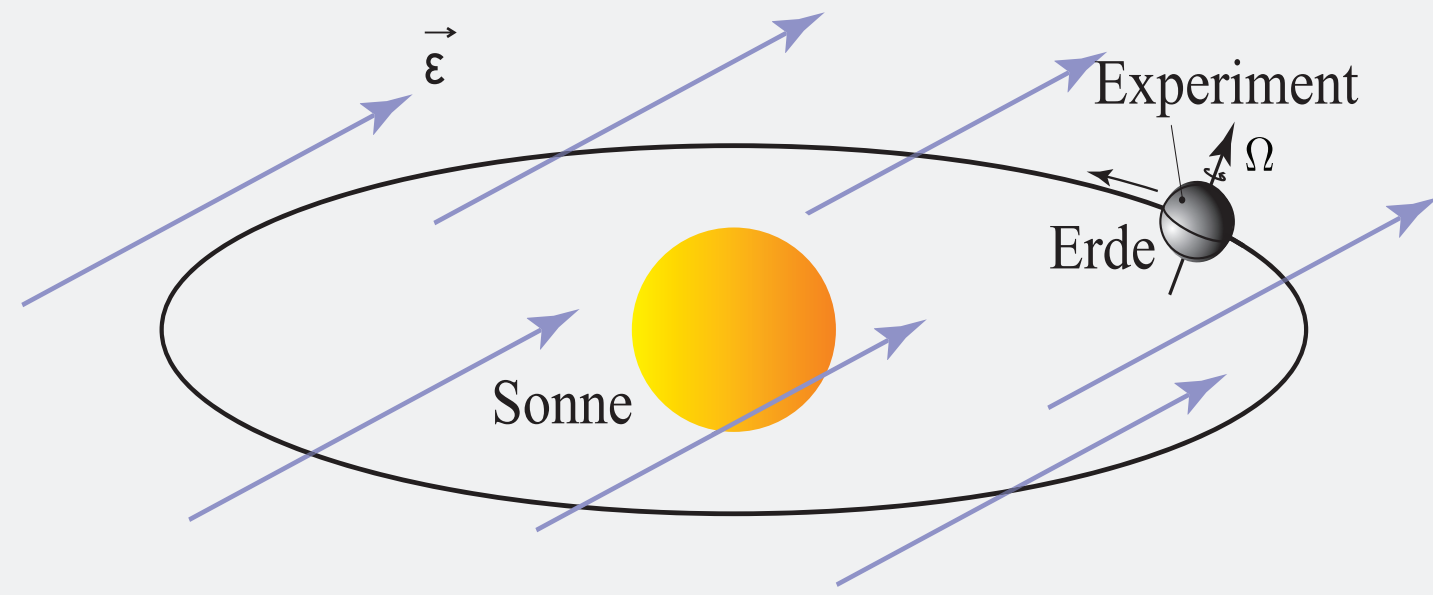


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Theoretical Motivation



With a clock comparison test the Lorentz Invariance of the Theory of Relativity can be checked experimentally. For this test, we search for a possible dependence of clock frequency of an atomic clock on its orientation.

Our atomic clock is the free precession of polarized Helium-3 in a constant magnetic field \vec{B} . An orientation dependence of the clock frequency can be written as a coupling to a background field $\vec{\epsilon}$, then the Hamiltonian is given by:

$$H = -\vec{\mu} \cdot \vec{B} - \beta \vec{\sigma} \cdot \vec{\epsilon}$$

The precession frequency is then given by

$$\nu = 2\mu B + 2\beta \cos(\vec{\epsilon}, \vec{B})$$

and changes with the time of the sidereal day as the clock moves with the surface of the earth.

The commonly used parametrization of Lorentz-Violating Effects from Kostelecky et al. allows to compare the sensitivity of different experiments. It also shows that Lorentz-Violation is likely accompanied by CPT violation.

Basic principle of the measurement

Our idea: Use the free precession of polarized ^3He gas in a glass cell as an atomic clock. Its frequency is given by the magnetic field and the contribution of a possible Lorentz-Violating effect.

- * Main reason to use ^3He :
 - longitudinal (T_1) and transverse (T_2^*) relaxation times can be made very long (> 100 h)
 - ^3He nuclear Polarization can be made very high, reaches $P_{\text{He}} \sim 80\%$

- * Precession frequency only depends on the magnetic field; no systematic effects due to light shifts, feedback etc.

- * Detection of Spin Precession by means of SQUIDS. Sensitivity $\sigma_{\text{B,SQUID}} \sim 4.5 \text{ fT}/\sqrt{\text{Hz}}$

- * Expected signal strength for a spherical cell ($p_{\text{He}} = 5 \text{ mbar}$, $P_{\text{He}} = 30\%$):

$$B_{\text{rot}} \sim \frac{\mu_0}{4\pi} \cdot \frac{2M}{r^3} \sim 30 \text{ pT @ } r = 5 \text{ cm (distance from the cell center)}$$

