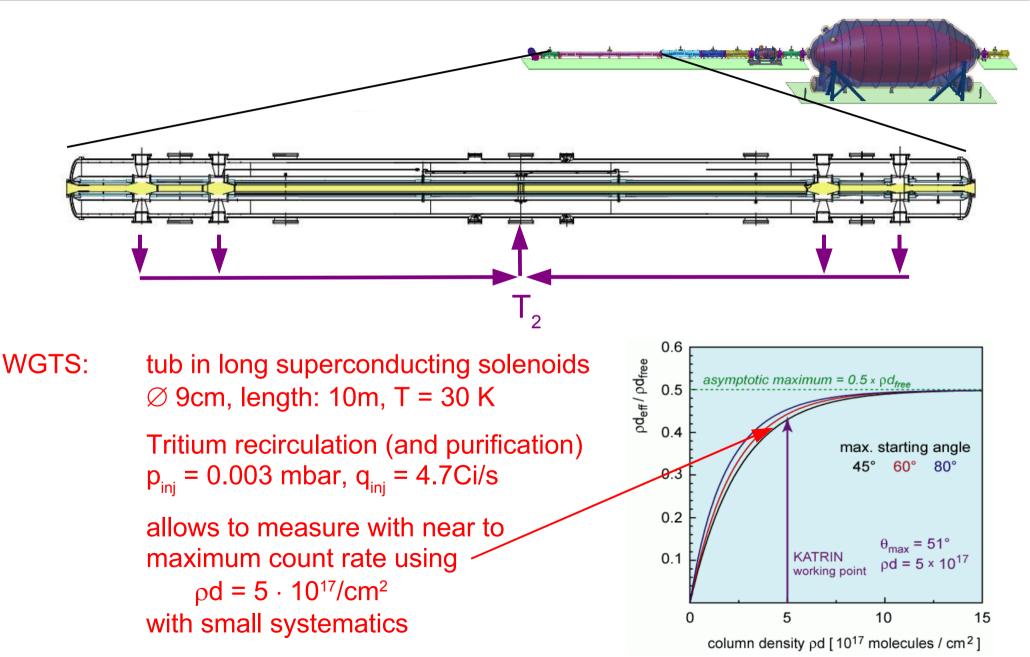


Molecular Windowless Gaseous Tritium Source WGTS

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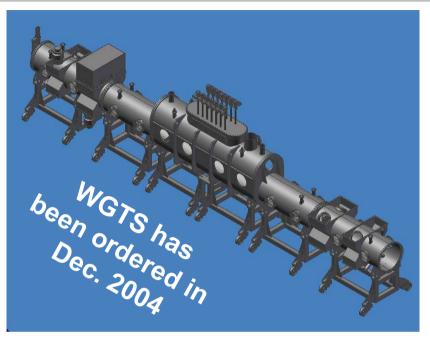




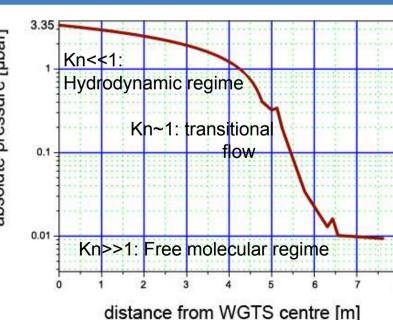


Molecular Windowless Gaseous Tritium Source WGTS





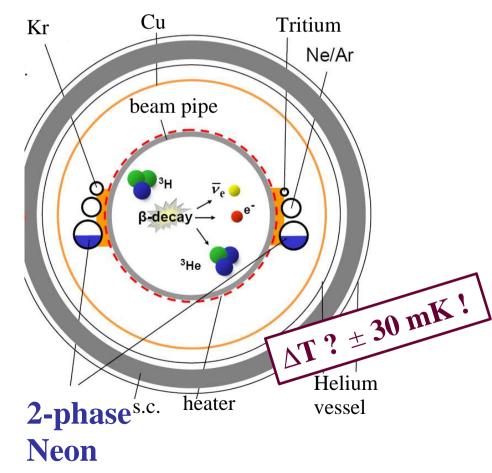
absolute pressure [µbar]