- * Longitudinal Relaxation Time T_1 :

$$\frac{1}{T_1} \sim \frac{1}{T_{1,\text{wall}}}$$

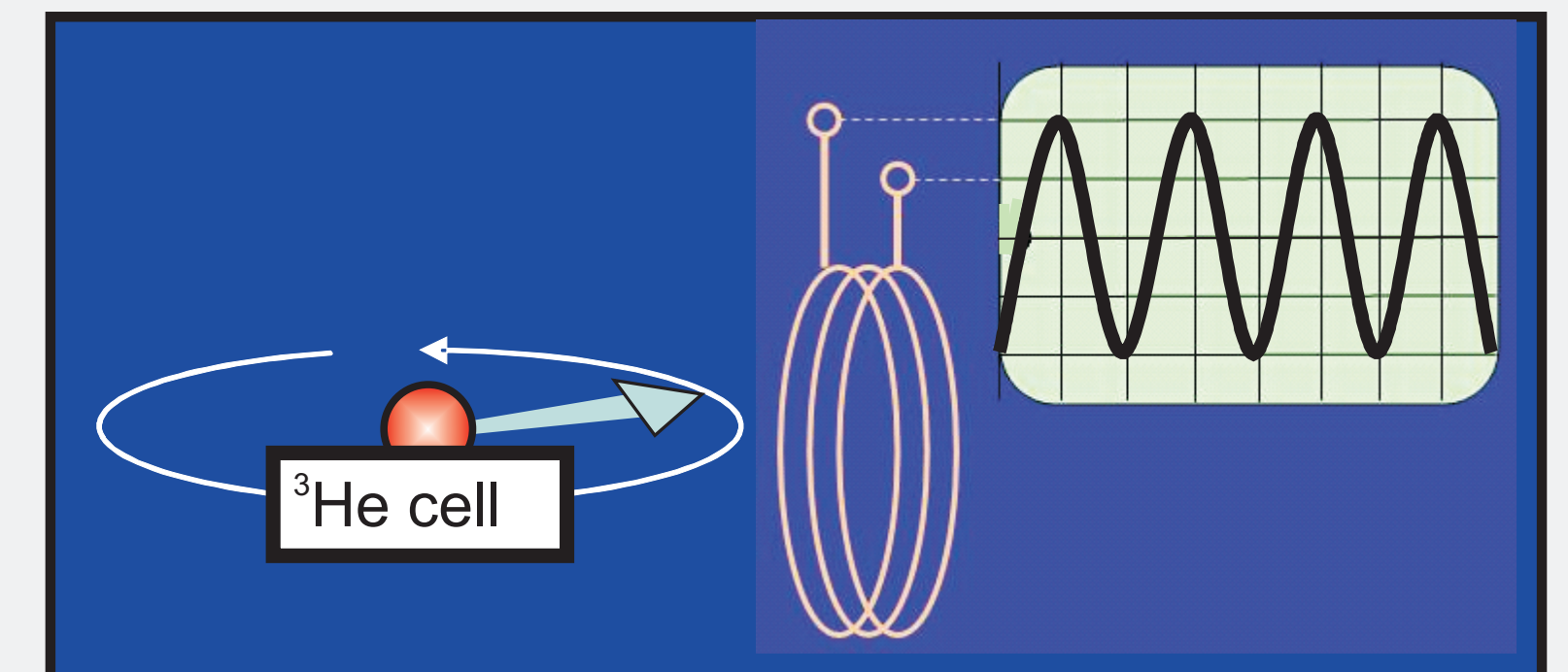
- * Transverse Relaxation Time T_2^* :

$$\frac{1}{T_2^*} = \frac{1}{2T_{1,\text{wall}}} + \frac{8\gamma^2 R^4}{175D} \left[\left(\frac{\partial B_z}{\partial x} \right)^2 + \left(\frac{\partial B_z}{\partial y} \right)^2 + \left(\frac{\partial B_z}{\partial z} \right)^2 \right] + D \cdot \frac{|\nabla B_x|^2 + |\nabla B_y|^2}{B_0^3} \cdot \frac{1}{(x_{11}^2 - 2)(1 + D^2 x_{11}^4 / \gamma^2 B_0^2 R^4)}$$

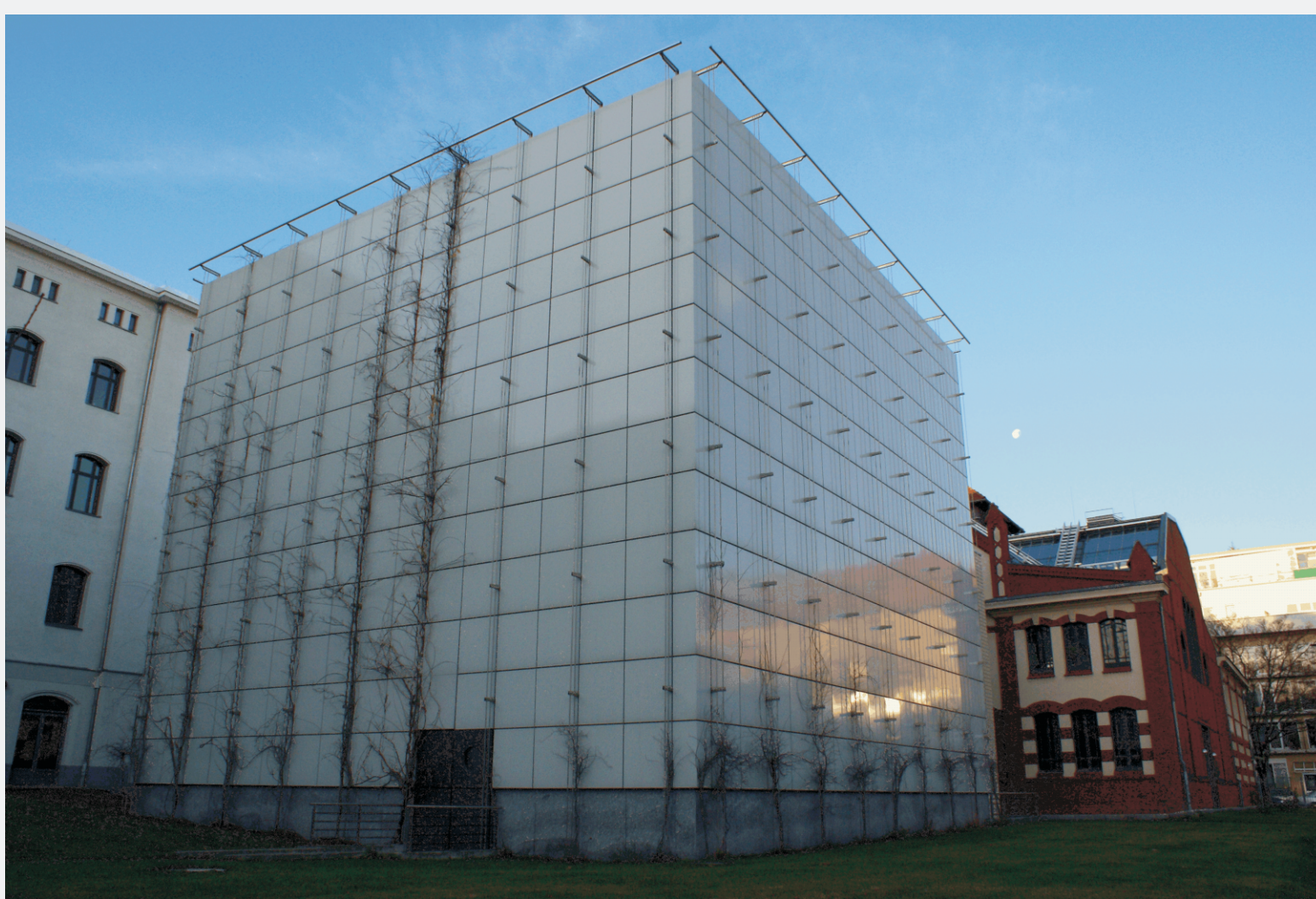
Here D is the diffusion coefficient, R is the cell radius, $x_{11} = 2.08$

- * Estimated value:

$T_2^* \sim 100 \text{ h @ absolute magnetic field gradient of some pT/cm, } B_0 \sim 1 \mu\text{T}$