Conceptional design

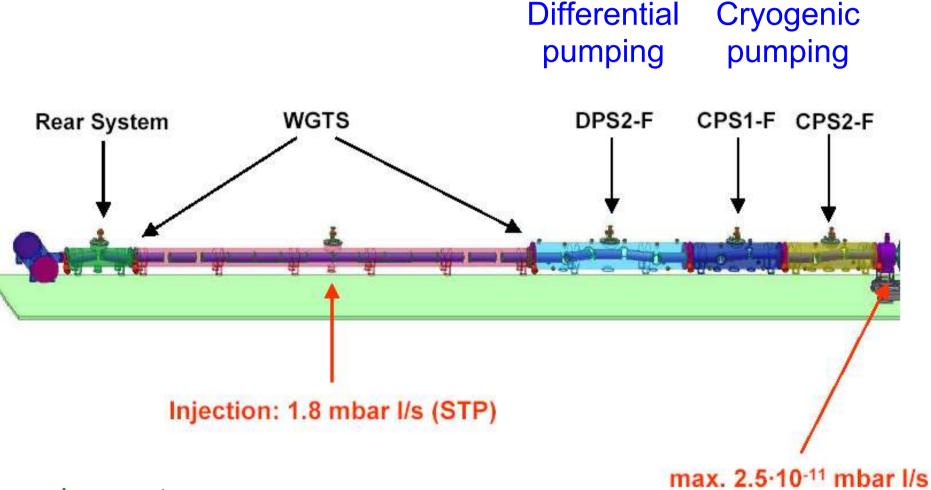
2 phase Neon cooling with operating temperature: 27-28 K

- **spatial** (homogeneity): ± 0.1%
- time (stability/hour): \pm **0.1%**



Transport and differential & cryo pumping sections





requirements:

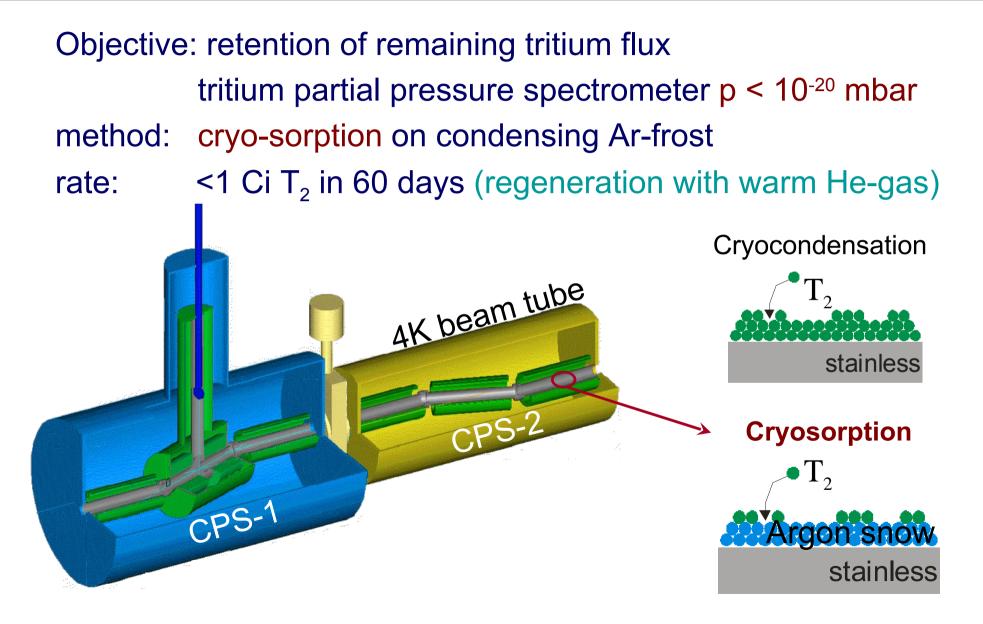
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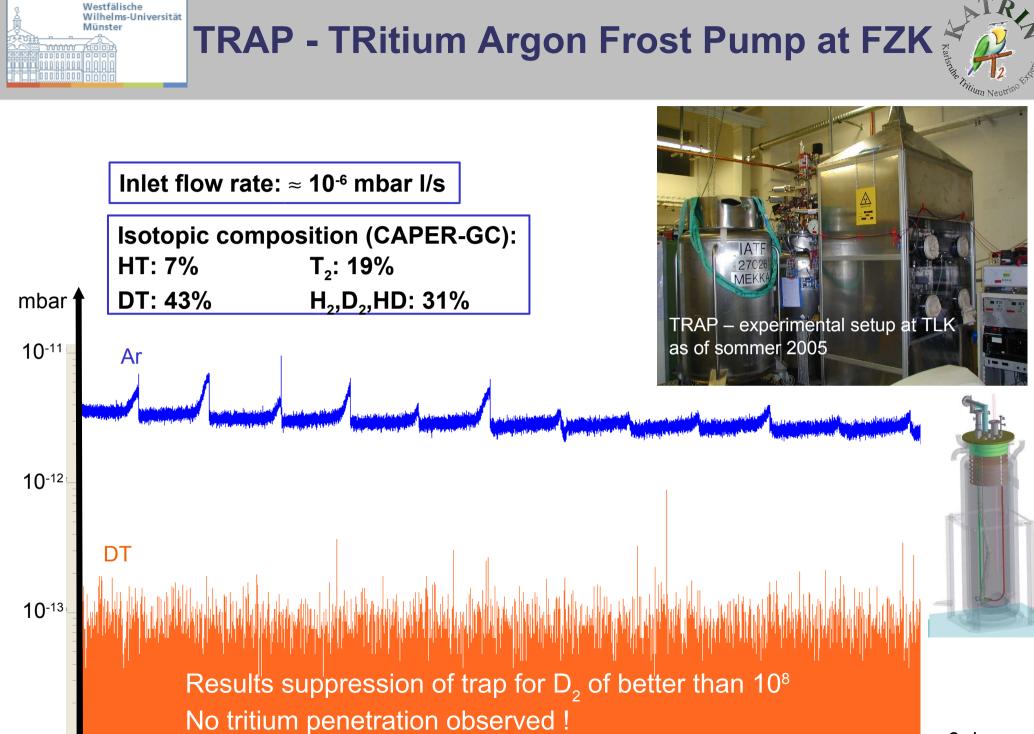
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- adiabatic electron guiding
- T_2 reduction factor of ~10¹¹