Our first experimental setup at the PTB Berlin

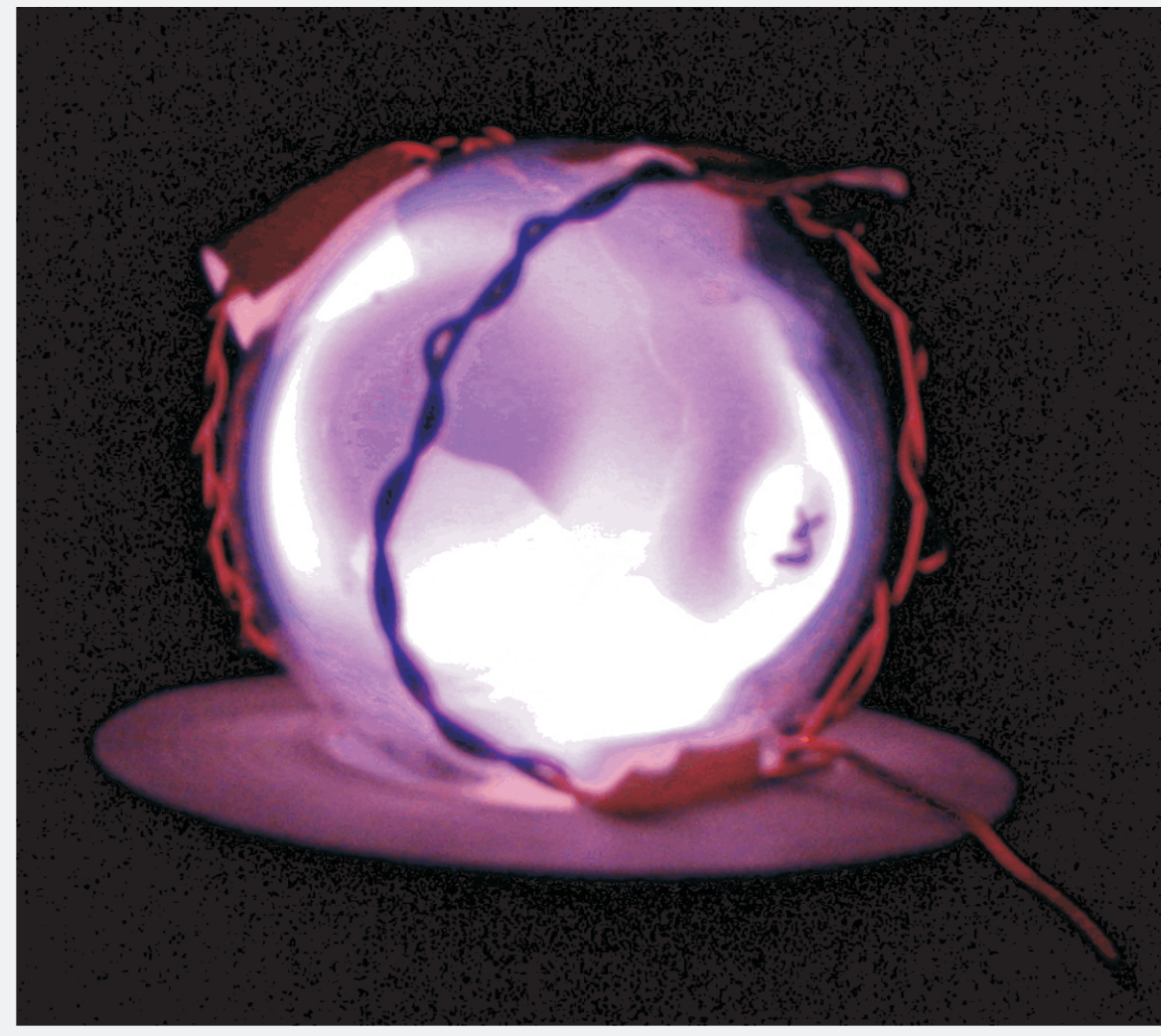


BMSR-2: Anti-Magnetic Screen (PTB Berlin)

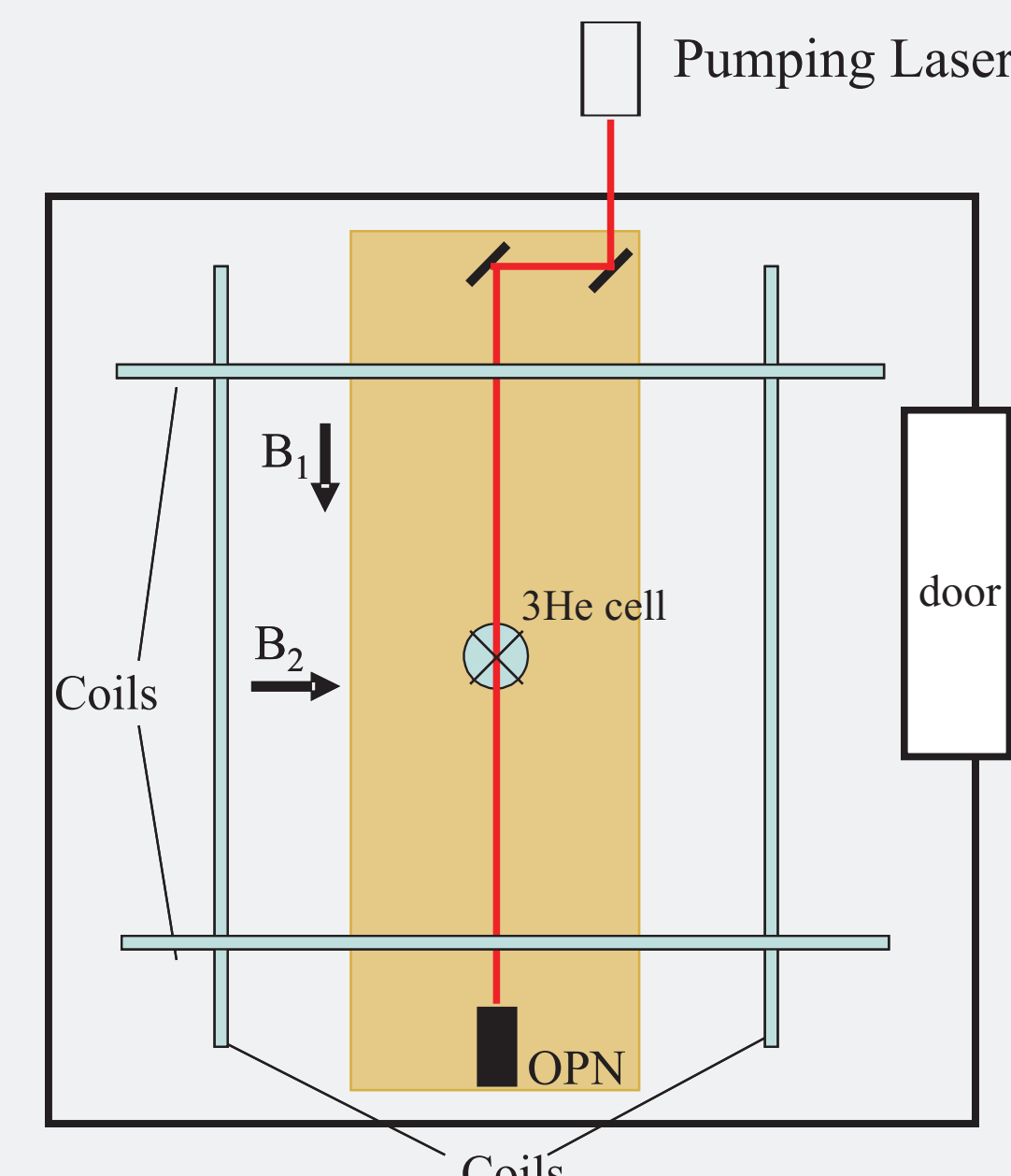
8 layers of μ -metal, HF shielding

rest field: 400 pT

gradient in rest field: several pT/cm in a volume of $\sim 2 \text{ m}^3$

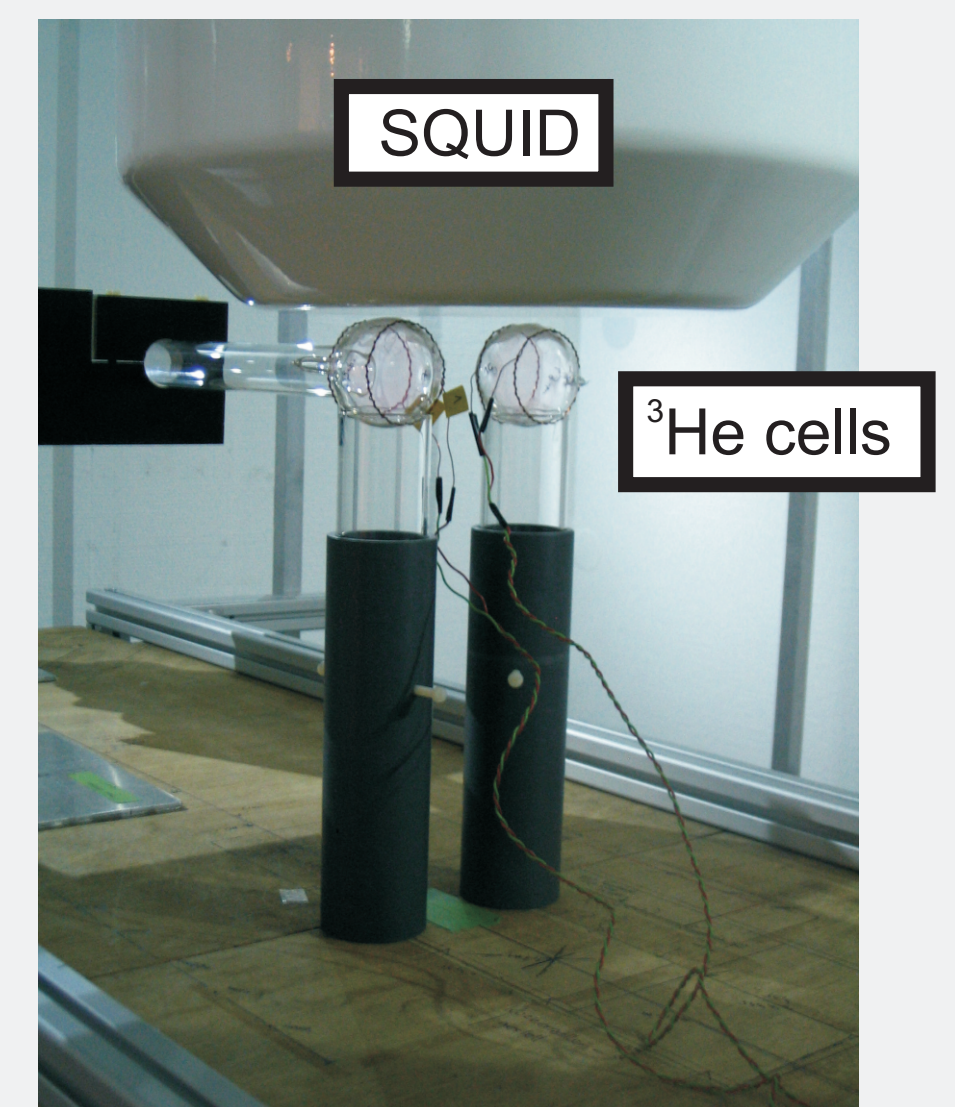


- * Glass cell filled with ^3He at $p = 5 \text{ mbar}$
- * Material: aluminosilicate glass Ge180
- * Size: 6 cm diameter
- * Gas discharge needed for Metastability Exchange Optical Pumping (MEOP)



Principle of the experiment:

- * ^3He gas is polarized by MEOP along B_1
- * Non-Adiabatic Field Switch ($\pi/2$ flip)
- * Spin precession around B_2