▶6 days



and main spectrometer

Pre

Main spectrometer:

- Ø Ø 10m, length 24m
 - \Rightarrow large energy resolution: $\Delta E = 0.93 \text{ eV}$
 - \Rightarrow high luminosity: L = A_{Seff} $\Delta\Omega/4\pi$ = A_{analyse} Δ E/(2E) = 20 cm²
- ultrahigh vacuum requirements (background) p < 10⁻¹¹ mbar
- "simple" construction: vacuum vessel at HV + "massless" screening electrode

1010

Pre spectrometer

- Transmission of electron with highest energy only (10⁻⁷ part in last 100 eV)
 - \Rightarrow Reduction of scattering probaility in main spectrometer
 - \Rightarrow Reduction of background
- only moderate energy resolution required: $\Delta E = 80 \text{ eV}$
- test of new ideas (XHV, shape of electrodes, avoid and remove of trapped particles, ...)

10.2



Detector



task: detection of transmitted ß-decay electrons segmented PIN-diode with high energy resolution ($\Delta E = 1 \text{ keV}$) 44 x 44 mm² record radial profile of flux tube 64 segments 5x5 mm² aim: background minimisation, systematic effects bonded onto ceramics with FET stage \Rightarrow post-acceleration to place signal line at lower intrinsic background detector magnet 3-6 20 flux tube -20 20 design: radially segmented shielding & electrode low-level Si-PIN diode array (kV-acceleration) detector veto ~150 pixels with A=100 cm²



Technical challenges

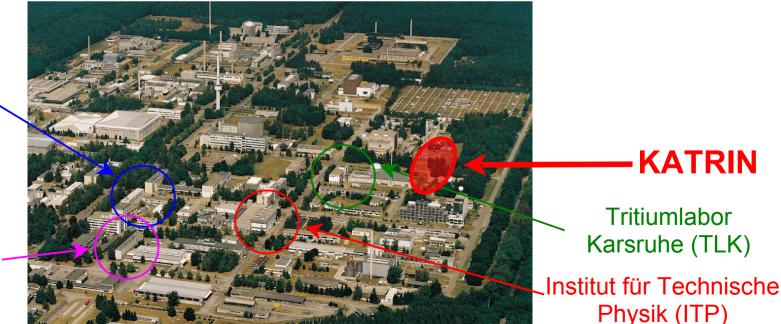


- Recirculation and purification of tritium to a large extent (kCi)
- \approx 30 superconducting solenoids
- UHV (< 10⁻¹¹ mbar) in huge volume (1000m²)
- HV calibration and stability on ppm level
- High resolution detectors
- •

 \Rightarrow ideal place: Forschungszentrum Karlsruhe/Germany

Inst. f. Kernphysik (IK)

Inst. f. Prozessdatenverarbeitung und Elektronik (IPE)