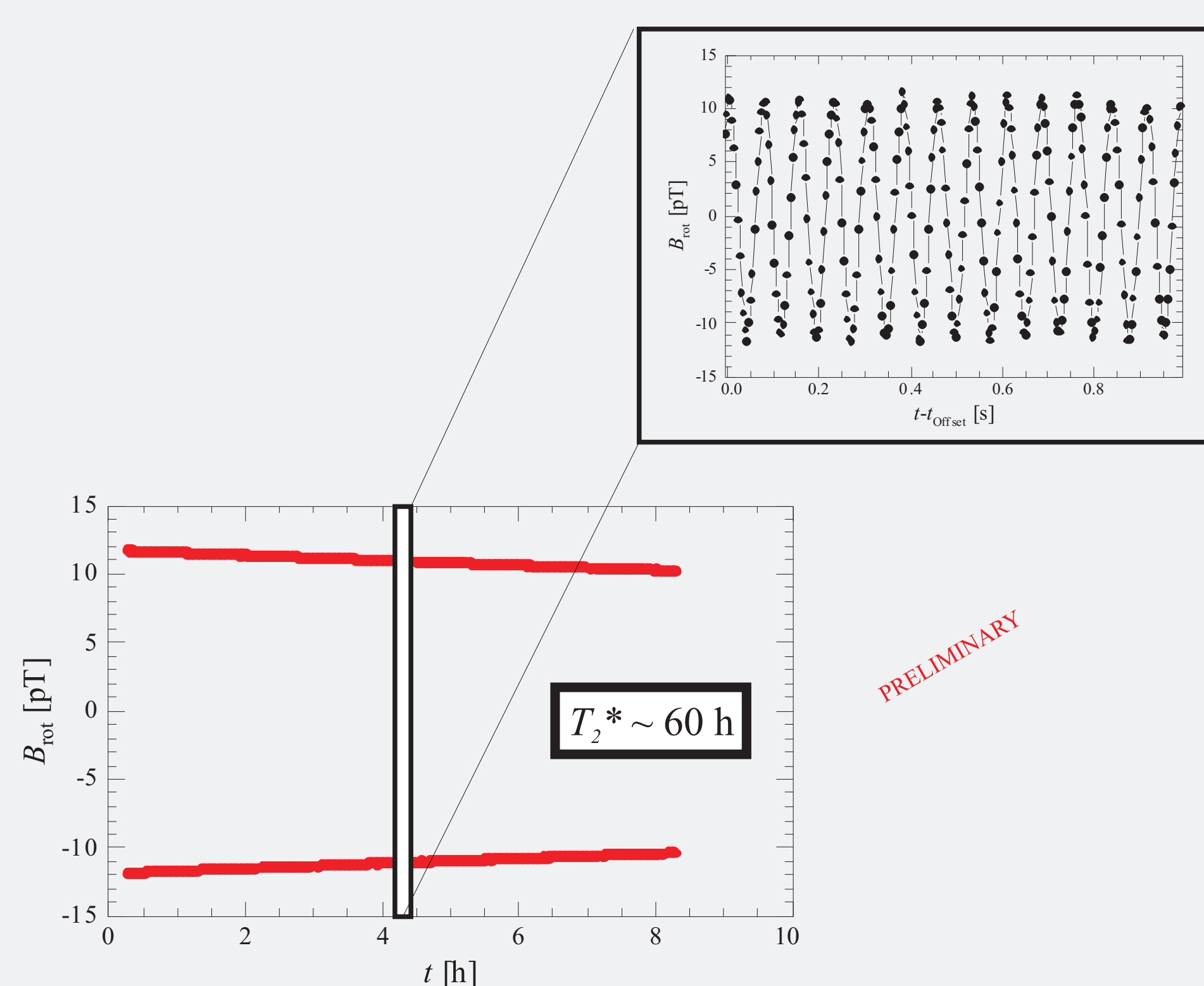


Detection of the free precession signal using a Low- T_c SQUID gradiometer.

In the picture, two cells were used to measure the magnetic field gradient between the cells.

Measurement of Precession

Two SQUIDS which are combined to form a gradiometer in order to decrease the background. The difference of the two is our signal. In a setup with one cell we get:



Measured signal strength:

$$B_{\text{rot}} \sim 13 \text{ pT}$$

Extracted transverse relaxation time T_2^* :