KATRIN location at Forschungszentrum Karlsruhe





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Vacuum tests and inner electrode system of pre spectrometer





ground electrode wire electrode + solid cones





dry air compartment to allow cooling at -20°C: outgasing rate < 10⁻¹³ mbar l/s cm² with getter pumps (NEG, 10000 l/s): p< 10⁻¹¹mbar ⇒ better than KATRIN requirements





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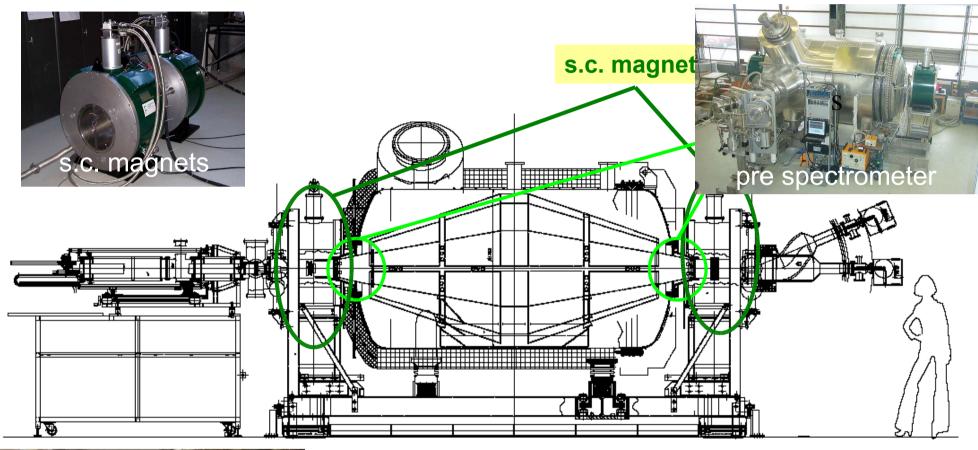
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Electromagnetic design tests Wilhelms-Universität at the pre spectrometer have just started

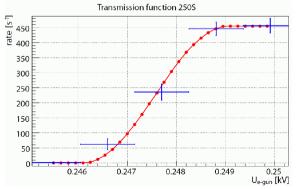






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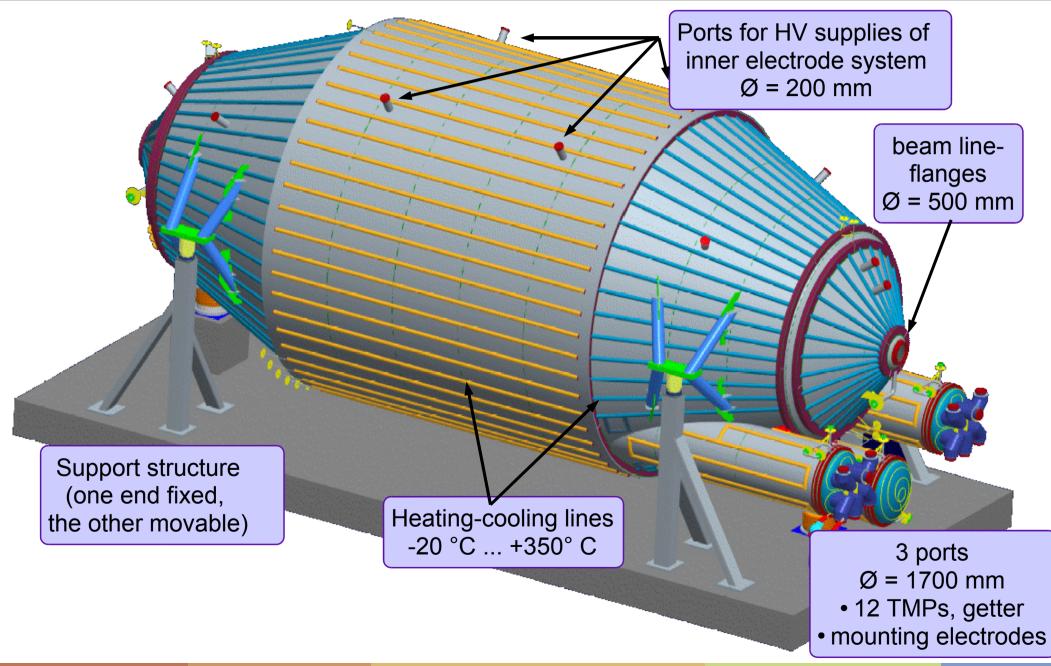
Main spectrometer 3dim model with heating-cooling system

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Main spectrometer vessel construction at MAN DWE

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Münster





Main spectrometer vessel construction at MAN DWE

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Main spectrometer vessel construction at MAN DWE

Wilhelms-Universität



August 2006: - construction of main spectrometer vessel has been finished - vessel has passed leak test successfully ! ⇒ world's biggest XHV vessel ever been build !

It is very big and heavy ... \Rightarrow a 8500 km long detour







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10

A KIN



29.10.06

29.10.06





1) Low background:

Mainz experiment: most background from spectrometer but KATRIN spectrometer is much bigger! ⇒ need something new !

2) Huge statistics: optimized source & large spectrometer

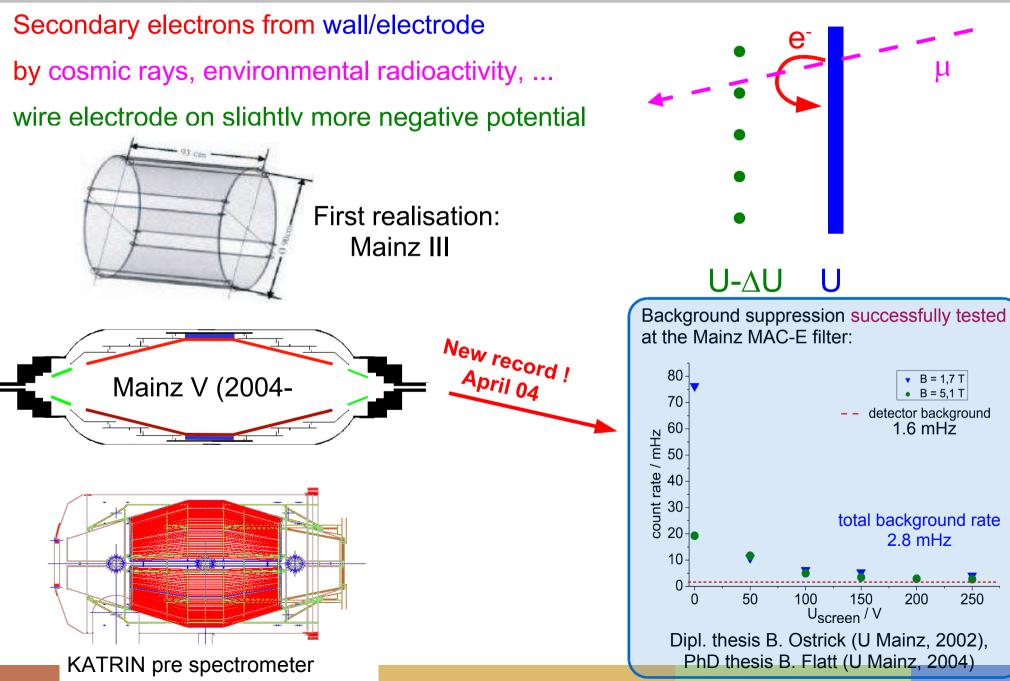
3) Systematic uncertainties:

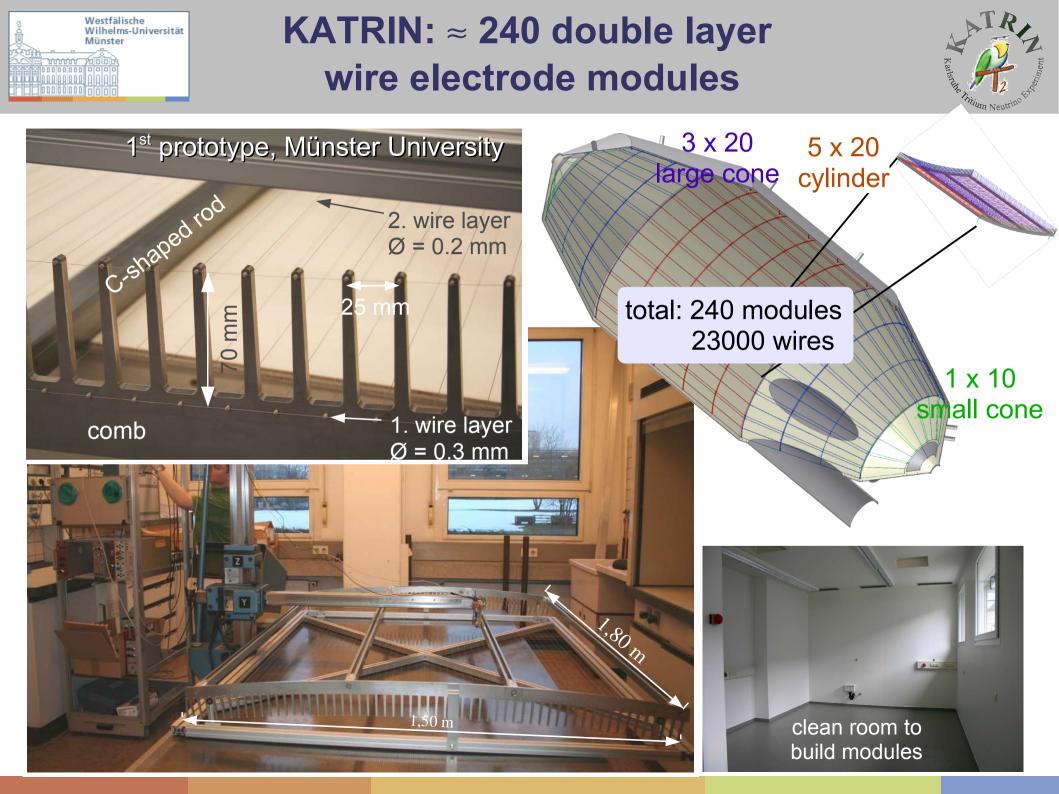
need to be very small !