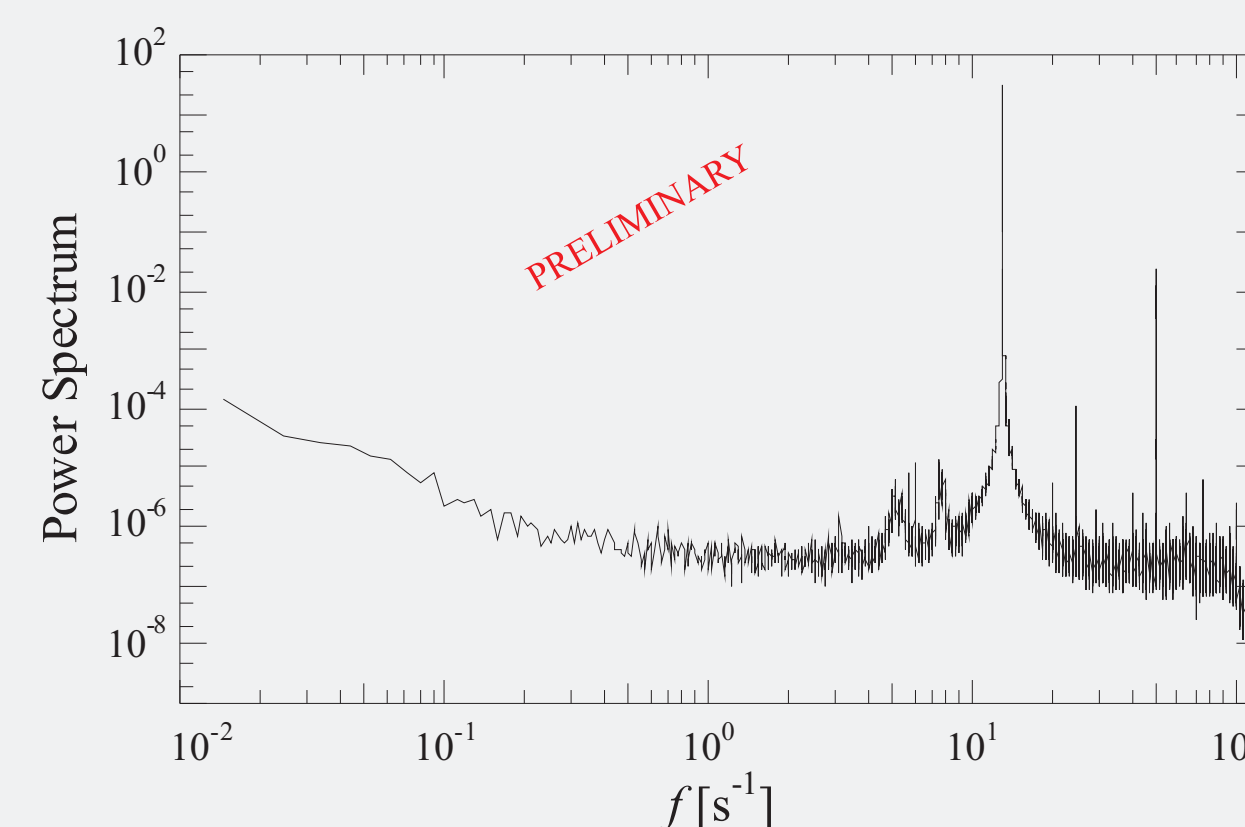
$$T_2^* \sim 60 \text{ h}$$

Statistical Sensitivity

The statistical precision is given by the so-called Cramér-Rao-Lower Bound, which is for a pure sinusoidal signal and white noise:

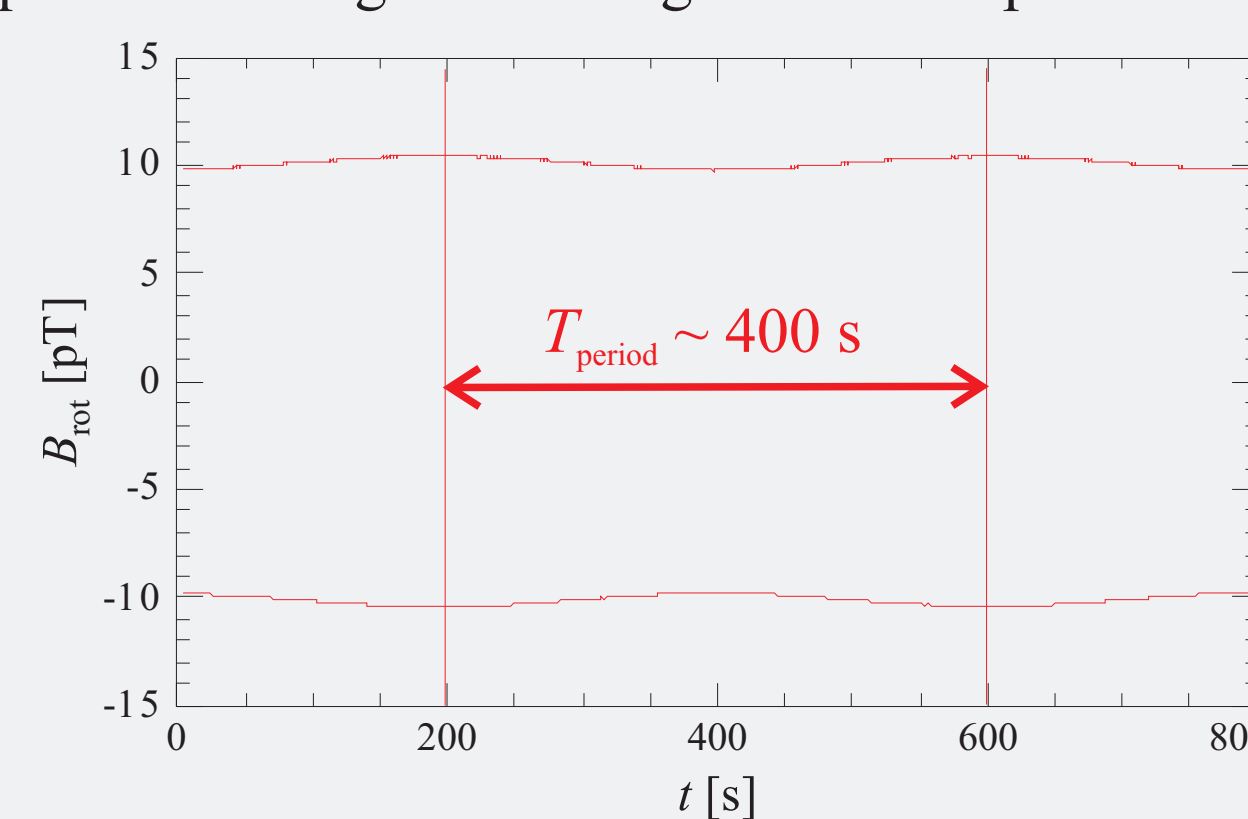
$$\delta B = \frac{\sqrt{6}}{\gamma_{\text{He}} \cdot \pi \cdot \text{SNR}} \cdot \frac{1}{\sqrt{T^3}}$$

With a SNR of about $3000/\sqrt{\text{Hz}}$ a field sensitivity of 1 fT is reached within $T \sim 100$ seconds.