Background reduction by a "massless" wire electrode



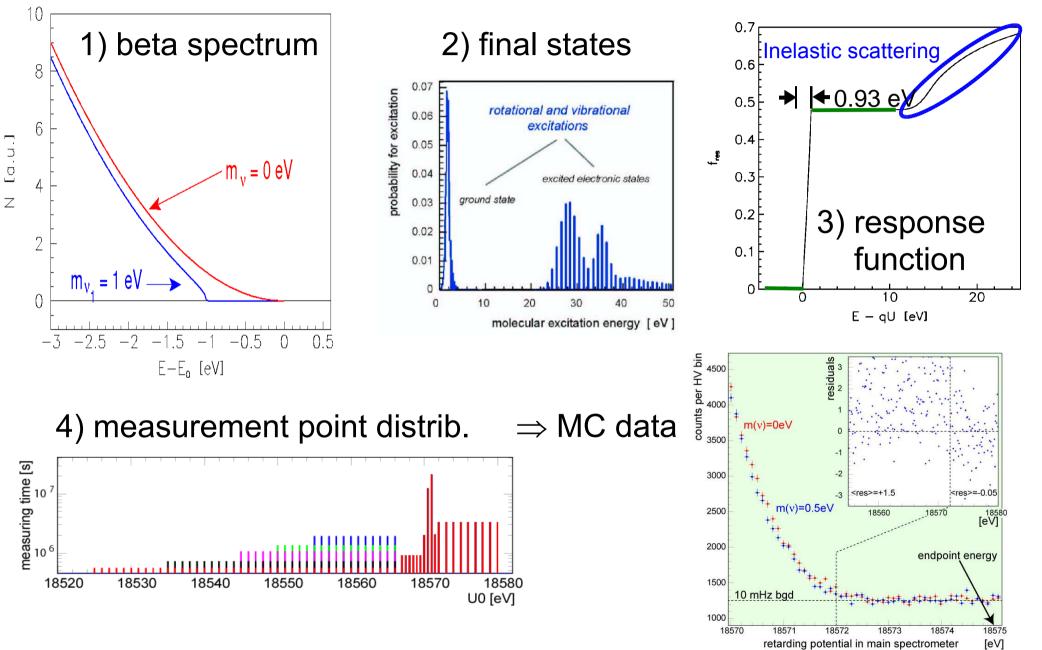






2) Statistics

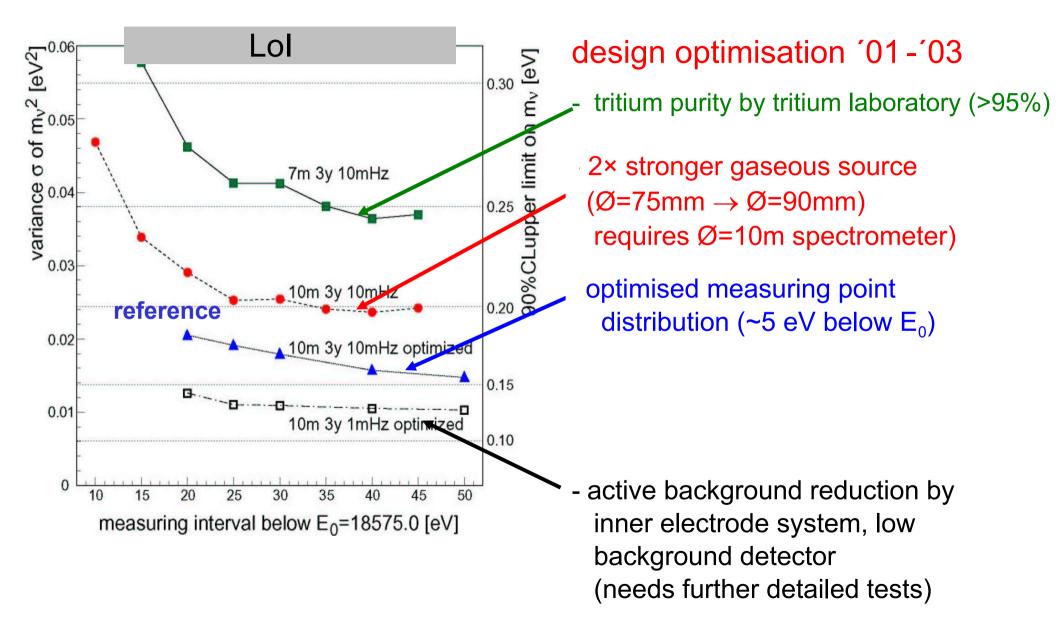






KATRIN's statistical uncertainty









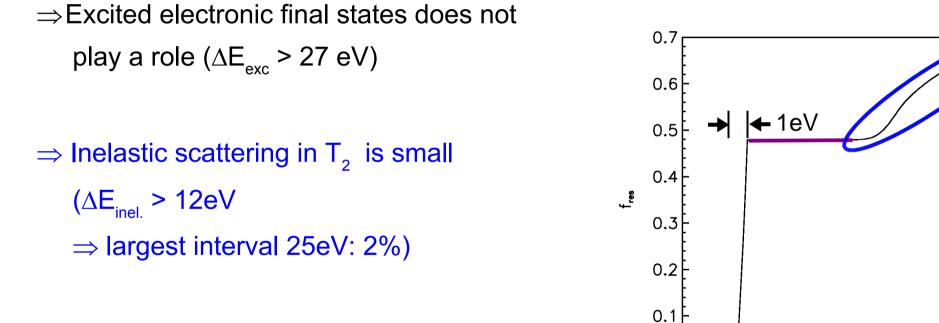
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10 E – qU [eV]

0

As smaller m(v)

as smaller the region of interest below endpoint E₀



⇒ One well-defined final state (similiar to cryo detectors)

Is only true, since MAC-E-Filter response function has no tails





any not accounted variance σ^2 leads to negative shift of m_v^2 : $\Delta m_v^2 = -2 \sigma^2$

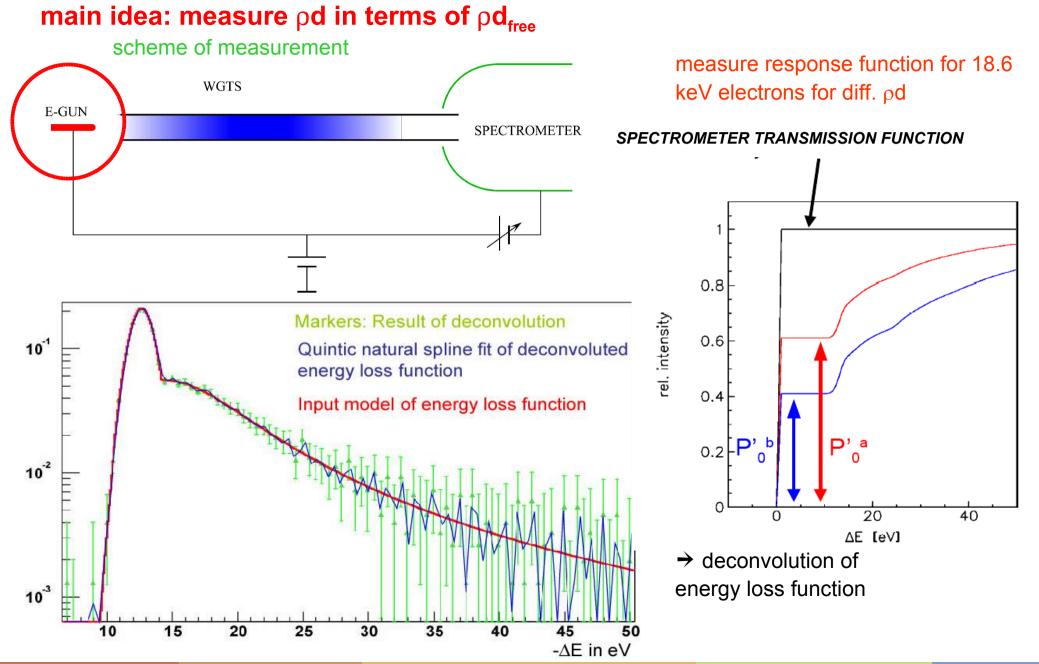
- 1. inelastic scatterings of ß's inside WGTS
 - requires dedicated e-gun measurements, unfolding techniques for response fct.
- 2. fluctuations of WGTS column density (required < 0.1%)
 - rear detector, Laser-Raman spectroscopy, T=30K stabilisation, e-gun measurements
- 3. transmission function
 - spatial resolved e-gun measurements
- 4. WGTS charging due to remaining ions (MC: φ < 20mV)
 inject low energy meV electrons from rear side, diagnostic tools available
- 5. final state distribution
 - reliable quantum chem. calculations
- 6. HV stability of retarding potential on ~3ppm level required
 precision HV divider (PTB), monitor spectrometer beamline

a few contributions with each: $\Delta m_v^2 \le 0.007 \text{ eV}^2$



Determine energy loss function and source column density $\rho\,\text{d}$

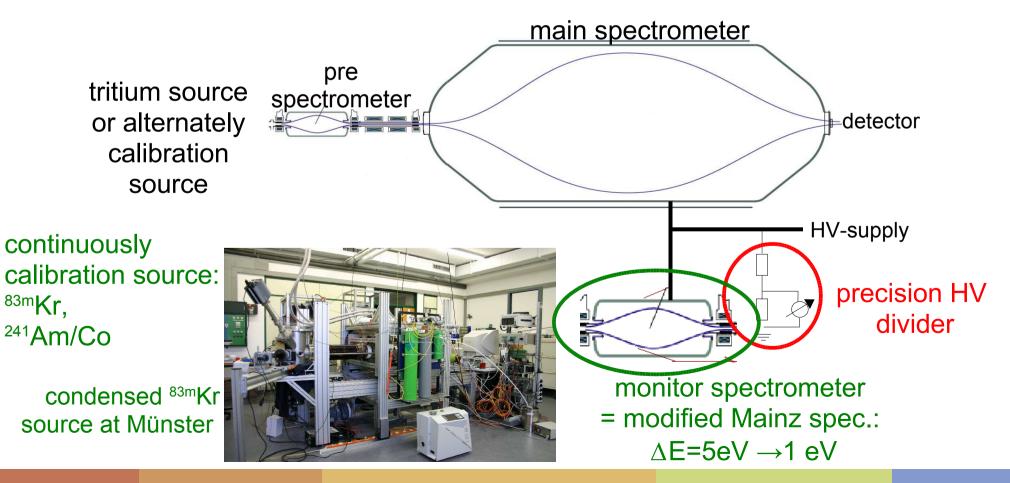








- Measure HV by precision HV divider
- Lock retarding HV by measuring energetically well-defined (atomic/nuclear standard) and sharp electron line with monitor spectrometer





KATRIN precision high voltage divider a accoration with DTD

Proupophuo

holtz-Gemeinschaf

1090 ICB Řež EXP-01/200

KATRIN Design Report 2004

KATRIN Collaboratic

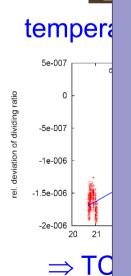
Februar 2005

der Heln



- higher T2 purity
 - larger statistics
 - optimized measurement point distribution
 - smaller systematic uncertainties

 \Rightarrow sensitivity on m(v) $\approx 0.20 \text{ eV/c}^2$



(about equal contribution from stat. and syst. uncertainties) (90% C.L. upper limit for $m(v_a) = 0$)

> $m(v_e) = 0.30 eV$ observable with 3σ $m(v_{a}) = 0.35eV$ observable with 5σ





Summary & Outlook



Absolute neutrino mass scale is needed for particle physics & astrophysics/cosmology by direct neutrino mass measurement (less model dependent & complementary)

KATRIN will become sensitive on the $m(v_e)$ down to 0.2 eV: $m(v_e) < 0.2 eV$ or $m(v_e) > 0 eV$ (for $m(v_e) \ge 0.30 eV$ @ 30)

- 2009/10 complete setup and commisioning
- 2010 start of data taking
- 2011 first results
- 2015 finish data taking