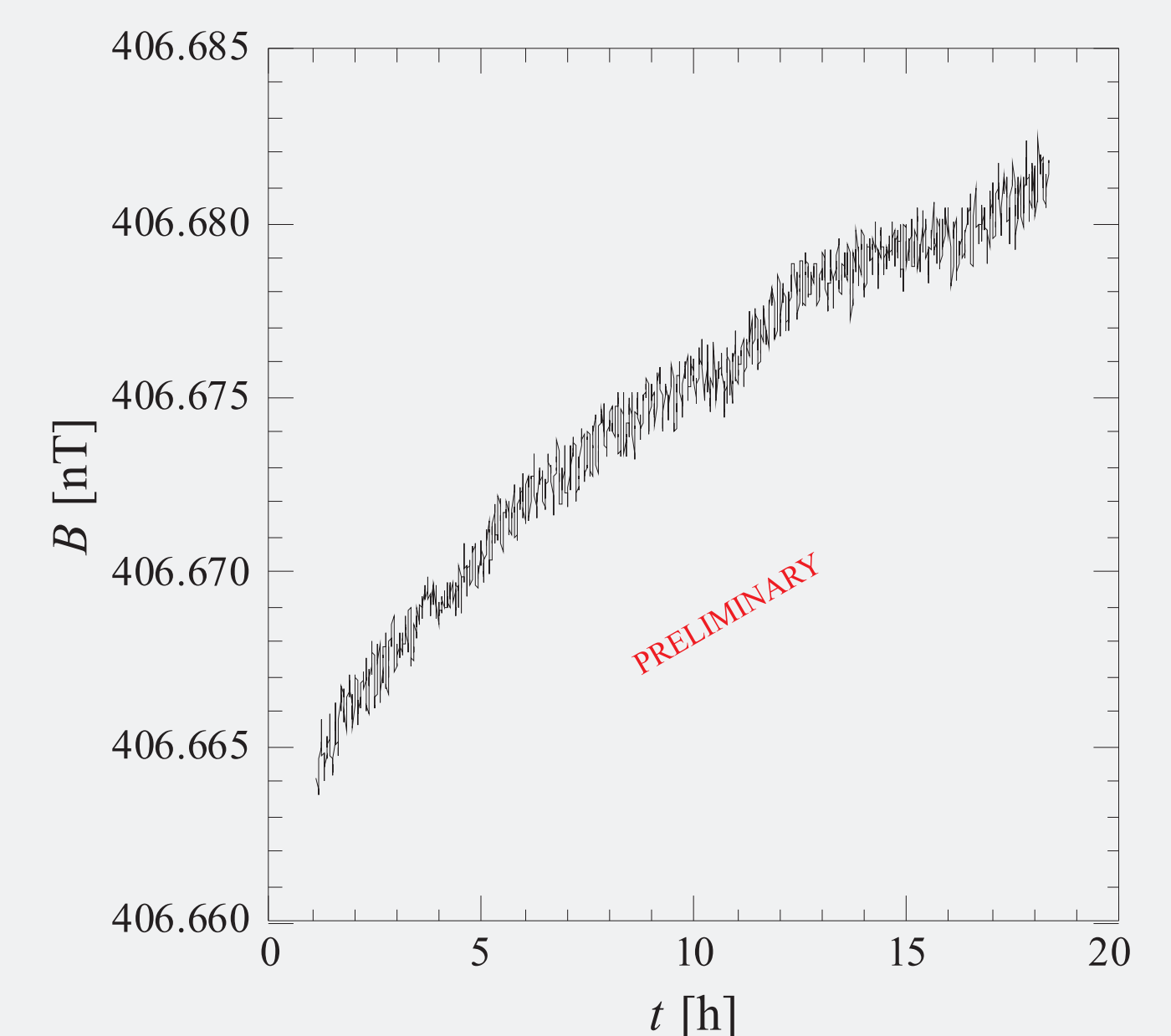
Measured Field Gradient

Measurement with 2 cells at a distance of 17.5 cm allow gradient measurements. The result, $T_{\text{period}} \sim 400$ s, corresponds to a magnetic field gradient of 4 pT/cm.



Magnetic field drifts and Comagnetometry

The variation of the measured magnetic field with time is shown in the picture below. We find that we have a magnetic field drift of about 1 pT/h. This number was independently confirmed in a direct measurement with the SQUID.



Problem: Magnetic field drift in BMSR-2 still too high. Solution:

- * Use of a comagnetometer (^3He , ^{129}Xe in the same volume)
- * ^3He , ^{129}Xe simultaneously pumped by spin exchange optical pumping (SEOP)
- * $T_2^*(^{129}\text{Xe})$ of 8000 s already demonstrated (Kilian et al.) \rightarrow improvements needed
- * Alternative: ^3He , ^{133}Cs comagnetometer: A very sensitive cesium magnetometer has been developed in the group of A. Weis at U. Fribourg. We don't expect the helium and the cesium to influence each other.